

# Developments on a massive planar pentabox with a differential equation method

C. Wever, in collaboration with C. Papadopoulos and D. Tommasini  
INPP, NCSR Demokritos

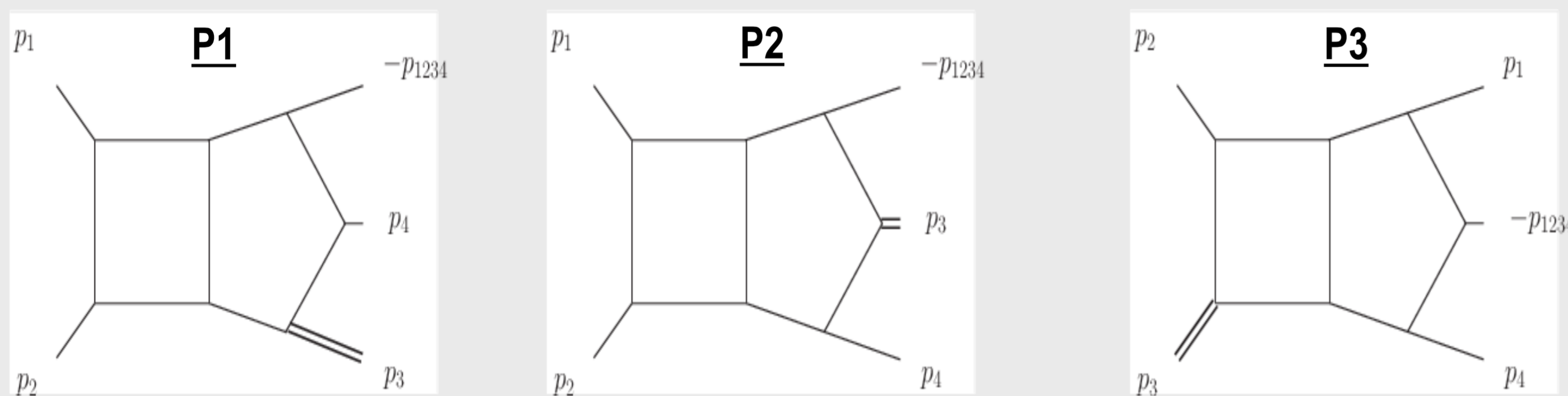


## 1. Motivation, introduction and goal

- Multi-loop calculations required for precision physics
  - NLO automation thanks to on-shell reduction methods
  - Next step: **NNLO** automation
- A finite basis of Master Integrals exists at two-loops:

$$\mathcal{A}^{2\text{-loop}} = \sum_{11\text{-prop}} \text{diagram} + \dots + \sum_{2\text{-prop}} \text{diagram} + \mathcal{R}$$

- Reduction to MI used for specific processes:
  - Integration by parts (IBP)** [1]
  - Missing ingredient: library of Master integrals (MI)**
- Interested next in two-loop, five-point diagrams with one external mass and massless propagators
  - Relevant e.g. for virtual-virtual contribution to  $2 \rightarrow 3$  LHC processes such as  $H + 2j, V + 2j, Vb\bar{b}$  (Les Houches Wishlist) at NNLO QCD
  - Three planar families:



**Goal:** compute all planar five-point MI with one external mass and massless internal propagators

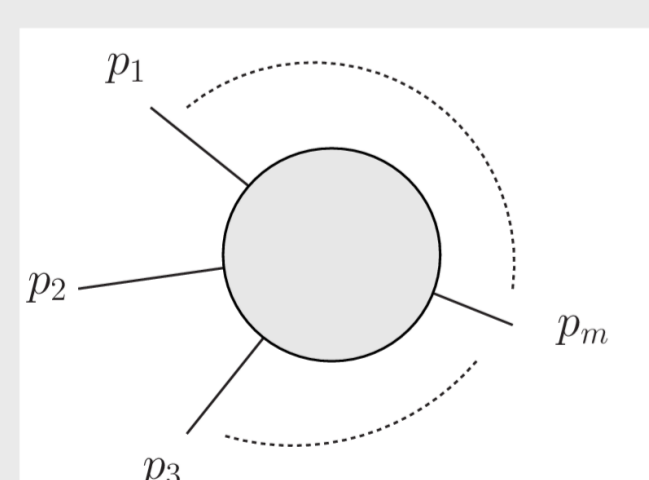
## 2. Theoretical basis: Simplified Differential Equations (SDE)

- Introduce auxiliary  $x$  in the denominators of loop integral [2]
- $x$ -parameter describes off-shellness of (some) external legs:

**Massive** legs:  $p_1 \xrightarrow{x\text{-parametrize}} p_1 + (1-x)q$

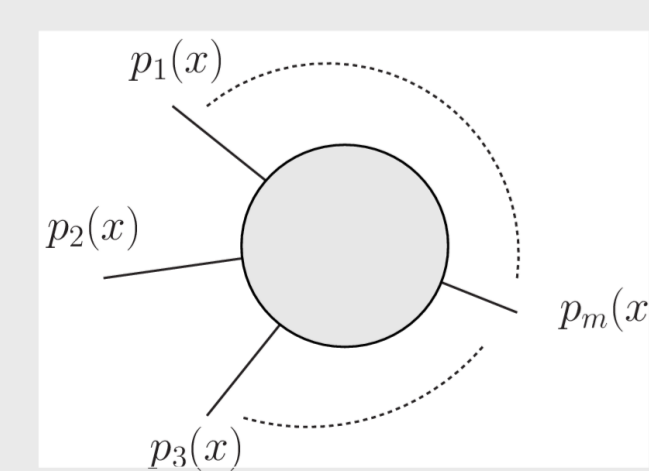
**Massless** legs:  $p_2 \xrightarrow{x=1} p_2 \text{ or } xp_2$

**General case:**



$$G_{a_1 \dots a_n}(s, \epsilon) = \int \left( \prod_i d^d k_i \right) \frac{1}{D_1^{2a_1}(k, p) \dots D_n^{2a_n}(k, p)}$$

$$x=1 \quad \begin{cases} p_i(x) = p_i + (1-x)q_i \\ \sum_i q_i = 0 \end{cases} \quad x\text{-parametrize}$$



$$G_{a_1 \dots a_n}(x, s, \epsilon) = \int \left( \prod_i d^d k_i \right) \frac{1}{D_1^{2a_1}(k, p(x)) \dots D_n^{2a_n}(k, p(x))}$$

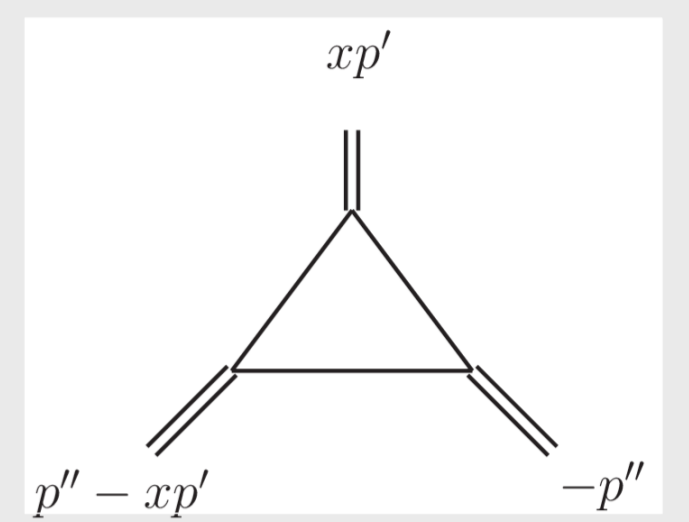
$$D_i(k, p) = c_{ij}k_j + d_{ij}p_j, \quad s = \{p_i \cdot p_j\}_{i,j}$$

$$\frac{\partial}{\partial x} \vec{G}^{MI}(x, s, \epsilon) \stackrel{IBP}{=} \vec{M}(x, s, \epsilon) \cdot \vec{G}^{MI}(x, s, \epsilon), \quad s = \{p_i \cdot p_j\}_{i,j}$$

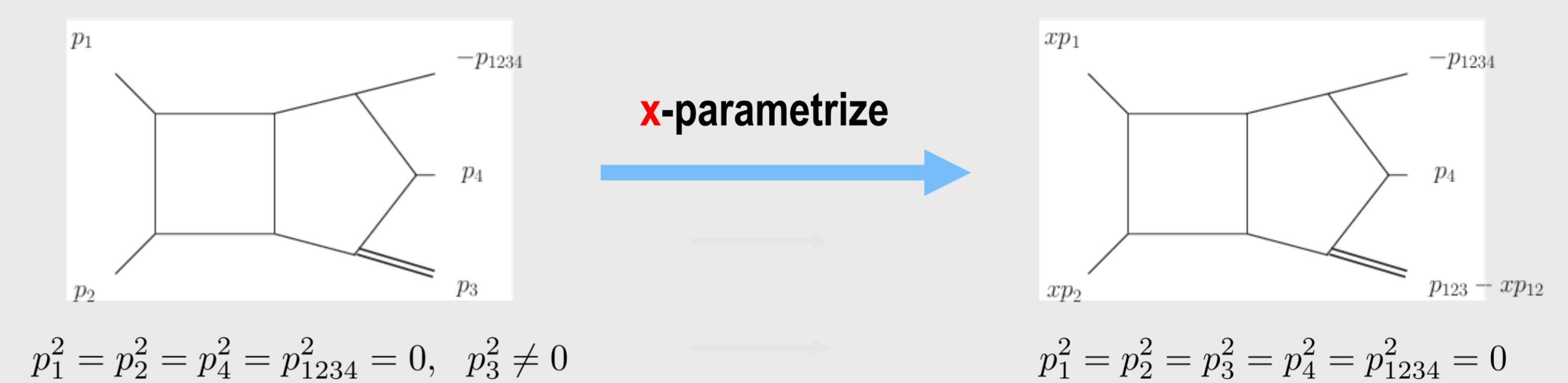
## 3. SDE method for the planar pentabox

Main criteria for choice of  $x$ -parametrization: **require Goncharov Polylog (GP) solution for DE**

In practice enough to choose the external legs such that the corresponding massive MI triangles (found by pinching external legs) are as follows:



**$x$ -parametrization for P1 family (74 MI in total):**



$$G_{a_1 \dots a_{11}}^{(P1)}(x) := \int \frac{d^d k_1 d^d k_2}{i\pi^{d/2} i\pi^{d/2}} \frac{1}{k_1^{2a_1} (k_1 + xp_1)^{2a_2} (k_1 + xp_{12})^{2a_3} (k_1 + p_{123})^{2a_4} (k_1 + p_{1234})^{2a_5}} \times \frac{1}{k_2^{2a_6} (k_2 - xp_1)^{2a_7} (k_2 - xp_{12})^{2a_8} (k_2 - p_{123})^{2a_9} (k_2 - p_{1234})^{2a_{10}} (k_1 + k_2)^{2a_{11}}}$$

- DE for P1 are known and integration underway in terms of GP's
- Reduction for P2 done (75 MI in total), P3 underway (bottleneck)

**Boundary term:** Integrands contain *branch points or poles* at  $x = \{x_1, x_2, \dots, \infty\}$  of form  $(x - x_i)^{m+n\epsilon}$

**Observation:** Boundary term always captured by integration from  $x = 0$  or appropriate  $x_i$

- Massless  $x = 1$  limit captured by resumming logs of  $(1 - x)$

## 4. Preliminary results for P1 and Outlook

	P1	$\epsilon^{-4}$	$\epsilon^{-3}$	$\epsilon^{-2}$	$\epsilon^{-1}$	$\epsilon^0$
Remainders:	$G_{1111000-11111}$	✓	✓	✓	✓	✓
	$G_{1111000011111}$	✓	✓	✓	✓	✓
	$G_{1111-100011111}$	✓	✓	✓	✓	
	$G_{1111001-11111}$	✓	✓	✓	✓	
	$G_{1111001011111}$	✓	✓	✓	✓	
	$G_{1111-101011111}$	✓	✓	✓	✓	

- Solutions expressed in terms of GPs with argument  $x$ :

$$G_{1111001011111}^{(P1)}(x, s, \epsilon) = \frac{C_0(s, \epsilon)}{x^4 (s_{12}(x-1)(x-s_{23}+s_{45}+s_{51})-s_{45}) + s_{34}(s_{45}(x-1)-s_{23}x)-s_{45}s_{51}(x-1)} \times \left\{ x^{-4\epsilon} C_1(x, s, \epsilon) + x^{-3\epsilon} C_2(x, s, \epsilon) + \frac{C_3(x, s)}{\epsilon^4} + \frac{1}{\epsilon^3} (GP(0;x)C_4(x, s) + GP(1;x)C_5(x, s) + GP(s_{45}/s_{12};x)C_6(x, s) \dots) + \dots \right\}$$

In **Euclidean region** agreement with SecDec [3]:

$$x = 1/13, \quad s_{12} = -2, \quad s_{23} = -3, \quad s_{34} = -5, \quad s_{45} = -7, \quad s_{51} = -11$$

**Analytical:**  $G_{1111001011111}^{(P1)} = \frac{1307.56}{\epsilon^4} + \frac{7834.53}{\epsilon^3} + \frac{22985.4}{\epsilon^2} + \frac{\dots}{\epsilon} + \dots + \mathcal{O}(\epsilon)$

**SecDec:**  $G_{1111001011111}^{(P1)} = \frac{1307.56}{\epsilon^4} + \frac{7833.34}{\epsilon^3} + \frac{22972.4}{\epsilon^2} + \frac{59772.6}{\epsilon} + 186628 + \mathcal{O}(\epsilon)$

**Outlook and Summary:**

- In progress: two-loop pentaboxes with one massive leg
- SDE method captures boundary terms by choosing the boundary at an appropriate branch point or pole

## 5. References

- Tkachov '81, Chetyrkin & Tkachov '81
- Papadopoulos '14, Papadopoulos, Tommasini, CW '14
- Borowka, Heinrich et al '11-'15