

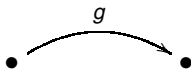
Higher Gauge Theory, Division Algebras and Superstrings

John Baez and John Huerta

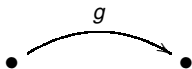
June 16, 2011
Quantum Theory and Gravitation

for more, see:
<http://math.ucr.edu/home/baez/susy/>

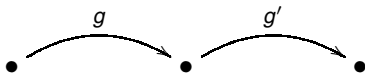
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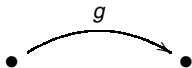
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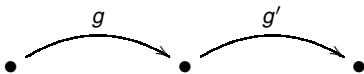
since composition of paths then corresponds to multiplication:



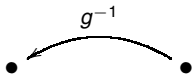
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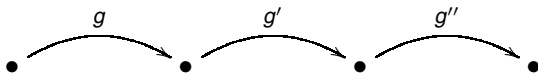
since composition of paths then corresponds to multiplication:



while reversing the direction corresponds to taking the inverse:



The associative law makes the holonomy along a triple composite unambiguous:



So: the topology dictates the algebra!

Higher gauge theory describes the parallel transport not only of point particles, but also higher-dimensional extended objects.

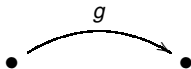
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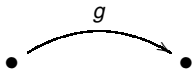
A '2-group' has objects:



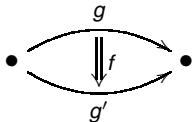
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A '2-group' has objects:



but also morphisms:



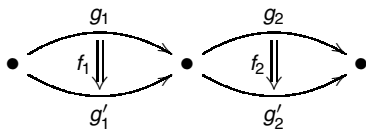
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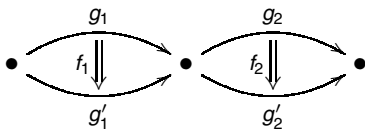
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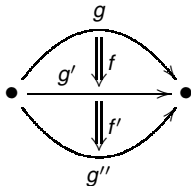
We can multiply objects:



multiply morphisms:



and also compose morphisms:



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Just as a group is a monoid where every element has an inverse, a **2-group** is a monoidal category where every object and every morphism has an inverse.

For higher gauge theory, we really want ‘Lie 2-groups’. By now there is an extensive theory of these.

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A spin foam model based on this Lie 2-group may serve as a ‘quantum model of flat 4d spacetime’, much as the Ponzano–Regge model does for 3d spacetime. See:

- Aristide Baratin and Derek Wise, 2-group representations of spin foams, arXiv:0910.1542.

for the Euclidean case.

Other examples show up in string theory. In his thesis, John Huerta showed that they explain this pattern:

- The only normed division algebras are \mathbb{R} , \mathbb{C} , \mathbb{H} and \mathbb{O} . They have dimensions $k = 1, 2, 4$ and 8 .
- The classical superstring makes sense only in dimensions $k + 2 = 3, 4, 6$ and 10 .
- The classical super-2-brane makes sense only in dimensions $k + 3 = 4, 5, 7$ and 11 .

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For superstrings we need Lie 2-supergroups; for super-2-branes we need Lie **3**-supergroups.

To get our hands on Lie n -supergroups, it's easiest to start with 'Lie n -superalgebras'. Let's see what those are.

An L_∞ -**algebra** is a chain complex

$$L_0 \xleftarrow{d} L_1 \xleftarrow{d} \cdots \xleftarrow{d} L_n \xleftarrow{d} \cdots$$

equipped with the structure of a Lie algebra ‘up to coherent chain homotopy’.

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- a map $d: L \rightarrow L$ of grade -1 with $d^2 = 0$

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Which Lie 2-superalgebras, if any, are relevant to superstrings?

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(Don’t worry, soon I’ll tell you what $\mathfrak{siso}(T)$ actually *is!*)

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Let's see how it works in detail.

If V is a finite-dimensional real vector space with a quadratic form Q , the **Clifford algebra** $\text{Cliff}(V)$ is the real associative algebra generated by V with relations

$$v^2 = -Q(v)$$

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In the first case $S_+ \not\cong S_-$. In the second, set $S_+ = S_- = S$. In either case, let's call S_+ and S_- **left- and right-handed spinors**.

$\text{Cliff}_0(V)$ acts on S_+ and S_- . But the whole Clifford algebra acts on $S_+ \oplus S_-$, with odd elements interchanging the two parts. So, we can ‘multiply’ a spinor by a vector and get a spinor of the other handedness:

$$\therefore V \otimes S_+ \rightarrow S_-$$

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So, let's see when $\dim(V) = \dim(S_+) = \dim(S_-)$.

V	S_{\pm}	normed division algebra?
\mathbb{R}^1	\mathbb{R}	YES: \mathbb{R}
\mathbb{R}^2	\mathbb{C}	YES: \mathbb{C}
\mathbb{R}^3	\mathbb{H}	NO
\mathbb{R}^4	\mathbb{H}	YES: \mathbb{H}
\mathbb{R}^5	\mathbb{H}^2	NO
\mathbb{R}^6	\mathbb{C}^2	NO
\mathbb{R}^7	\mathbb{R}^8	NO
\mathbb{R}^8	\mathbb{R}^8	YES: \mathbb{O}

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\mathbb{R}^4	\mathbb{H}	YES: \mathbb{H}
\mathbb{R}^5	\mathbb{H}^2	NO
\mathbb{R}^6	\mathbb{C}^2	NO
\mathbb{R}^7	\mathbb{R}^8	NO
\mathbb{R}^8	\mathbb{R}^8	YES: \mathbb{O}

Increasing k by 8 multiplies $\dim(S_{\pm})$ by 16, so these are the *only* normed division algebras!

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Now vectors V are the 2×2 Hermitian matrices with entries in \mathbb{K} :

$$V = \left\{ \begin{pmatrix} t+x & \bar{y} \\ y & t-x \end{pmatrix} : t, x \in \mathbb{R}, y \in \mathbb{K} \right\}.$$

Now our quadratic form Q comes from the determinant:

$$\det \begin{pmatrix} t+x & \bar{y} \\ y & t-x \end{pmatrix} = t^2 - x^2 - |y|^2$$

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Concretely:

$$[\psi, \phi] = \psi\phi^\dagger + \phi\psi^\dagger$$

It’s symmetric!

So, we can define the **translation Lie superalgebra**

$$T = V \oplus S_+$$

with V as its even part and S_+ as its odd part. We define the bracket to be zero except for $[-, -]: S_+ \otimes S_+ \rightarrow V$. The Jacobi identity holds trivially.

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$\text{Spin}(V)$ acts on everything, and its Lie algebra is $\mathfrak{so}(V)$, so we can form the semidirect product

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The corresponding Lie supergroup acts as symmetries of 'Minkowski superspace'.

We get a Poincaré Lie superalgebra whenever we have an invariant symmetric bracket that takes two spinors and gives a vector. What's so special about the dimensions 3, 4, 6 and 10?

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In fact, *only* for Minkowski spacetimes of dimension 3, 4, 6, and 10 does this identity hold!

The identity $[\psi, \psi] \cdot \psi = 0$ lets us extend the Poincaré Lie superalgebra $\mathfrak{siso}(T)$ to a Lie 2-superalgebra

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The idea:

- d is zero
- $[-, -]$ is zero except for the bracket in $\mathfrak{siso}(T)$
- $[-, -, -]$ is zero unless two arguments are spinors and one is a vector in $\mathfrak{siso}(T)$, and

$$[\psi, \phi, \nu] = g([\psi, \phi], \nu) \in \mathbb{R}$$

where $\psi, \phi \in S_+$, $\nu \in V$, and $g: V \otimes V \rightarrow R$ is the **Minkowski metric**: the bilinear form corresponding to Q .

To get a Lie 2-superalgebra this way, the ternary bracket must obey an equation. This says that

$$[-, -, -]: \mathfrak{siso}(T)^{\otimes 3} \rightarrow \mathbb{R}$$

is a *3-cocycle* in Lie superalgebra cohomology.

To get a Lie 2-superalgebra this way, the ternary bracket must obey an equation. This says that

$$[-, -, -]: \mathfrak{siso}(T)^{\otimes 3} \rightarrow \mathbb{R}$$

is a *3-cocycle* in Lie superalgebra cohomology.

The equation

$$[\psi, \psi] \cdot \psi = 0$$

is this cocycle condition in disguise.

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Theorem (John Huerta)

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This theorem takes real work to prove. Not every Lie 2-superalgebra has a corresponding ‘Lie 2-supergroup’ in such a simple-minded sense! There are important finite-dimensional Lie 2-algebras that don’t come from 2-groups in the category of manifolds—instead, they come from ‘stacky’ Lie 2-groups.

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Theorem (Huerta)

In Minkowski spacetimes of dimensions 4, 5, 7 and 11, we can use division algebras to construct a 4-cocycle on the Poincaré Lie superalgebra.

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*Moreover, there is a 3-group in the category of supermanifolds, **2-Brane**(T), whose Lie 3-superalgebra is **2-brane**(T).*

2-Brane(T) is relevant to the theory of supersymmetric 2-branes in dimension 4, 5, 7 and 11. And so, the octonionic case is relevant to M -theory (whatever that is).

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Summary: For theories that include gravity and describe the parallel transport of extended objects, we want Lie n -groups extending the Lorentz or Poincaré group.

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Summary: For theories that include gravity and describe the parallel transport of extended objects, we want Lie n -groups extending the Lorentz or Poincaré group. Lie n -*supergroups* extending the Poincaré *supergroup* only exist in special dimensions, thanks to special properties of the normed division algebras. In 10 and 11 dimensions, the octonions play a crucial role.

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For more see:

- John Huerta, *Division Algebras, Supersymmetry and Higher Gauge Theory*, at <http://math.ucr.edu/home/baez/susy/>