### Pure Spinor Helicity Methods

**Rutger Boels** 

Niels Bohr International Academy, Copenhagen

R.B., arXiv:0908.0738 [hep-th]

Rutger Boels (NBIA)

Pure Spinor Helicity Methods

15th String Workshop, Zürich 1 / 25

The greatest common denominator of the audience?

The greatest common denominator of the audience in books?

The greatest common denominator of the audience in books:



・ 同 ト ・ ヨ ト ・ ヨ ト ・ ヨ

The greatest common denominator of the audience in science books?





12

イロト イポト イヨト イヨト ヨ

# Why you should pay attention

### subject:

calculation of scattering amplitudes in D > 4 with many legs

Rutaer	Bool	e (NIRIA)
ruugei	Duci	

# Why you should pay attention

### subject:

calculation of scattering amplitudes in D > 4 with many legs

pure spinor helicity methods:

precise control over Poincaré and Susy quantum numbers

# Why you should pay attention

### subject:

calculation of scattering amplitudes in D > 4 with many legs

### pure spinor helicity methods:

precise control over Poincaré and Susy quantum numbers for all legs simultaneously

(日) (周) (王) (王) (王)

# Outline

# **Motivation**

### Covariant representation theory of 2

- Poincaré algebra ۲
- Spin algebra ٠
- Susy algebra ۲



3 -

### Every talk on amplitudes should mention ....

Rutger Boels (NBIA)

Pure Spinor Helicity Methods

15th String Workshop, Zürich 5 / 25

### Every talk on amplitudes should mention ....



Rutger Boels (NBIA)

Pure Spinor Helicity Methods

15th String Workshop, Zürich 5 / 25

• loops in four dimensions  $\leftarrow$  dimensional regularization

Rutger Boels (NBIA)

Pure Spinor Helicity Methods

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □

- loops in four dimensions  $\leftarrow$  dimensional regularization
- one-loop: pure Yang-Mills [Giele, Kunszt, and Melnikov, 08]
  - uses 6D trees
- high loops:  $\mathcal{N} = 4$ ,  $\mathcal{N} = 8$  [Bern, Dixon, Kosower et. al.], [others]
  - uses 10D, 11D trees

- loops in four dimensions  $\leftarrow$  dimensional regularization
- one-loop: pure Yang-Mills [Giele, Kunszt, and Melnikov, 08]
  - uses 6D trees
- high loops:  $\mathcal{N} = 4$ ,  $\mathcal{N} = 8$  [Bern, Dixon, Kosower et. al.], [others]
  - uses 10D, 11D trees
- recent quantum leaps in four: what is special about four?
- string theory (analytic S-matrix type techniques)

- loops in four dimensions  $\leftarrow$  dimensional regularization
- one-loop: pure Yang-Mills [Giele, Kunszt, and Melnikov, 08]
  - uses 6D trees
- high loops:  $\mathcal{N} = 4$ ,  $\mathcal{N} = 8$  [Bern, Dixon, Kosower et. al.], [others]
  - uses 10D, 11D trees
- recent quantum leaps in four: what is special about four?
- string theory (analytic S-matrix type techniques)
- little is known

- $\bullet$  loops in four dimensions  $\leftarrow$  dimensional regularization
- one-loop: pure Yang-Mills [Giele, Kunszt, and Melnikov, 08]
  - uses 6D trees
- high loops:  $\mathcal{N} = 4$ ,  $\mathcal{N} = 8$  [Bern, Dixon, Kosower et. al.], [others]
  - uses 10D, 11D trees
- recent quantum leaps in four: what is special about four?
- string theory (analytic S-matrix type techniques)
- little is known

### is there a D > 4 analogue of:

- MHV amplitudes?
- (any recent buzzword in *D* = 4?)

- $\bullet$  loops in four dimensions  $\leftarrow$  dimensional regularization
- one-loop: pure Yang-Mills [Giele, Kunszt, and Melnikov, 08]
  - uses 6D trees
- high loops:  $\mathcal{N} = 4$ ,  $\mathcal{N} = 8$  [Bern, Dixon, Kosower et. al.], [others]
  - uses 10D, 11D trees
- recent quantum leaps in four: what is special about four?
- string theory (analytic S-matrix type techniques)
- little is known

### is there a D > 4 analogue of:

- MHV amplitudes?
- (any recent buzzword in *D* = 4?)
- pure spinor spaces are higher dimensional twistor spaces

# Outline

### Covariant representation theory of 2

- Poincaré algebra ۲

- higher D massless vectors decompose into massive D = 4
- massive spinor helicity: e.g. [Kleiss, Stirling, 85], [Dittmaier, 98]

- higher D massless vectors decompose into massive D = 4
- massive spinor helicity: e.g. [Kleiss, Stirling, 85], [Dittmaier, 98]
- maximal set of commuting operators:  $k_{\mu}$ ,  $W_{\mu}$

- higher D massless vectors decompose into massive D = 4
- massive spinor helicity: e.g. [Kleiss, Stirling, 85], [Dittmaier, 98]
- maximal set of commuting operators:  $k_{\mu}$ ,  $W_{\mu}$
- $k_{\mu}k^{\mu} = m^2$ , pick spin axis through light-like vector q

$$R_z = R_q^1 = rac{q^\mu W_\mu}{2q \cdot k} = rac{\epsilon_{\mu
u
ho\sigma}q^\mu k^
u \Sigma^{
ho\sigma}}{2q \cdot k}$$

- higher D massless vectors decompose into massive D = 4
- massive spinor helicity: e.g. [Kleiss, Stirling, 85], [Dittmaier, 98]
- maximal set of commuting operators:  $k_{\mu}$ ,  $W_{\mu}$
- $k_{\mu}k^{\mu} = m^2$ , pick spin axis through light-like vector q

$$\boldsymbol{R}_{\boldsymbol{z}} = \boldsymbol{R}_{\boldsymbol{q}}^{1} = \frac{q^{\mu} \boldsymbol{W}_{\mu}}{2\boldsymbol{q} \cdot \boldsymbol{k}} = \frac{\epsilon_{\mu\nu\rho\sigma} q^{\mu} \boldsymbol{k}^{\nu} \boldsymbol{\Sigma}^{\rho\sigma}}{2\boldsymbol{q} \cdot \boldsymbol{k}}$$

•  $\exists$  vectors  $n_1$  and  $n_2$  such that  $q, \hat{q}, n_1, n_2$  span  $R^{1,3}, q \cdot n_i = 0$ 

- higher D massless vectors decompose into massive D = 4
- massive spinor helicity: e.g. [Kleiss, Stirling, 85], [Dittmaier, 98]
- maximal set of commuting operators:  $k_{\mu}$ ,  $W_{\mu}$
- $k_{\mu}k^{\mu} = m^2$ , pick spin axis through light-like vector q

$$\boldsymbol{R}_{\boldsymbol{z}} = \boldsymbol{R}_{\boldsymbol{q}}^{1} = \frac{q^{\mu} \boldsymbol{W}_{\mu}}{2\boldsymbol{q} \cdot \boldsymbol{k}} = \frac{\epsilon_{\mu\nu\rho\sigma} q^{\mu} \boldsymbol{k}^{\nu} \boldsymbol{\Sigma}^{\rho\sigma}}{2\boldsymbol{q} \cdot \boldsymbol{k}}$$

•  $\exists$  vectors  $n_1$  and  $n_2$  such that  $q, \hat{q}, n_1, n_2$  span  $R^{1,3}, q \cdot n_i = 0$ 

### massive polarization vectors

$$e_{\pm}^{\mu} = \frac{1}{\sqrt{2}} \left( n_{1}^{\mu} \pm i n_{2}^{\mu} - q^{\mu} \frac{((n_{1} \pm i n_{2}) \cdot k)}{q \cdot k} \right) \qquad e_{0}^{\mu} = \frac{k^{\mu}}{m} - \frac{m q^{\mu}}{q \cdot k}$$

- higher D massless vectors decompose into massive D = 4
- massive spinor helicity: e.g. [Kleiss, Stirling, 85], [Dittmaier, 98]
- maximal set of commuting operators:  $k_{\mu}$ ,  $W_{\mu}$
- $k_{\mu}k^{\mu} = m^2$ , pick spin axis through light-like vector q

$$\boldsymbol{R}_{\boldsymbol{z}} = \boldsymbol{R}_{\boldsymbol{q}}^{1} = \frac{q^{\mu} \boldsymbol{W}_{\mu}}{2\boldsymbol{q} \cdot \boldsymbol{k}} = \frac{\epsilon_{\mu\nu\rho\sigma} q^{\mu} \boldsymbol{k}^{\nu} \boldsymbol{\Sigma}^{\rho\sigma}}{2\boldsymbol{q} \cdot \boldsymbol{k}}$$

•  $\exists$  vectors  $n_1$  and  $n_2$  such that  $q, \hat{q}, n_1, n_2$  span  $R^{1,3}, q \cdot n_i = 0$ 

### massive polarization vectors

$$\boldsymbol{e}_{\pm}^{\mu} = \frac{1}{\sqrt{2}} \left( \boldsymbol{n}_{1}^{\mu} \pm \mathrm{i} \boldsymbol{n}_{2}^{\mu} - \boldsymbol{q}^{\mu} \frac{\left( (\boldsymbol{n}_{1} \pm \mathrm{i} \boldsymbol{n}_{2}) \cdot \boldsymbol{k} \right)}{\boldsymbol{q} \cdot \boldsymbol{k}} \right) \qquad \boldsymbol{e}_{0}^{\mu} = \frac{k^{\mu}}{m} - \frac{m \boldsymbol{q}^{\mu}}{\boldsymbol{q} \cdot \boldsymbol{k}}$$

- higher D massless vectors decompose into massive D = 4
- massive spinor helicity: e.g. [Kleiss, Stirling, 85], [Dittmaier, 98]
- maximal set of commuting operators:  $k_{\mu}$ ,  $W_{\mu}$
- $k_{\mu}k^{\mu} = m^2$ , pick spin axis through light-like vector q

$$\boldsymbol{R}_{z} = \boldsymbol{R}_{q}^{1} = \frac{q^{\mu} W_{\mu}}{2q \cdot k} = \frac{\epsilon_{\mu\nu\rho\sigma} q^{\mu} k^{\nu} \Sigma^{\rho\sigma}}{2q \cdot k}$$

•  $\exists$  vectors  $n_1$  and  $n_2$  such that  $q, \hat{q}, n_1, n_2$  span  $R^{1,3}, q \cdot n_i = 0$ 

massive polarization vectors

$$e^{\mu}_{\pm} = rac{1}{\sqrt{2}} \left( ilde{n}^{\mu}_1 \pm \mathrm{i} ilde{n}^{\mu}_2 
ight) (k) \qquad e^{\mu}_0 = rac{k^{\mu}}{m} - rac{mq^{\mu}}{q \cdot k}$$

- higher D massless vectors decompose into massive D = 4
- massive spinor helicity: e.g. [Kleiss, Stirling, 85], [Dittmaier, 98]
- maximal set of commuting operators:  $k_{\mu}$ ,  $W_{\mu}$
- $k_{\mu}k^{\mu} = m^2$ , pick spin axis through light-like vector q

$$\boldsymbol{R}_{z} = \boldsymbol{R}_{q}^{1} = \frac{q^{\mu} W_{\mu}}{2q \cdot k} = \frac{\epsilon_{\mu\nu\rho\sigma} q^{\mu} k^{\nu} \Sigma^{\rho\sigma}}{2q \cdot k}$$

•  $\exists$  vectors  $n_1$  and  $n_2$  such that  $q, \hat{q}, n_1, n_2$  span  $R^{1,3}, q \cdot n_i = 0$ 

massive polarization vectors

$$e^{\mu}_{\pm} = rac{1}{\sqrt{2}} \left( ilde{n}^{\mu}_1 \pm \mathrm{i} ilde{n}^{\mu}_2 
ight) (k) \qquad e^{\mu}_0 = rac{k^{\mu}}{m} - rac{mq^{\mu}}{q \cdot k}$$

• broken gauge theory in  $D = 4 \rightarrow [R.B., Christian Schwinn, to appear]$ 

• given  $k_{\mu}$ , little group is ISO(D-2)

• given  $k_{\mu}$ , little group is SO(D-2)

- given  $k_{\mu}$ , little group is SO(D-2)
- Pauli-Lubanski vector tensors  $k_{[\mu} \Sigma_{\nu \rho]}$

- given  $k_{\mu}$ , little group is SO(D-2)
- Pauli-Lubanski vector tensors  $k_{[\mu} \Sigma_{\nu \rho]}$
- from previous use:
  - choose q such that  $q^2 = 0$
  - choose  $q, \hat{q}, n_i$  to form an ortho-normal basis of  $R^{1,D-1}$

• construct 
$$\tilde{n}_i = n_i - q \frac{(n_i \cdot k)}{(q \cdot k)}$$

- given  $k_{\mu}$ , little group is SO(D-2)
- Pauli-Lubanski vector tensors  $k_{[\mu} \Sigma_{\nu \rho]}$
- from previous use:
  - choose q such that  $q^2 = 0$
  - choose  $q, \hat{q}, n_i$  to form an ortho-normal basis of  $R^{1,D-1}$
  - construct  $\tilde{n}_i = n_i q \frac{(n_i \cdot k)}{(q \cdot k)}$
- for representation theory choose Cartan generators as

$$R_q^j \equiv rac{1}{2} ilde{n}_{2j-1} ilde{n}_{2j} \Sigma = q n_1 n_2 (k_{[\mu} \Sigma_{
u 
ho]})$$

• eigenvalues form a weight vector  $\vec{h}$ :  $R_q^j e = h^j e$ ,  $\vec{h} = (0, ..., \pm 1, ...)$ 

- given  $k_{\mu}$ , little group is SO(D-2)
- Pauli-Lubanski vector tensors  $k_{[\mu} \Sigma_{\nu \rho]}$
- from previous use:
  - choose q such that  $q^2 = 0$
  - choose  $q, \hat{q}, n_i$  to form an ortho-normal basis of  $R^{1,D-1}$
  - construct  $\tilde{n}_i = n_i q \frac{(n_i \cdot k)}{(q \cdot k)}$
- for representation theory choose Cartan generators as

$$\mathcal{R}_q^j\equivrac{1}{2} ilde{n}_{2j-1} ilde{n}_{2j}\Sigma=qn_1n_2(k_{[\mu}\Sigma_{
u
ho]})$$

• eigenvalues form a weight vector  $\vec{h}$ :  $R_q^j e = h^j e$ ,  $\vec{h} = (0, ..., \pm 1, ...)$ 

polarization vectors (in q lightcone gauge)

$$e^{\mu}(0,\ldots,\pm_j,\ldots)=rac{1}{\sqrt{2}}\left( ilde{n}^{\mu}_{2j-1}\pm\mathrm{i} ilde{n}^{\mu}_{2j}
ight)(k)$$

• requires choice of complex structure

Rutger Boels (NBIA)

Pure Spinor Helicity Methods

소리 에 소문에 이 제 문어 소문에 드릴 것

- massive vectors: SO(D − 1)
- ∃ generalization to all integer spins (e.g. gravity)

- massive vectors: SO(D-1)
- ∃ generalization to all integer spins (e.g. gravity)
- q: choice of Cartan generators and choice of gauge
- *D* = 4: little group is Abelian

- massive vectors: SO(D − 1)
- ∃ generalization to all integer spins (e.g. gravity)
- q: choice of Cartan generators and choice of gauge
- D = 4: little group is Abelian
- can compare different legs:

$$oldsymbol{e}_{\mu}\left(oldsymbol{k}_{i},oldsymbol{\vec{h}}_{i}
ight)\cdotoldsymbol{e}^{\mu}\left(oldsymbol{k}_{j},oldsymbol{\vec{h}}_{j}
ight)=-\delta\left(oldsymbol{\vec{h}}_{i}+oldsymbol{\vec{h}}_{j}
ight)$$

- massive vectors: SO(D − 1)
- ∃ generalization to all integer spins (e.g. gravity)
- q: choice of Cartan generators and choice of gauge
- D = 4: little group is Abelian
- can compare different legs:

$$oldsymbol{e}_{\mu}\left(oldsymbol{k}_{i},oldsymbol{\vec{h}}_{i}
ight)\cdotoldsymbol{e}^{\mu}\left(oldsymbol{k}_{j},oldsymbol{\vec{h}}_{j}
ight)=-\delta\left(oldsymbol{\vec{h}}_{i}+oldsymbol{\vec{h}}_{j}
ight)$$

Application: higher dimensional YM amplitudes at tree level  $\langle (\pm 1, 0, \ldots)^{i_1} (0, \pm 1, \ldots)^{i_2} \dots (0, \ldots, \pm 1)^{i_{D/2}} \rangle = 0$
### Remarks and a quick application

- massive vectors: SO(D 1)
- ∃ generalization to all integer spins (e.g. gravity)
- q: choice of Cartan generators and choice of gauge
- D = 4: little group is Abelian
- can compare different legs:

$$oldsymbol{e}_{\mu}\left(oldsymbol{k}_{i},oldsymbol{\vec{h}}_{i}
ight)\cdotoldsymbol{e}^{\mu}\left(oldsymbol{k}_{j},oldsymbol{\vec{h}}_{j}
ight)=-\delta\left(oldsymbol{\vec{h}}_{i}+oldsymbol{\vec{h}}_{j}
ight)$$

Application: higher dimensional YM amplitudes at tree level

$$\langle \left(\pm 1,0,\ldots\right)^{i_1} (0,\pm 1,\ldots)^{i_2} \ldots (0,\ldots,\pm 1)^{i_{D/2}} \rangle = 0$$

- 'helicity equal' in D = 4
- 'one helicity unequal' does not vanish in D > 4
- class of Einstein gravity amplitudes through KLT
- six gluon open string amplitude [Oprisa, Stieberger, 2005]

#### Outline

#### Motivation

## Covariant representation theory of Poincaré algebra

- Spin algebra
- Susy algebra

#### 3 Outlook

- polarization vectors and q and k span  $\mathbb{R}^{1,D-1}$
- needed: similar basis of spinors
- would like: vectors in terms of spinors

▲□ ▶ ▲ ■ ▶ ▲ ■ ▶ ■ ■ ■ ● ● ●

- polarization vectors and q and k span  $\mathbb{R}^{1,D-1}$
- needed: similar basis of spinors
- would like: vectors in terms of spinors

guess:  $e^{\mu}(ec{h})\sim \overline{\xi}\gamma^{\mu}\psi$ 

want:

$$q_\mu e^\mu = k_\mu e^\mu = \ldots = 0$$

therefore:

$$\overline{\xi}\left(\pmb{q}_{\mu}\gamma^{\mu}
ight)=\left(\pmb{k}_{\mu}\gamma^{\mu}
ight)\psi=0 \quad ext{or} \quad \overline{\xi}\left(\pmb{k}_{\mu}\gamma^{\mu}
ight)=\left(\pmb{q}_{\mu}\gamma^{\mu}
ight)\psi=0$$

other inner products  $\rightarrow$  total  $\frac{D}{2}$  annihilation conditions

▲□ ▶ ▲ ■ ▶ ▲ ■ ▶ ■ ■ ■ ● ● ●

- polarization vectors and q and k span  $\mathbb{R}^{1,D-1}$
- needed: similar basis of spinors
- would like: vectors in terms of spinors

guess:  $e^{\mu}(ec{h})\sim \overline{\xi}\gamma^{\mu}\psi$ 

want:

$$q_\mu e^\mu = k_\mu e^\mu = \ldots = 0$$

therefore:

$$\overline{\xi}\left( \pmb{q}_{\mu}\gamma^{\mu}
ight) =\left( \pmb{k}_{\mu}\gamma^{\mu}
ight)\psi =0 \quad ext{or} \quad \overline{\xi}\left( \pmb{k}_{\mu}\gamma^{\mu}
ight) =\left( \pmb{q}_{\mu}\gamma^{\mu}
ight)\psi =0$$

other inner products  $\rightarrow$  total  $\frac{D}{2}$  annihilation conditions

• need spinors annihilated by  $\frac{D}{2}$  generators  $v^i_\mu \gamma^\mu$  such that  $\langle v^i, v^j \rangle = 0 \; \forall i, j$ 

- polarization vectors and q and k span  $\mathbb{R}^{1,D-1}$
- needed: similar basis of spinors
- would like: vectors in terms of spinors

guess:  $e^{\mu}(\vec{h}) \sim \overline{\xi} \gamma^{\mu} \psi$ 

want:

$$q_\mu e^\mu = k_\mu e^\mu = \ldots = 0$$

therefore:

$$\overline{\xi}\left( \pmb{q}_{\mu}\gamma^{\mu}
ight) =\left( \pmb{k}_{\mu}\gamma^{\mu}
ight)\psi =0 \quad ext{or} \quad \overline{\xi}\left( \pmb{k}_{\mu}\gamma^{\mu}
ight) =\left( \pmb{q}_{\mu}\gamma^{\mu}
ight)\psi =0$$

other inner products  $\rightarrow$  total  $\frac{D}{2}$  annihilation conditions

• need spinors annihilated by  $\frac{D}{2}$  generators  $v_{\mu}^{i}\gamma^{\mu}$  such that  $\langle v^{i}, v^{j} \rangle = 0 \ \forall i, j \rightarrow \text{definition of pure spinors}$ 

• spinors transform under Clifford algebra  $\{\gamma^{\mu}, \gamma^{\nu}\} = 2\eta^{\mu\nu}$ 

- spinors transform under Clifford algebra  $\{\gamma^{\mu}, \gamma^{\nu}\} = 2\eta^{\mu\nu}$
- o define

$$\begin{split} \gamma_0^+ &\equiv \frac{1}{\sqrt{2}} k_\mu \gamma^\mu \qquad \gamma_0^- &\equiv \frac{1}{\sqrt{2}} \frac{q_\mu}{q \cdot k} \gamma^\mu \\ \gamma_i^+ &\equiv \frac{i}{\sqrt{2}} e_\mu^{+,i} \gamma^\mu \qquad \gamma_i^- &\equiv \frac{i}{\sqrt{2}} e_\mu^{-,i} \gamma^\mu \end{split}$$

- spinors transform under Clifford algebra  $\{\gamma^{\mu}, \gamma^{\nu}\} = 2\eta^{\mu\nu}$
- define

$$\gamma_0^+ \equiv \frac{1}{\sqrt{2}} k_\mu \gamma^\mu \qquad \gamma_0^- \equiv \frac{1}{\sqrt{2}} \frac{q_\mu}{q \cdot k} \gamma^\mu$$
$$\gamma_i^+ \equiv \frac{i}{\sqrt{2}} e_\mu^{+,i} \gamma^\mu \qquad \gamma_i^- \equiv \frac{i}{\sqrt{2}} e_\mu^{-,i} \gamma^\mu$$

• Clifford algebra:  $\frac{D}{2}$  copies of fermionic harmonic oscillator

$$\{\gamma_i^a, \gamma_j^b\} = \delta_{ij}\delta_{a,-b}$$

- spinors transform under Clifford algebra  $\{\gamma^{\mu}, \gamma^{\nu}\} = 2\eta^{\mu\nu}$
- define

$$\gamma_0^+ \equiv \frac{1}{\sqrt{2}} k_\mu \gamma^\mu \qquad \gamma_0^- \equiv \frac{1}{\sqrt{2}} \frac{q_\mu}{q \cdot k} \gamma^\mu$$
$$\gamma_i^+ \equiv \frac{i}{\sqrt{2}} e_\mu^{+,i} \gamma^\mu \qquad \gamma_i^- \equiv \frac{i}{\sqrt{2}} e_\mu^{-,i} \gamma^\mu$$

• Clifford algebra:  $\frac{D}{2}$  copies of fermionic harmonic oscillator

$$\{\gamma_i^a, \gamma_j^b\} = \delta_{ij}\delta_{a,-b}$$

- quantum numbers  $R_q^j = \frac{1}{2}[\gamma_j^+, \gamma_j^-] = \left(\gamma_j^+ \gamma_j^- \frac{1}{2}\right)$
- quantum numbers  $\leftrightarrow$  annihilation conditions,

$$\begin{aligned} \gamma_i^{2h^i}\psi(\vec{h}) &= 0 & \text{no sum }, & \gamma_i^{2h^i}\xi(\vec{h}) &= 0 & \text{no sum} \\ \gamma_\mu k^\mu \psi(\vec{h}) &= 0 & & \gamma_\mu q^\mu \xi(\vec{h}) &= 0 . \end{aligned}$$

- spinors transform under Clifford algebra  $\{\gamma^{\mu}, \gamma^{\nu}\} = 2\eta^{\mu\nu}$
- define

$$\gamma_0^+ \equiv \frac{1}{\sqrt{2}} k_\mu \gamma^\mu \qquad \gamma_0^- \equiv \frac{1}{\sqrt{2}} \frac{q_\mu}{q \cdot k} \gamma^\mu$$
$$\gamma_i^+ \equiv \frac{i}{\sqrt{2}} e_\mu^{+,i} \gamma^\mu \qquad \gamma_i^- \equiv \frac{i}{\sqrt{2}} e_\mu^{-,i} \gamma^\mu$$

• Clifford algebra:  $\frac{D}{2}$  copies of fermionic harmonic oscillator

$$\{\gamma_i^{a}, \gamma_j^{b}\} = \delta_{ij}\delta_{a,-b}$$

- quantum numbers  $R_q^j = \frac{1}{2}[\gamma_j^+, \gamma_j^-] = \left(\gamma_j^+ \gamma_j^- \frac{1}{2}\right)$
- quantum numbers ↔ annihilation conditions,

$$\begin{aligned} \gamma_i^{2h^i}\psi(\vec{h}) &= 0 \quad \text{no sum} , \quad \gamma_i^{2h^i}\xi(\vec{h}) &= 0 \quad \text{no sum} \\ \gamma_\mu k^\mu \psi(\vec{h}) &= 0 \quad \gamma_\mu q^\mu \xi(\vec{h}) &= 0 . \end{aligned}$$

phases, schmases

▲□ ▶ ▲ □ ▶ ▲ □ ▶ □ □ ● ● ●

• there is a natural spinor inner product

$$\overline{\psi\left(h_{0},\vec{h}\right)}\psi'\left(h_{0}',\vec{h}'\right)\sim\ \delta\left(h_{0}+h_{0}'\right)\delta\left(\vec{h}-\vec{h}'\right) \text{ where } \overline{\psi'}\psi\equiv\psi'^{\dagger}\gamma_{0}\psi$$

• there is a natural spinor inner product

$$\overline{\psi\left(\mathbf{h}_{0},\vec{\mathbf{h}}
ight)}\psi^{\prime}\left(\mathbf{h}_{0}^{\prime},\vec{\mathbf{h}}^{\prime}
ight)\sim\,\delta\left(\mathbf{h}_{0}+\mathbf{h}_{0}^{\prime}
ight)\delta\left(\vec{\mathbf{h}}-\vec{\mathbf{h}}^{\prime}
ight)$$
 where  $\overline{\psi^{\prime}}\psi\equiv\psi^{\prime\dagger}\gamma_{0}\psi$ 

 for real momenta, fixing one spinor product fixes all by algebra, dependent on phase conventions

• there is a natural spinor inner product

$$\overline{\psi\left(h_{0},\vec{h}
ight)}\psi'\left(h_{0}',\vec{h}'
ight)\sim\,\delta\left(h_{0}+h_{0}'
ight)\delta\left(\vec{h}-\vec{h}'
ight)$$
 where  $\overline{\psi'}\psi\equiv\psi'^{\dagger}\gamma_{0}\psi$ 

 for real momenta, fixing one spinor product fixes all by algebra, dependent on phase conventions

#### spinor helicity

There is a phase convention for which

$$egin{aligned} & k_{\mu}\gamma^{\mu} &= \sum_{ec{h}}\psi_{ec{h}}\overline{\psi_{ec{h}}} & q_{\mu}\gamma^{\mu} &= \sum_{ec{h}}\xi_{ec{h}}\overline{\xi_{ec{h}}} \ & q_{\mu} &= rac{1}{2}\overline{\xi_{ec{h}}}\gamma_{\mu}\xi_{ec{h}} & k_{\mu} &= rac{1}{2}\overline{\psi_{ec{h}}}\gamma_{\mu}\psi_{ec{h}} \end{aligned}$$

for any  $\vec{h}$ . Similar formulae for polarization vectors.

• there is a natural spinor inner product

$$\overline{\psi\left(h_{0},\vec{h}
ight)}\psi'\left(h_{0}',\vec{h}'
ight)\sim\,\delta\left(h_{0}+h_{0}'
ight)\delta\left(\vec{h}-\vec{h}'
ight)$$
 where  $\overline{\psi'}\psi\equiv\psi'^{\dagger}\gamma_{0}\psi$ 

 for real momenta, fixing one spinor product fixes all by algebra, dependent on phase conventions

#### spinor helicity

There is a phase convention for which

$$egin{aligned} & k_{\mu}\gamma^{\mu} &= \sum_{ec{h}}\psi_{ec{h}}\overline{\psi_{ec{h}}} & q_{\mu}\gamma^{\mu} &= \sum_{ec{h}}\xi_{ec{h}}\overline{\xi_{ec{h}}} \ & q_{\mu} &= rac{1}{2}\overline{\xi_{ec{h}}}\gamma_{\mu}\xi_{ec{h}} & k_{\mu} &= rac{1}{2}\overline{\psi_{ec{h}}}\gamma_{\mu}\psi_{ec{h}} \end{aligned}$$

for any  $\vec{h}$ . Similar formulae for polarization vectors.

• representation is redundant (D = 4, 1), (D = 6, 2), (D = 10, 4)

- choose highest weight  $\xi(\frac{1}{2}, \dots, \frac{1}{2})$   $(q\xi = 0)$
- define (phases of) other states using (ordered) lowering operators, e.g.

$$\psi(-\frac{1}{2}) \equiv \gamma_0^- \gamma_1^- \xi(\frac{1}{2}) = -\gamma_1^- \gamma_0^- \xi(\frac{1}{2})$$

- choose highest weight  $\xi(\frac{1}{2}, \dots, \frac{1}{2})$   $(q\xi = 0)$
- define (phases of) other states using (ordered) lowering operators, e.g.

$$\psi(-\frac{1}{2}) \equiv \gamma_0^- \gamma_1^- \xi(\frac{1}{2}) = -\gamma_1^- \gamma_0^- \xi(\frac{1}{2})$$

- this fixes the action of all generators on all states
- obtained states form a basis of all spinors

- choose highest weight  $\xi(\frac{1}{2}, \dots, \frac{1}{2})$   $(q\xi = 0)$
- define (phases of) other states using (ordered) lowering operators, e.g.

$$\psi(-\frac{1}{2}) \equiv \gamma_0^- \gamma_1^- \xi(\frac{1}{2}) = -\gamma_1^- \gamma_0^- \xi(\frac{1}{2})$$

- this fixes the action of all generators on all states
- obtained states form a basis of all spinors

numerics & lightcone analysis...

- choose highest weight  $\xi(\frac{1}{2}, \dots, \frac{1}{2})$   $(q\xi = 0)$
- define (phases of) other states using (ordered) lowering operators, e.g.

$$\psi(-\frac{1}{2}) \equiv \gamma_0^- \gamma_1^- \xi(\frac{1}{2}) = -\gamma_1^- \gamma_0^- \xi(\frac{1}{2})$$

- this fixes the action of all generators on all states
- obtained states form a basis of all spinors

#### numerics & lightcone analysis...

- pick a frame in which q = (1, 0, ..., 0, 1),  $\tilde{n}_i = (0, ..., 1_i, 0, ...)$ .
- pick a gamma matrix representation.
- find eigenvectors of  $q_{\mu}\gamma^{\mu}\xi=0$
- other quantum numbers are then easy ← momentum independent
- all solutions to the massless Dirac equation by applying  $\gamma_\mu k^\mu$

Some remarks Checks

Extensions

Intriguing further structure

To do

Rutger Boels (NBIA)

Pure Spinor Helicity Methods

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □

Checks

- 4D spinor helicity
- proposal for 6D [Cheung, O' Donnal, 09]
- lightcone frame

Extensions

Intriguing further structure

To do

Checks

- 4D spinor helicity
- proposal for 6D [Cheung, O' Donnal, 09]
- lightcone frame

Extensions

- generalization to all spinor representations (e.g. gravitini)
- solutions to the massive Dirac equation, use

$$k-qrac{k^2}{2q\cdot k}=k^{\flat}$$

Intriguing further structure

To do

Checks

- 4*D* spinor helicity
- proposal for 6D [Cheung, O' Donnal, 09]
- lightcone frame

Extensions

- generalization to all spinor representations (e.g. gravitini)
- solutions to the massive Dirac equation, use

$$k-q\frac{k^2}{2q\cdot k}=k^\flat$$

Intriguing further structure

 $\bullet\,$  choice of complex structure?  $\rightarrow\,$  pure spinor spaces To do

Checks

- 4*D* spinor helicity
- proposal for 6D [Cheung, O' Donnal, 09]
- lightcone frame

Extensions

- generalization to all spinor representations (e.g. gravitini)
- solutions to the massive Dirac equation, use

$$k-q\frac{k^2}{2q\cdot k}=k^\flat$$

Intriguing further structure

 $\bullet\,$  choice of complex structure?  $\rightarrow\,$  pure spinor spaces To do

- study supersymmetry
- find more amplitudes (on-shell recursion!)

### Outline

#### Motivation

#### 2 Covariant representation theory of

- Poincaré algebra
- Spin algebra
- Susy algebra

#### 3 Outlook

 susy relates scattering amplitudes with different particles independent of coupling constants

- susy relates scattering amplitudes with different particles independent of coupling constants
- 4*D*: derived from
  - Lagrangian [Grisaru, Pendleton, Van Nieuwenhuizen, 76]
    - on-shell susy algebra [Grisaru, Pendleton, 77]
- further only 4*D* fundamental massive multiplet [Schwinn, Weinzierl, 06]: derivation 1

- susy relates scattering amplitudes with different particles independent of coupling constants
- 4*D*: derived from
  - Lagrangian [Grisaru, Pendleton, Van Nieuwenhuizen, 76]
  - on-shell susy algebra [Grisaru, Pendleton, 77]
- further only 4*D* fundamental massive multiplet [Schwinn, Weinzierl, 06]: derivation 1
- should not require off-shell information

- susy relates scattering amplitudes with different particles independent of coupling constants
- 4*D*: derived from
  - Lagrangian [Grisaru, Pendleton, Van Nieuwenhuizen, 76]
  - on-shell susy algebra [Grisaru, Pendleton, 77]
- further only 4D fundamental massive multiplet [Schwinn, Weinzierl, 06]: derivation 1
- should not require off-shell information

#### In any susy S-matrix theory

 $\mathbf{0} = \langle \mathbf{0} | \textit{S} \, \Phi_{\text{in}} \, \textit{Q} \, | \mathbf{0} \rangle = \langle \mathbf{0} | \textit{S} \, \textit{[} \textit{Q}, \Phi_{\text{in}} \textit{]} | \mathbf{0} \rangle = \langle \mathbf{0} | \textit{S} \, \textit{Q} \, | \text{in} \rangle$ 

need: knowledge of action of supersymmetry on free in-state

- susy relates scattering amplitudes with different particles independent of coupling constants
- 4*D*: derived from
  - Lagrangian [Grisaru, Pendleton, Van Nieuwenhuizen, 76]
  - on-shell susy algebra [Grisaru, Pendleton, 77]
- further only 4D fundamental massive multiplet [Schwinn, Weinzierl, 06]: derivation 1
- should not require off-shell information

#### In any susy S-matrix theory

 $0 = \langle 0 | \textit{S} \, \Phi_{\text{in}} \, \textit{Q} \, | 0 \rangle = \langle 0 | \textit{S} \, [\textit{Q}, \Phi_{\text{in}}] | 0 \rangle = \langle 0 | \textit{S} \, \textit{Q} \, | \text{in} \rangle$ 

- need: knowledge of action of supersymmetry on free in-state
- need: covariant supersymmetry representation theory

Rutger Boels (NBIA)

$$\{Q, \overline{Q}\} = 2k_{\mu}\gamma^{\mu}$$

$$\{ {\it Q}, \overline{\it Q} \} = 2 {\it k}_\mu \gamma^\mu$$

and expand everything into the pure spinor basis, e.g.

$$Q \equiv \sum_{\vec{h}} \left( Q_{\vec{h}} \psi(k, \vec{h}) + Q'_{\vec{h}} \xi(q, \vec{h}) \right)$$

$$\{ {\it Q}, \overline{\it Q} \} = 2 {\it k}_\mu \gamma^\mu$$

and expand everything into the pure spinor basis, e.g.

$$Q \equiv \sum_{\vec{h}} \left( Q_{\vec{h}} \psi(k, \vec{h}) + Q'_{\vec{h}} \xi(q, \vec{h}) \right)$$

from which we get  $\rightarrow$ 

$$\{Q_{\vec{h}}, \overline{Q_{\vec{h}}}\} = 2 \qquad Q'_{\vec{h}} = 0$$

$$\{ {\it Q}, \overline{\it Q} \} = 2 {\it k}_\mu \gamma^\mu$$

and expand everything into the pure spinor basis, e.g.

$$Q \equiv \sum_{\vec{h}} \left( Q_{\vec{h}} \psi(k, \vec{h}) + Q'_{\vec{h}} \xi(q, \vec{h}) \right)$$

from which we get  $\rightarrow$ 

$$\{Q_{\vec{h}}, \overline{Q_{\vec{h}}}\} = 2 \qquad Q'_{\vec{h}} = 0$$

in terms of Lorentz invariant generators

$$|Q_{ec{h}}|k,ec{g}
angle = |k,ec{g}+ec{h}
angle$$

 $\rightarrow$  action of any Q known covariantly

Rutger Boels (NBIA)

(I) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1)) < ((1))

#### Example: $\mathcal{N} = 1, D = 10$

- bosonic states are e.g.  $\vec{h} = (\pm 1, 0, 0, 0) \ (\# = 8)$
- fermionic ones are  $\vec{h} = (\pm \frac{1}{2}, \pm \frac{1}{2}, \pm \frac{1}{2}, a)$ , last 'bit' fixed by chirality
- drop last bit

#### Example: $\mathcal{N} = 1, D = 10$

- bosonic states are e.g.  $\vec{h} = (\pm 1, 0, 0, 0) \ (\# = 8)$
- fermionic ones are  $\vec{h} = (\pm \frac{1}{2}, \pm \frac{1}{2}, \pm \frac{1}{2}, a)$ , last 'bit' fixed by chirality
- drop last bit



Rutger Boels (NBIA)

금 등 김 금 등 - 금 등
• every fermionic harmonic oscillator has a coherent state representation,

$$|k\eta_i
angle=m{e}^{\left(m{Q}^i\eta_i
ight)}|k,ec{h}_{ ext{top}}
angle$$

• every fermionic harmonic oscillator has a coherent state representation,

$$|k\eta_i
angle=oldsymbol{e}^{ig(Q^i\eta_iig)}|k,ec{h}_{ ext{top}}
angle$$

• choice of highest weight state component  $\rightarrow$  split algebra into creation and annihilation operators

 every fermionic harmonic oscillator has a coherent state representation,

$$|k\eta_i\rangle = e^{\left(Q^i\eta_i\right)}|k, \vec{h}_{\mathrm{top}}\rangle$$

- choice of highest weight state component  $\rightarrow$  split algebra into creation and annihilation operators
- Susy now acts naturally on the coherent state

$$\begin{split} \boldsymbol{e}^{\left(\boldsymbol{Q}_{c}^{i}\boldsymbol{\theta}_{i}\right)}|\boldsymbol{k},\eta_{i}\rangle &= \boldsymbol{e}^{\sum_{i}\eta_{i}\boldsymbol{\theta}_{i}}|\boldsymbol{k},\eta_{i}+\eta_{i}'\rangle\\ \boldsymbol{e}^{\left(\boldsymbol{Q}_{a}^{i}\boldsymbol{\theta}_{i}\right)}|\boldsymbol{k},\eta_{i}\rangle &= |\boldsymbol{k},\eta_{i}+\theta_{i}\rangle \end{split}$$

<<p>(日本)

• every fermionic harmonic oscillator has a coherent state representation,

$$|k\eta_i\rangle = e^{\left(Q^i\eta_i\right)}|k, \vec{h}_{
m top}
ight
angle$$

- choice of highest weight state component  $\rightarrow$  split algebra into creation and annihilation operators
- Susy now acts naturally on the coherent state

$$m{e}^{m{\left(Q_{c}^{i} heta_{i}
ight)}}|m{k},\eta_{i}
angle=m{e}^{\sum_{i}\eta_{i} heta_{i}}|m{k},\eta_{i}+\eta_{i}'
angle$$

$$\boldsymbol{e}^{\left(\boldsymbol{Q}_{a}^{i}\boldsymbol{\theta}_{i}\right)}|\boldsymbol{k},\eta_{i}\rangle=|\boldsymbol{k},\eta_{i}+\boldsymbol{\theta}_{i}\rangle$$

- compare [Arkani-Hamed,Cachazo,Kaplan, 08] for  $\mathcal{N}=(4,8), D=4$  (8 + 6 additional states)
- coherent state scattering amplitudes naturally supersymmetric

• every fermionic harmonic oscillator has a coherent state representation,

$$|k\eta_i\rangle = e^{\left(Q^i\eta_i\right)}|k, \vec{h}_{\mathrm{top}}\rangle$$

- choice of highest weight state component  $\rightarrow$  split algebra into creation and annihilation operators
- Susy now acts naturally on the coherent state

$$e^{\left(\mathcal{Q}_{c}^{i} heta_{i}
ight)}|k,\eta_{i}
angle=e^{\sum_{i}\eta_{i} heta_{i}}|k,\eta_{i}+\eta_{i}^{\prime}
angle$$

$$e^{\left(Q_{a}^{i}\theta_{i}\right)}|k,\eta_{i}\rangle = |k,\eta_{i}+\theta_{i}\rangle$$

- compare [Arkani-Hamed,Cachazo,Kaplan, 08] for  $\mathcal{N}=(4,8), D=4$  (8 + 6 additional states)
- coherent state scattering amplitudes naturally supersymmetric
- same conclusions?  $\rightarrow$  next slide

• consider *n* particle 'all plus' amplitude,

$$A(k_i, \vec{h}) = \int \left(\prod_{i=1}^n d\eta_i^4\right) A(k_i, \eta_i)$$

• consider *n* particle 'all plus' amplitude,

$$A(k_i, \vec{h}) = \int \left(\prod_{i=1}^n d\eta_i^4\right) A(k_i, \eta_i)$$

 can use susy to shift one η<sub>i</sub> coherent state parameter to zero, without phase factors

• consider *n* particle 'all plus' amplitude,

$$A(k_i, \vec{h}) = \int \left(\prod_{i=1}^n d\eta_i^4\right) A(k_i, \eta_i)$$

- can use susy to shift one η<sub>i</sub> coherent state parameter to zero, without phase factors
- amplitudes vanish to all orders in  $\hbar, \alpha', g_s, \ldots$

• consider *n* particle 'all plus' amplitude,

$$A(k_i, \vec{h}) = \int \left(\prod_{i=1}^n d\eta_i^4\right) A(k_i, \eta_i)$$

- can use susy to shift one η<sub>i</sub> coherent state parameter to zero, without phase factors
- amplitudes vanish to all orders in  $\hbar, \alpha', g_s, \dots$
- applies to all top states which share a coherent state parameter
- Yang-Mills, Einstein gravity, string theory ...

• consider *n* particle 'all plus' amplitude,

$$A(k_i, \vec{h}) = \int \left(\prod_{i=1}^n d\eta_i^4\right) A(k_i, \eta_i)$$

- can use susy to shift one η<sub>i</sub> coherent state parameter to zero, without phase factors
- amplitudes vanish to all orders in  $\hbar, \alpha', g_s, \dots$
- applies to all top states which share a coherent state parameter
- Yang-Mills, Einstein gravity, string theory ...

an important subtlety

• consider *n* particle 'all plus' amplitude,

$$A(k_i, \vec{h}) = \int \left(\prod_{i=1}^n d\eta_i^4\right) A(k_i, \eta_i)$$

- can use susy to shift one η<sub>i</sub> coherent state parameter to zero, without phase factors
- amplitudes vanish to all orders in  $\hbar, \alpha', g_s, \ldots$
- applies to all top states which share a coherent state parameter
- Yang-Mills, Einstein gravity, string theory ...

#### an important subtlety

does the susy transformation exist?

• consider *n* particle 'all plus' amplitude,

$$A(k_i, \vec{h}) = \int \left(\prod_{i=1}^n d\eta_i^4\right) A(k_i, \eta_i)$$

- can use susy to shift one η<sub>i</sub> coherent state parameter to zero, without phase factors
- amplitudes vanish to all orders in  $\hbar, \alpha', g_s, \ldots$
- applies to all top states which share a coherent state parameter
- Yang-Mills, Einstein gravity, string theory ....

#### an important subtlety

does the susy transformation exist?

- D = 4: free parameters ('one helicity unequal' = 0)
  - any chiral/antichiral susy transformation (allowed if  $U(1)_R$  is unbroken)

• consider *n* particle 'all plus' amplitude,

$$A(k_i, \vec{h}) = \int \left(\prod_{i=1}^n d\eta_i^4\right) A(k_i, \eta_i)$$

- can use susy to shift one  $\eta_i$  coherent state parameter to zero, without phase factors
- amplitudes vanish to all orders in  $\hbar, \alpha', g_s, \dots$
- applies to all top states which share a coherent state parameter
- Yang-Mills, Einstein gravity, string theory ...

#### an important subtlety

does the susy transformation exist?

- D = 4: free parameters ('one helicity unequal' = 0)
  - any chiral/antichiral susy transformation (allowed if  $U(1)_R$  is unbroken)

• D > 4: fixes susy transformation ('one helicity unequal'  $\neq 0$ )

crosscheck

extensions

to do

Ruto	ier E	Boels	(NB	IA)

crosscheck

•  $\exists$  field theory derivation for  $\mathcal{N} = 1$  D = 10

extensions

to do

crosscheck

•  $\exists$  field theory derivation for  $\mathcal{N} = 1$  D = 10

extensions

- ∃ extension to massive reps (reproduces *D* = 4 fundamental massive)
- BPS representations by decomposition [Fayet, 78]

to do

<<p>(日本)

crosscheck

•  $\exists$  field theory derivation for  $\mathcal{N} = 1$  D = 10

extensions

- ∃ extension to massive reps (reproduces *D* = 4 fundamental massive)
- BPS representations by decomposition [Fayet, 78]

to do

• generic central charges ('technicality')

# Conclusions and outlook

- complete spinor helicity construction from covariant representation theory
- quantum numbers under control
  - class of vanishing amplitudes

# Conclusions and outlook

- complete spinor helicity construction from covariant representation theory
- quantum numbers under control
  - class of vanishing amplitudes

- join the fun!
  - ▶ non-zero amplitudes? → on-shell recursion (work in progress)
  - less nuts and bolts?
  - connection to Berkovits' pure spinors?
  - spontaneous symmetry breaking?

Example:  $\mathcal{N} = 1, D = 11$ 



Rutger Boels (NBIA)

Pure Spinor Helicity Methods