

Dipole CFT and Integrability

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Non-relativistic holography

Schrödinger background:

$$ds^2 = -\frac{\ell^2 du^2}{z^4} + \frac{-2du dv + d\mathbf{x}^2 + dz^2}{z^2}$$

Son'08

Balasubramanian,McGreevy'08

Scale-invariant but time and space scale differently:

$$u \rightarrow \lambda^2 u$$

$$z \rightarrow \lambda z$$

$$\mathbf{x} \rightarrow \lambda \mathbf{x}$$

$$\ell = \frac{\sqrt{\lambda}}{2\pi} L$$

AdS length scale

CFT length scale

Holographic dictionary: $m_\phi \longleftrightarrow \Delta$

$$\phi(v, u, \mathbf{x}, z) \sim e^{iMv} \phi(u, \mathbf{x}, z)$$

Solving Klein-Gordon equation:

$$\begin{aligned}\Delta &= \frac{d}{2} + \sqrt{\frac{d^2}{4} + m_\phi^2 + \ell^2 M^2} \\ &= \frac{d}{2} + \sqrt{\frac{d^2}{4} + m_\phi^2 + \frac{\lambda}{4\pi^2} L^2 M^2}\end{aligned}$$

Volovich, Wen'08
Blau, Hartong, Rollier'09

not protected!

NR CFT & DLCQ

$$P_- |M\rangle = M |M\rangle \quad \Longrightarrow \quad P_+ = \frac{\mathbf{P}^2}{2M}$$

Unbroken symmetry (Schrödinger algebra):

$$\mathfrak{Sch} = \{T_a \in \mathfrak{so}(d, 2) | [T_a, P_-] = 0\}$$

$$\langle \underbrace{P_+, P_i, L_{ij}, L_{i-}}_{\text{Galilean}}; D - L_{+-}; K_-; M \rangle$$

Galilean

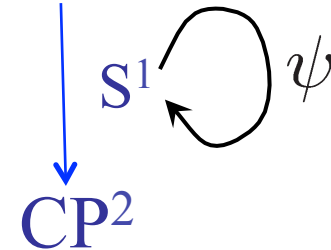
NR dilatation

Mass

Schrödinger from light-cone TsT

TsT transformation of $\text{AdS}_5 \times S^5$:

$$\begin{aligned} \text{T} - \text{duality} : \quad & \psi \rightarrow \tilde{\psi} \\ \text{s} \text{hift} : \quad & v \rightarrow v + \ell \tilde{\psi} \\ \text{T} - \text{duality} : \quad & \tilde{\psi} \rightarrow \tilde{\tilde{\psi}} \end{aligned}$$



$$ds^2 = -\frac{\ell^2 du^2}{z^4} + \frac{-2du dv + d\mathbf{x}^2 + dz^2}{z^2} + d\Omega_{S^5}^2$$

$$F_5 = (1 + *) \text{Vol}(S^5)$$

$$B = \frac{1}{z^2} du \wedge \eta$$

Herzog, Rangamani, Ross'08
Maldacena, Martelli, Tachikawa'08
Adams, Balasubramanian, McGreevy'08

vertical 1-form of Hopf fibration

TsT $\Leftrightarrow$ Twisted boundary conditions:

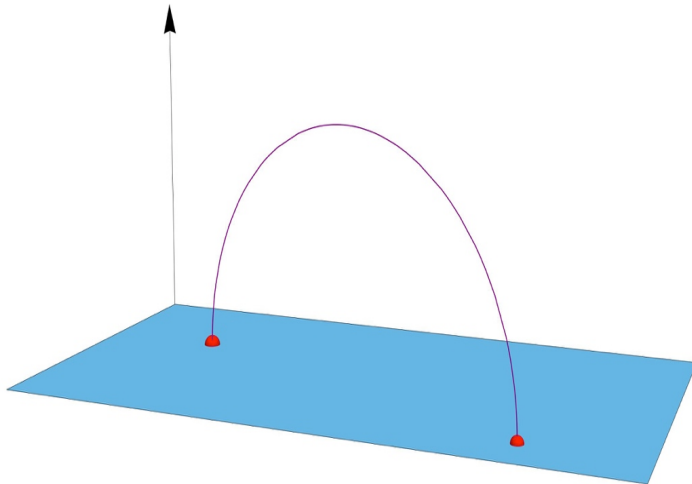
$$\begin{aligned} v(\sigma + 2\pi, \tau) &= v(\sigma, \tau) + LJ \\ \psi(\sigma + 2\pi, \tau) &= \psi(\sigma, \tau) - LM \end{aligned}$$

total ang. momentum

total light-cone momentum

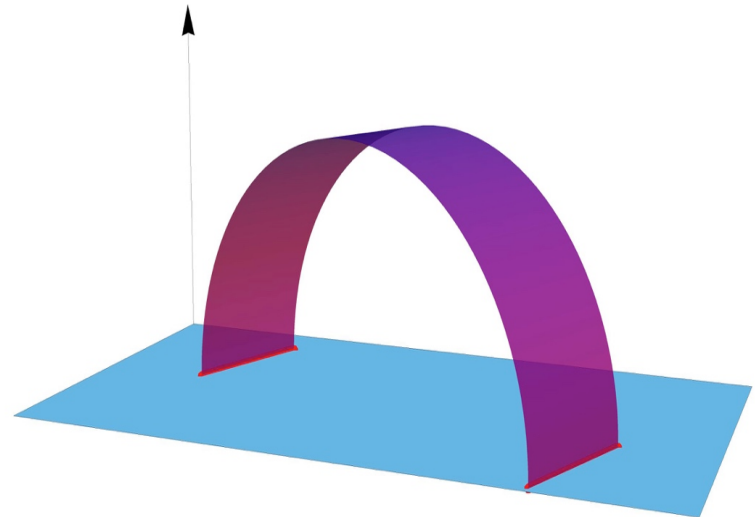
Example: BMN string

Closed-string b.c.'s:



$$\Delta = J$$

Twisted b.c.'s:



$$\Delta = \sqrt{J^2 + \frac{\lambda}{4\pi^2} L^2 M^2}$$

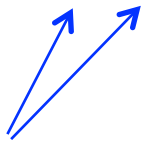
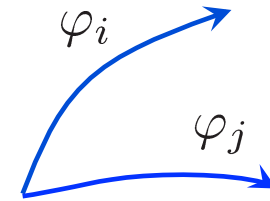
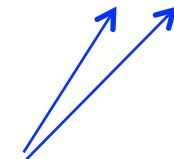
- dual op. is non-local
- has anomalous dim.

AdS/CFT triality

Lunin,Maldacena'05

Beisert,Roiban'05

Frolov'05

N=4 SYM	AdS ₅ ×S ⁵	Spin chain
Star Product	TsT	Drinfeld-Reshetikhin twist
$A * B = e^{\frac{i}{2} \gamma_{ij} Q_A^i Q_B^j} AB$  commuting charges	 commuting isometries	$h_{ab} \rightarrow F_{ab} h_{ab} F_{ba}$ $F_{ab} = e^{\frac{i}{2} \gamma_{ij} Q_a^i Q_b^j}$  Cartan generators

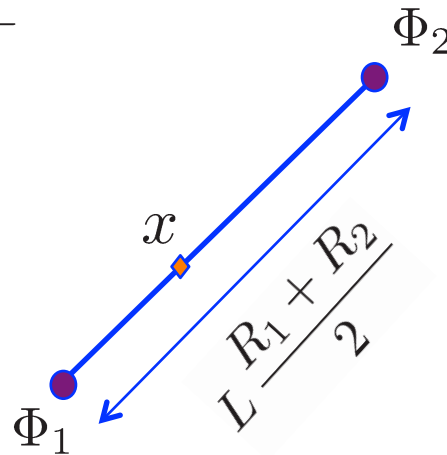
Dipole CFT

Bergman, Ganor '00

Star product:

$$\Phi_1 * \Phi_2(x) = e^{\frac{L}{2}(\partial_{-1} R_2 - R_1 \partial_{-2})} \Phi_1(x_1) \Phi_2(x_2) \Big|_{x_{1,2}=x} = \Phi_1(x + L_2) \Phi_2(x - L_1)$$

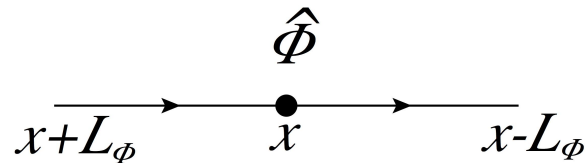
Dipole length: $L_a = \frac{L}{2} R_a n_-$



$$\mathcal{L} = \frac{2}{g^2} \text{tr} \left(-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \mathcal{D}_\mu Z_j^\dagger * \mathcal{D}^\mu Z_j - \frac{1}{2} [Z_j^\dagger, Z_j]_*^2 + [Z_j^\dagger, Z_k^\dagger]_* [Z_j, Z_k]_* \right)$$

Seiberg-Witten map

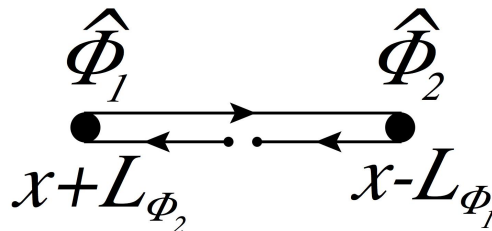
$$\Phi(x) \rightarrow U^{-1} * \Phi * U(x) = U^{-1}(x + L_\Phi) \Phi(x) U(x - L_\Phi)$$



$$\Phi(x) = [x + L_\Phi, x] \hat{\Phi}(x) [x, x - L_\Phi], \quad [x, y] = \text{P exp} \int_x^y dv A_-(v)$$

transforms locally

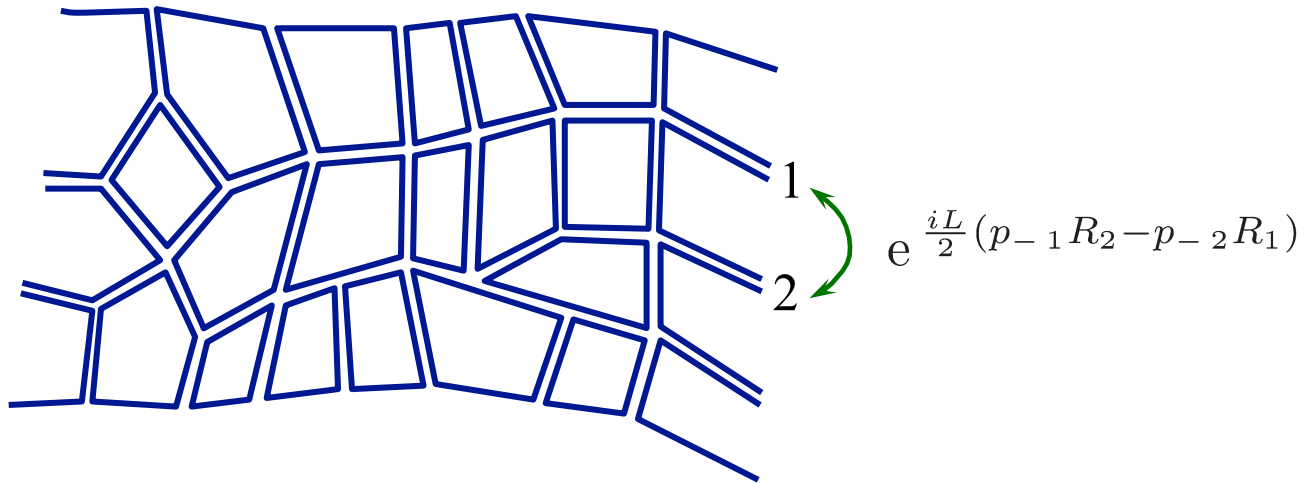
$$\hat{\Phi}_1 \hat{*} \hat{\Phi}_2 = e^{\frac{iL}{2}(P_{-1}R_2 - R_1P_{-2})} \hat{\Phi}_1 \hat{\Phi}_2, \quad P_\mu = -iD_\mu$$



Planar equivalence

$\Phi_1 \Phi_2 \rightarrow \Phi_1 * \Phi_2 \quad \longrightarrow \quad \text{Phase factors in interaction vertices}$

Equivalence theorem: phases cancel in planar diagrams
(or can be moved onto external lines)



Operators and spin chains

$$\mathcal{O} = \text{tr } \hat{\Phi}_1 \dots \hat{\Phi}_J$$

- Hamiltonian is twisted by a phase
- Boundary conditions are periodic

$$\mathcal{O} = \text{tr } \Phi_1 * \dots * \Phi_J$$

- Hamiltonian is the same as in N=4 SYM
- Boundary conditions are twisted

Drinfeld-Reshetikhin twist

Drinfeld'90
Reshetikhin'90

$$H = \sum_{l=1}^J h_{l,l+1}$$

$$\tilde{h}_{ab} = F_{ab} h_{ab} F_{ba}$$

Beisert,Roiban,05
Ahn,Bajnok,Bombardelli,Nepomechie'10

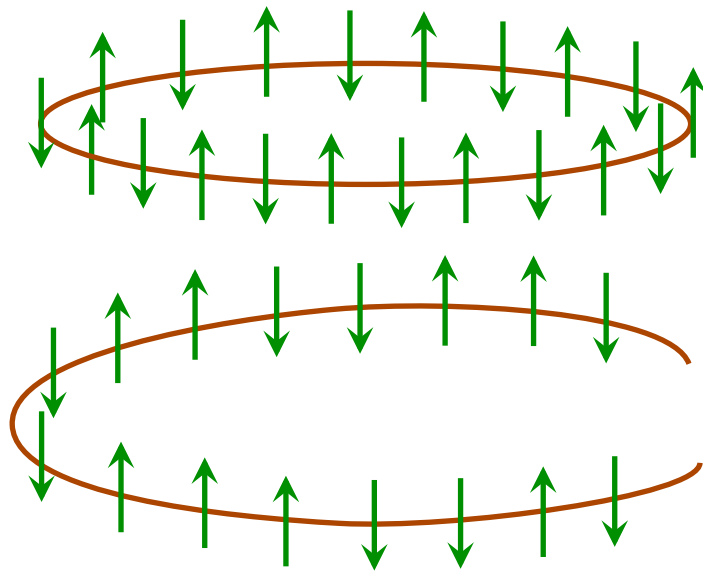
$$F_{ab} = e^{\frac{iL}{2}(P_{-a}R_b - R_aP_{-b})}$$

Model remains integrable:

$$\tilde{L}_{al}(u) = F_{al} L_{al}(u) F_{al}$$

satisfies YBE with the same R-matrix

Twist can be transferred from Hamiltonian to boundary conditions:



$$\mathcal{O} = \text{tr } \hat{\Phi}_1 \dots \hat{\Phi}_J$$

$$\tilde{h}_{l,l+1}$$

$$\mathcal{O} = \text{tr } \Phi_1 * \dots * \Phi_J$$

$$h_{l,l+1}$$

twist



closed
strings



twisted
B.C.'s

$$\begin{aligned}\varphi_j(\tau, \sigma + 2\pi) - \varphi_j(\tau, \sigma) &= -LM \\ v(\tau, \sigma + 2\pi) - v(\tau, \sigma) &= LR\end{aligned}$$

Frolov'05

Frolov,Roiban,Tseytlin'05

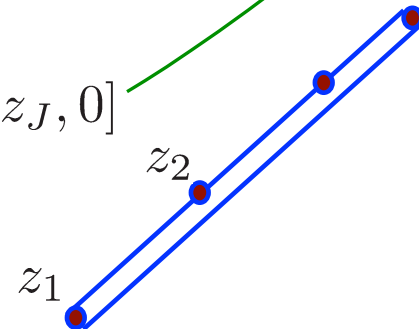
sl(2) sector

$$\mathcal{O} = \text{tr } D_-^{S_1} Z \dots D_-^{S_J} Z$$

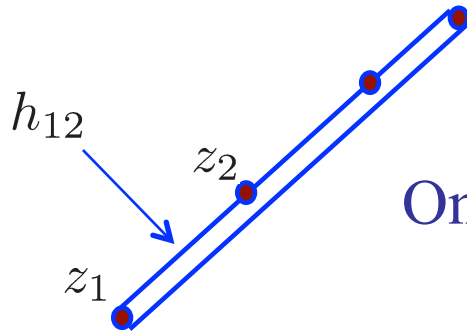
Taylor expansion

- Undeformed theory: finite-dim mixing of local ops
- Dipole CFT:
 - mixing: quantum-mechanical problem
 - eigenstates: non-local ops

$$\mathcal{O}(z_1, \dots, z_J) = \text{tr}[0, z_1] Z(z_1) [z_1, 0] \dots [0, z_J] Z(z_J) [z_J, 0]$$



Spin chain



One-loop dilatation operator:

$$D = \sum_{\ell=1}^L h_{\ell, \ell+1}$$

Untwisted:

$$h_{ab} \mathcal{O}(z_a, z_b) = \frac{\lambda}{8\pi^2} \int_0^1 \frac{du}{u} (2\mathcal{O}(z_a, z_b) - \mathcal{O}(z_a + uz_{ba}, z_b) - \mathcal{O}(z_a, z_b + uz_{ab}))$$

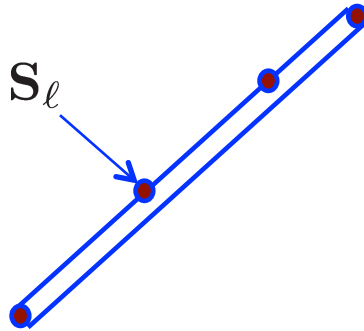
Balitsky, Braun'89

Belitsky, Derkachov, Korchemsky, Manashov'04

Twisted:

$$h_{ab} \mathcal{O}(z_a, z_b) = \frac{\lambda}{8\pi^2} \int_0^1 \frac{du}{u} (2\mathcal{O}(z_a, z_b) - \mathcal{O}(z_a + uz_{ba} + uL, z_b) - \mathcal{O}(z_a, z_b + uz_{ab} - uL))$$

Algebraic Bethe Ansatz



Spin $\frac{1}{2}$ representation of $\mathfrak{sl}(2)$ at each site:

$$S^0 = z\partial + \frac{1}{2}, \quad S^- = \partial, \quad S^+ = z^2\partial + z$$

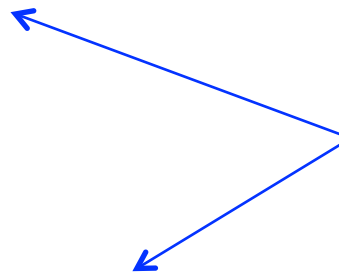
Lax operator (2x2 matrix):

$$L_\ell(u) = u + i\boldsymbol{\sigma} \cdot \mathbf{S}_\ell$$

Monodromy matrix:

$$T(u) = L_1(u) \dots L_J(u)$$

Satisfy YBE



$$T(u) = \begin{pmatrix} A(u) & B(u) \\ C(u) & D(u) \end{pmatrix}$$

$$|0\rangle = \delta(z_1) \dots \delta(z_J) \quad \Longleftrightarrow \quad \text{tr } Z^J(0)$$

$$C(u) |0\rangle = 0$$

Eigenstates: $B(u_1) \dots B(u_S) |0\rangle$

Bethe equations: $\left(\frac{u_j + \frac{i}{2}}{u_j - \frac{i}{2}} \right)^J = \prod_{k \neq j} \frac{u_j - u_k - i}{u_j - u_k + i}$

Anomalous dimension: $\Delta = J + S + \frac{\lambda}{8\pi^2} \sum_j \frac{1}{u_j^2 + \frac{1}{4}}$

Drinfeld-Reshetikhin twist

$$\tilde{T}(u) = L_1(u) \dots L_J(u) K$$

$$K = e^{L\sigma^-} = \begin{pmatrix} 1 & 0 \\ LJ & 1 \end{pmatrix}$$

Rep. of P_- in auxiliary space:

$$\tilde{T}(u) = \begin{pmatrix} A(u) + LJB(u) & B(u) \\ C(u) + LJD(u) & D(u) \end{pmatrix}$$

$$(C(u) + LJD(u)) |\Omega\rangle = 0 \quad \forall u$$

doesn't exist

No Bethe ansatz!

Baxter equation

$$\left(\frac{u_j + \frac{i}{2}}{u_j + \frac{i}{2}} \right)^J = - \prod_k \frac{u_j - u_k - i}{u_j - u_k + i}$$

Baxter polynomial: $Q(u) = \prod_j (u - u_j)$

$$(u + i/2)^J Q(u + i) + (u - i/2)^J Q(u - i) \equiv P(u)$$

Bethe eqns $\Rightarrow P(u_j) = 0 \Rightarrow P(u) = t(u)Q(u)$

Baxter equation:

$$(u + i/2)^J Q(u + i) + (u - i/2)^J Q(u - i) = t(u)Q(u)$$

$$t(u) = 2u^J + \dots$$

$$\Delta^{1-\text{loop}} = \frac{i\lambda}{8\pi^2} \partial_u \log \frac{Q(u + i/2)}{Q(u - i/2)} \Big|_{u=0}$$

Untwisted theory:

$$t(u) = 2u^J + 0u^{J-1} + \dots$$

Q(u) polynomial

Twisted theory:

$$t(u) = 2u^J + LMu^{J-1} + \dots$$

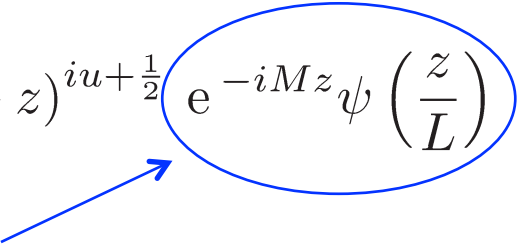
Q(u) non-polynomial

Example: twist-2 operators

$$\mathcal{O} = \text{tr } Z D^S Z$$

$$(u + i/2)^J Q(u + i) + (u - i/2)^J Q(u - i) = (2u^2 + LMu + t_0)Q(u)$$

Mellin transform: Faddeev, Korchemsky '93

$$Q(u) = \cosh \pi u \int_{-L}^L dz (L - z)^{-iu - \frac{1}{2}} (L + z)^{iu + \frac{1}{2}} e^{-iMz} \psi\left(\frac{z}{L}\right)$$


Coincides with coordinate-space wave function

$$(1 - x^2)\psi'' - 2x\psi' + (L^2 M^2 - x^2)\psi = (1/2 + t_0)\psi$$

Solution at M=0:

Untwisted

$$\psi(z) = \left(\frac{d}{dz} \right)^S \delta(z)$$

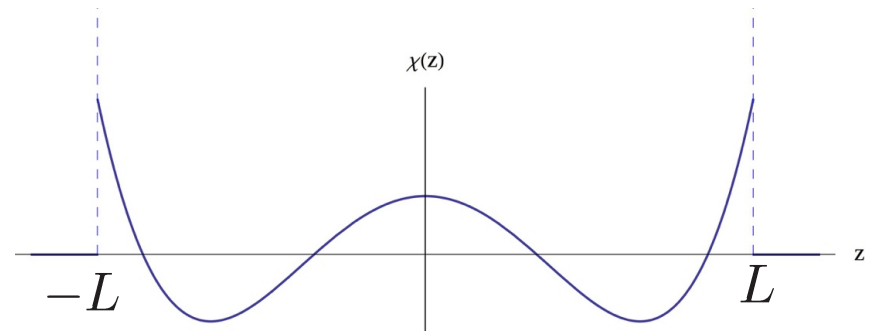
Twisted

$$\psi(z) = \sqrt{L} P_S(z/L) \times \text{Step}(z/L)$$

Legendre polynomials

Eigenstates – local ops. $\text{tr } Z D_-^S Z$

Eigenstates – non-local



Eigenvalues:

$$\delta\Delta = \frac{\lambda}{2\pi^2} h(S)$$

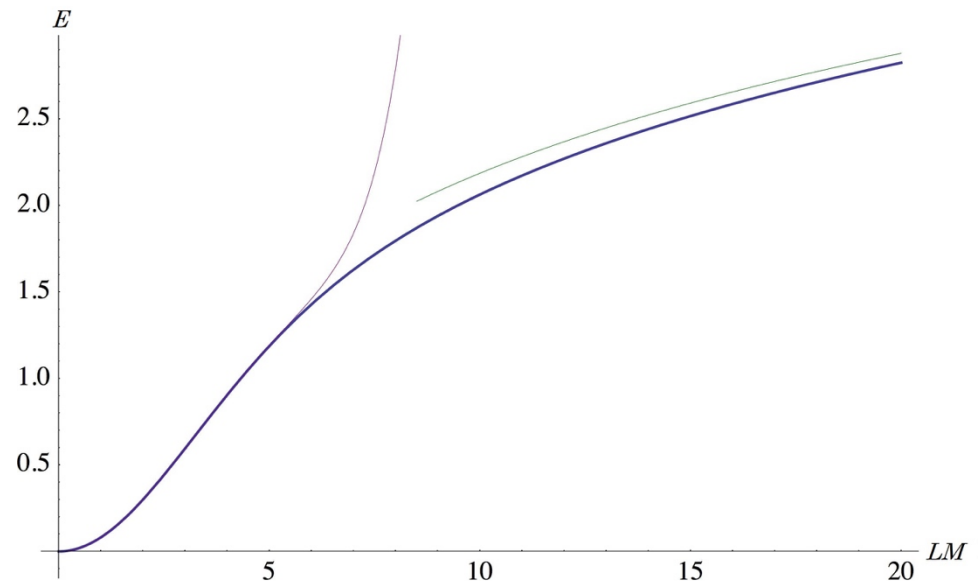
General solution

$$(1 - x^2)\psi'' - 2x\psi' + (L^2M^2 - x^2)\psi = (1/2 + t_0)\psi$$

Spheroidal wave equation

$$t_0 + 1/2 = -\lambda_{S,0} \left(\frac{LM}{2} \right)$$

$$\psi(x) = \text{PS}_{S,0} \left(\frac{LM}{2}, x \right)$$



Non-zero momentum

- States with different S mix
- Ground state energy is lifted

To leading order in LM:

$$\delta\Delta_0 = \frac{\lambda L^2 M^2}{24\pi^2} + \mathcal{O}(L^4 M^4)$$

More general result for any J:

$$\delta\Delta_0 = \frac{\lambda L^2 M^2}{8\pi^2(J+2)} + \mathcal{O}(L^4 M^4)$$

Agrees with BMN string energy at $J \gg 1$:

$$\Delta = \sqrt{J^2 + \frac{\lambda}{4\pi^2} L^2 M^2}$$

Conclusions

- Dipole CFT is integrable at planar level
- Full system of TQ (or QQ) equations at one loop?
- Quantum Spectral Curve:
 - ✓ same as in N=4 SYM, only asymptotics of Q-functions is different
- Dipole deformations of $\text{AdS}_3/\text{CFT}_2$

Azeyanagi, Hofman, Song, Strominger '12

Spinning BHs, Kerr/CFT,...

