

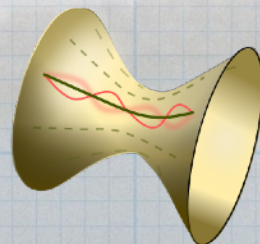
Wormholes in Quantum Ergodicity

Altland, Bagrets, PN, Sonner, Vielma
arXiv: 2105.12129

Belin, de Boer, PN, Sonner
arXiv: 2012.07875

Belin, de Boer, PN, Sonner
arXiv: 2111.xxxx

& some work in progress with J. Sonner and R. Baumgartner



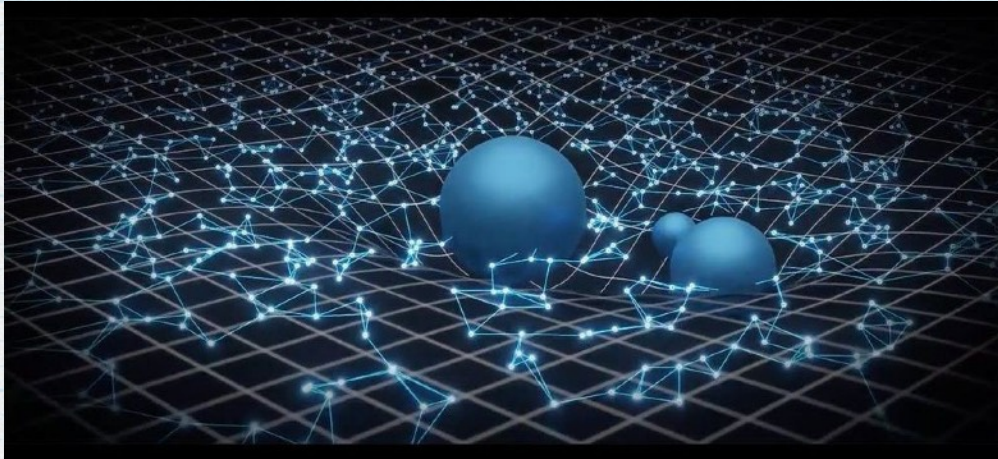
Ascona conference



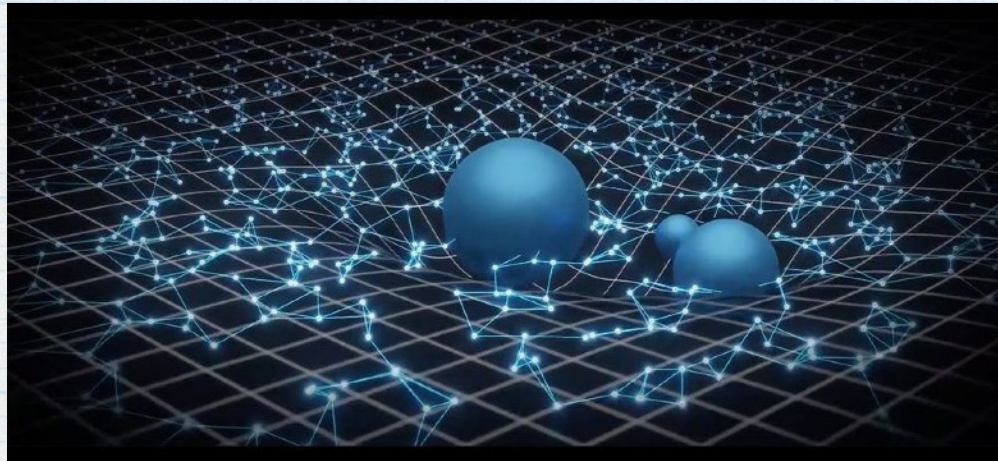
Good Old Dream of Holography



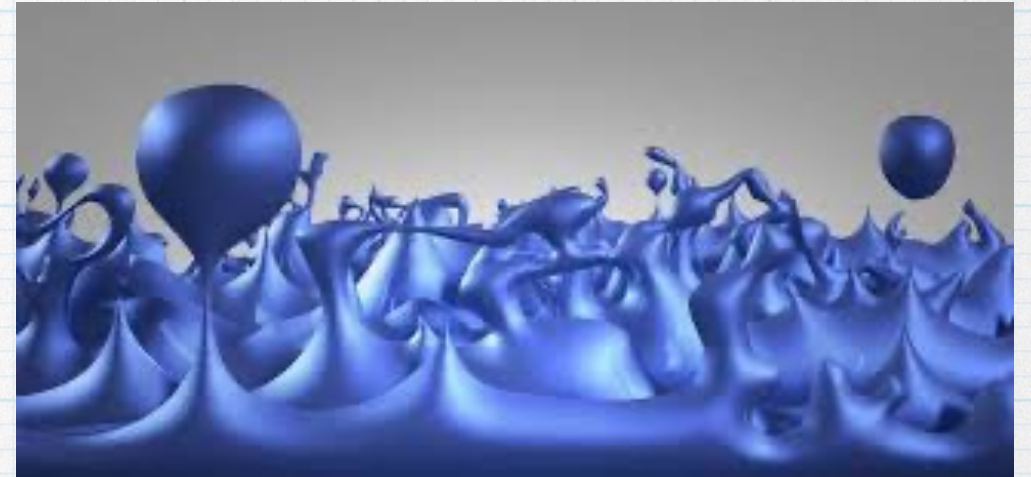
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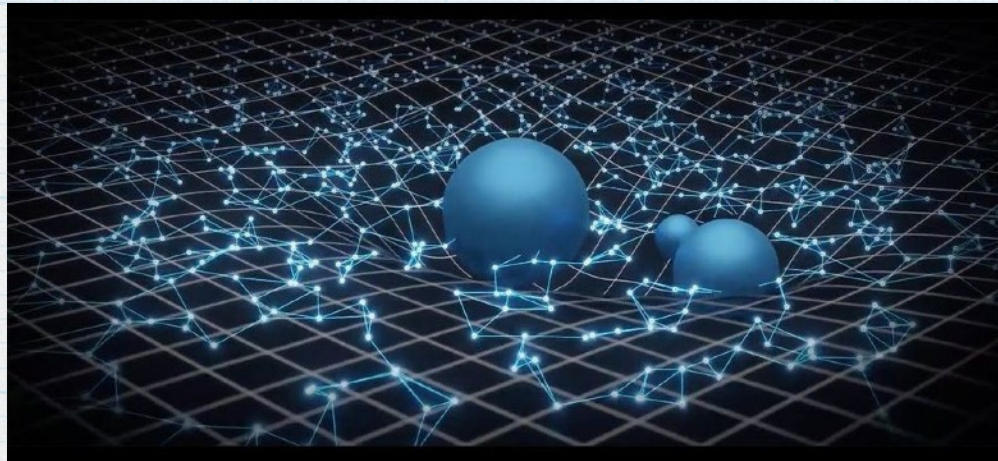
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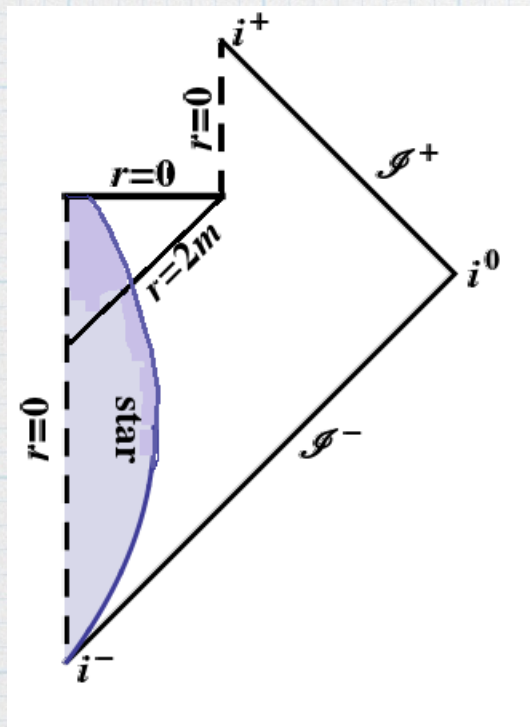
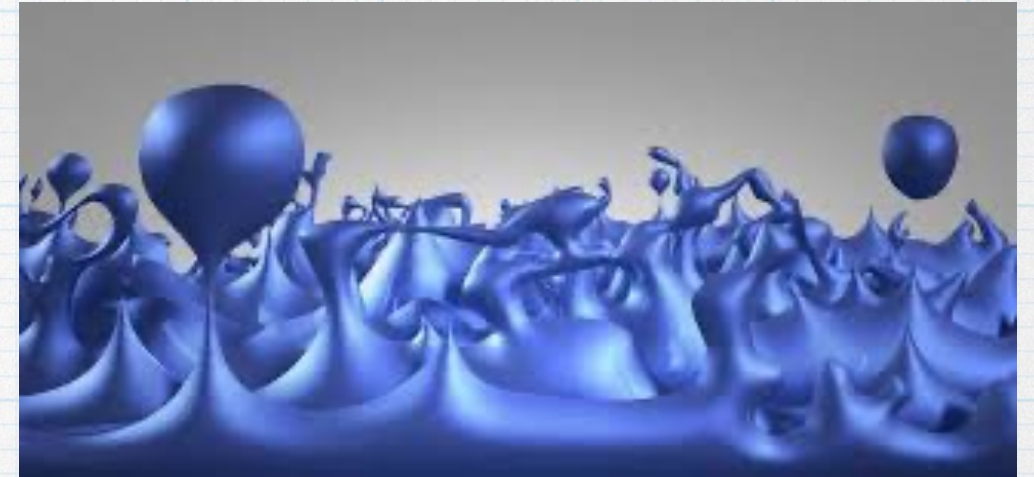
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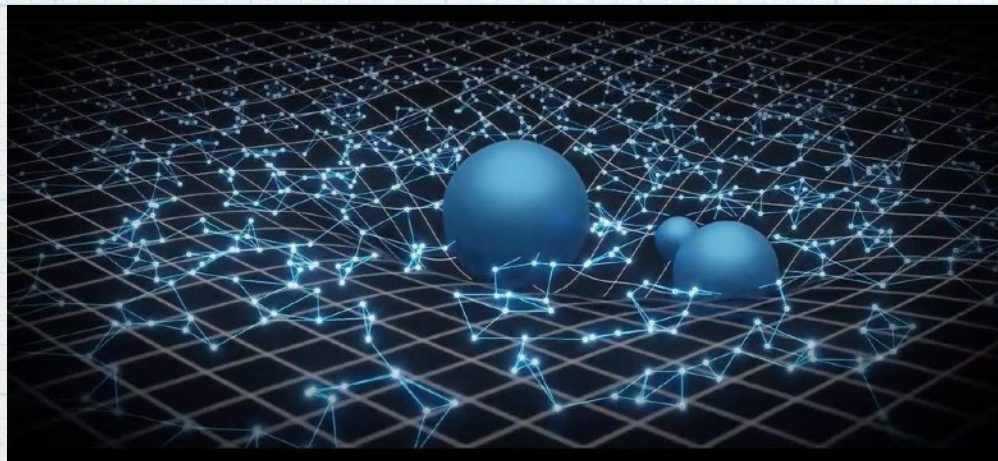


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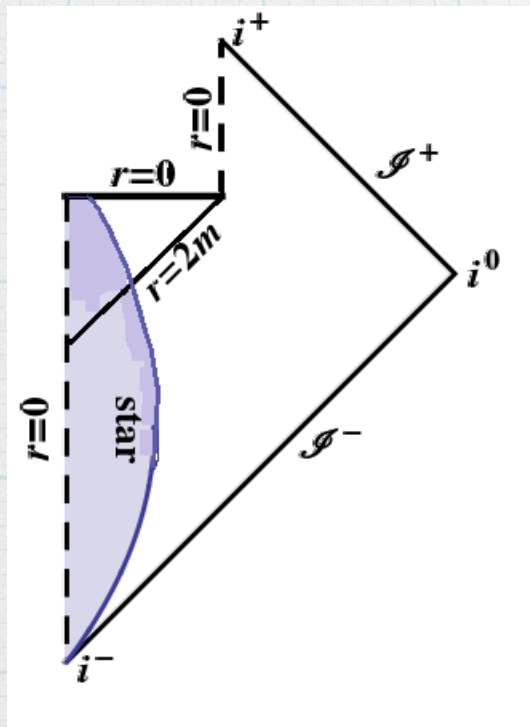
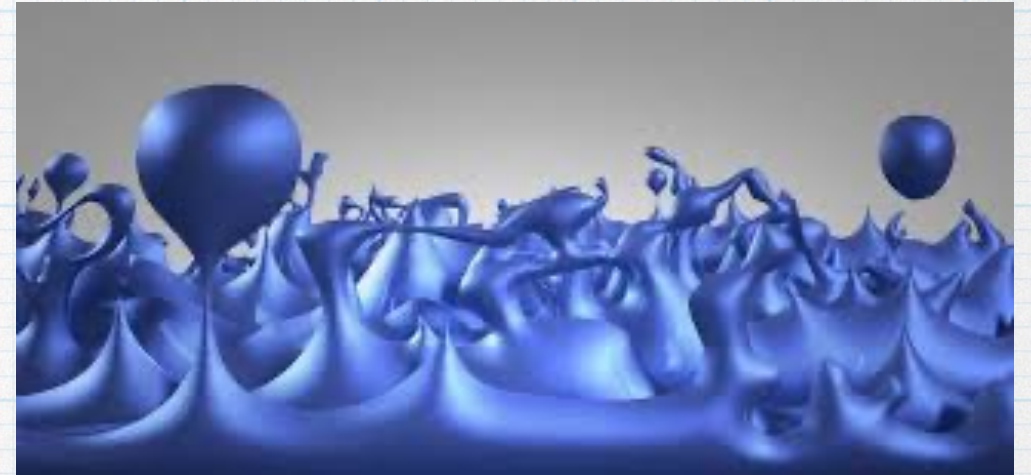


Information
Loss
Paradox

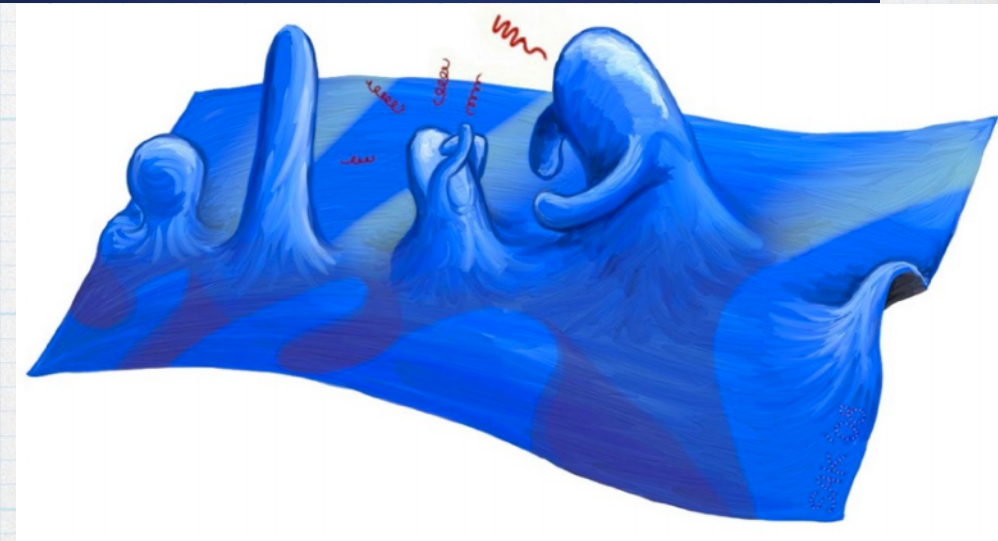
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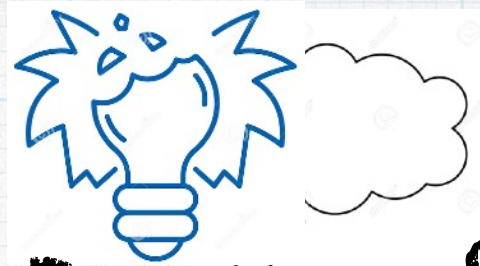
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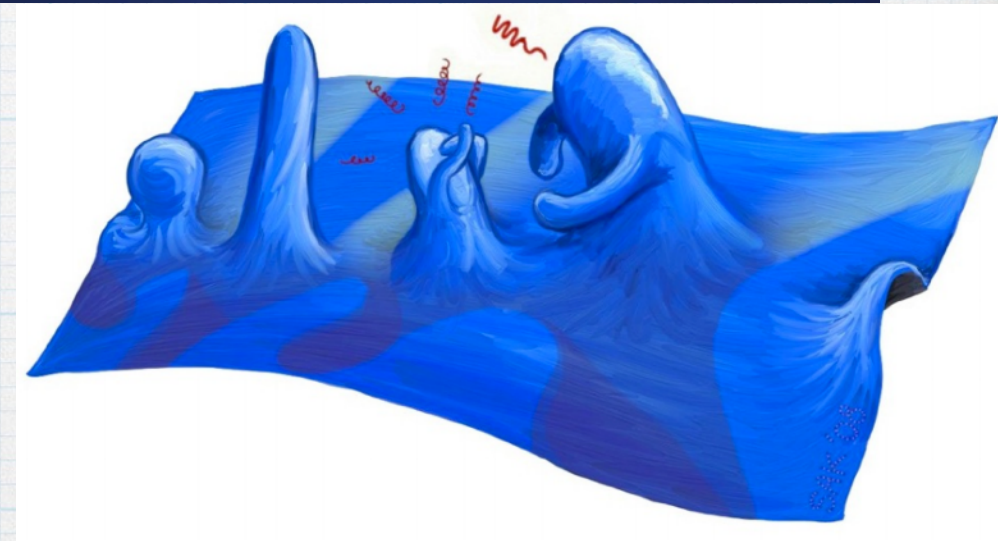
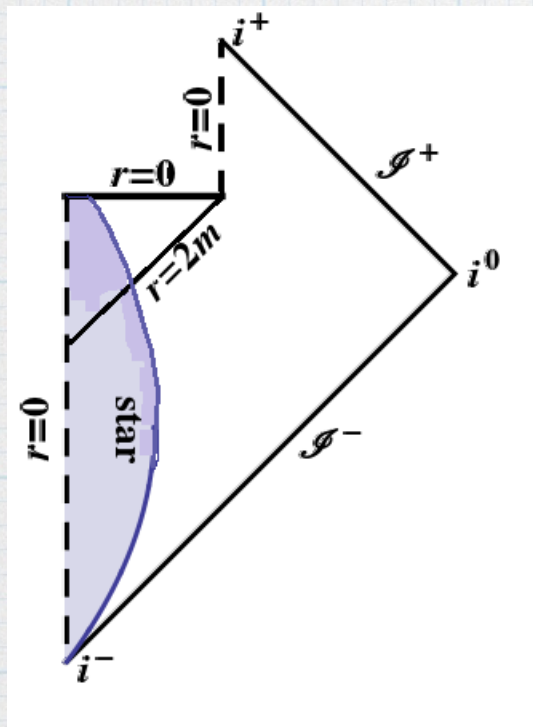
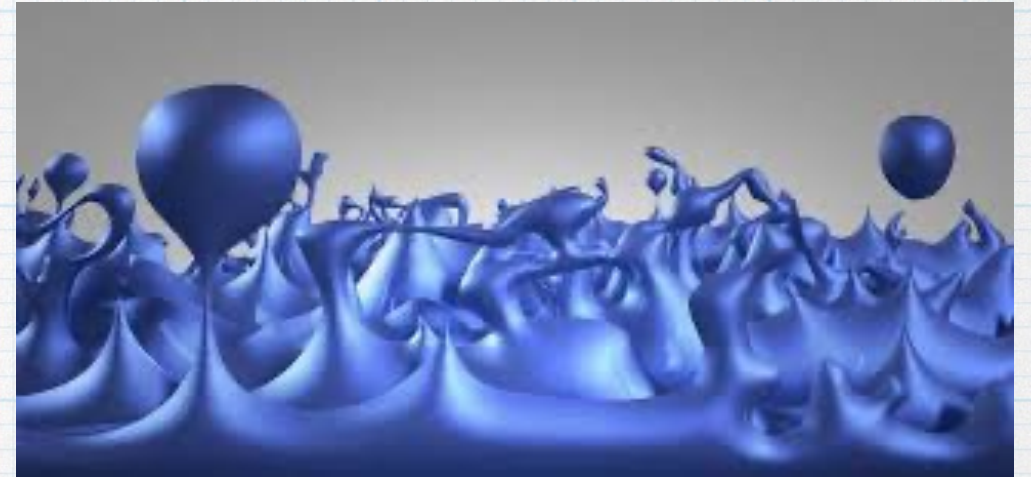
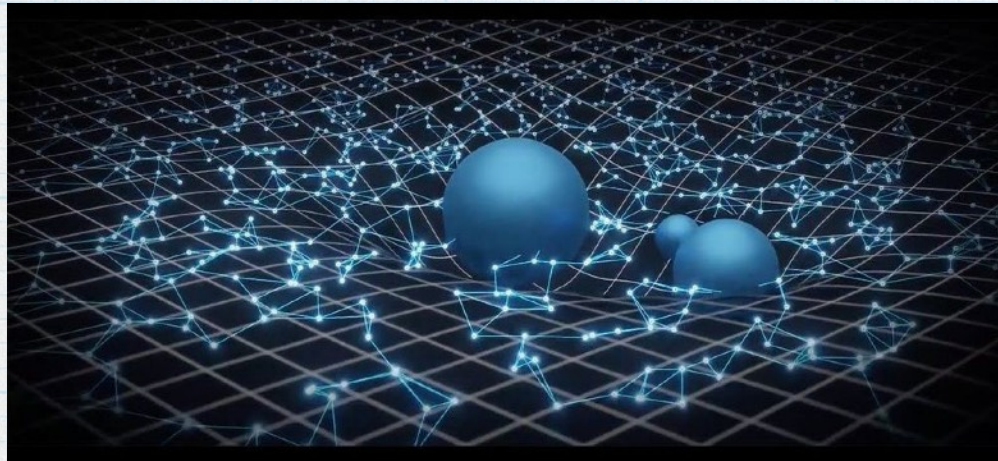
Information
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Strongly-coupled
QFT



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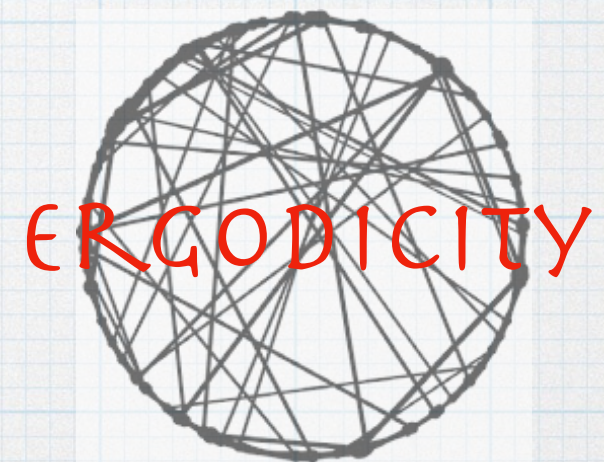


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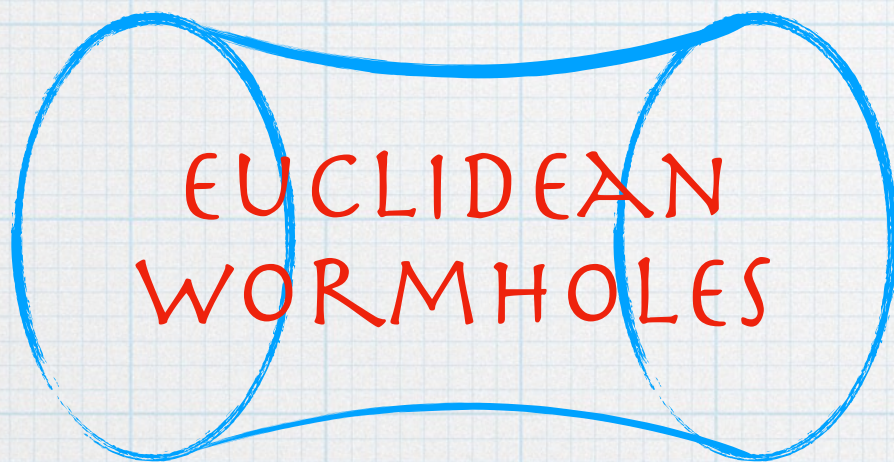
A Partial Resolution?!

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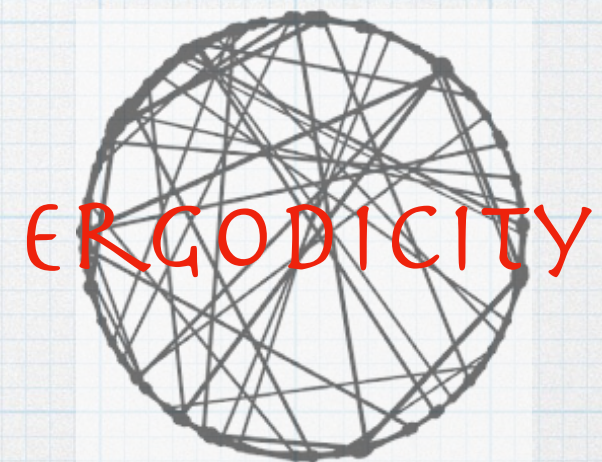


Altland, Bagrets '17;
Altland, Sonner '20;
Altland, Bagrets, *PN*, Sonner, Vielma '21;
Belin, de Boer, *PN*, Sonner '21;
more...

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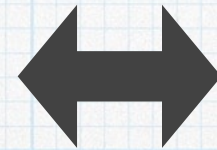
Saad, Shenker, Stanford '18 & '19;
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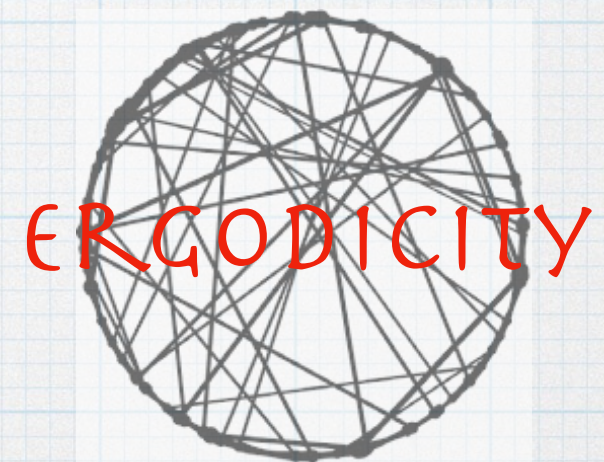
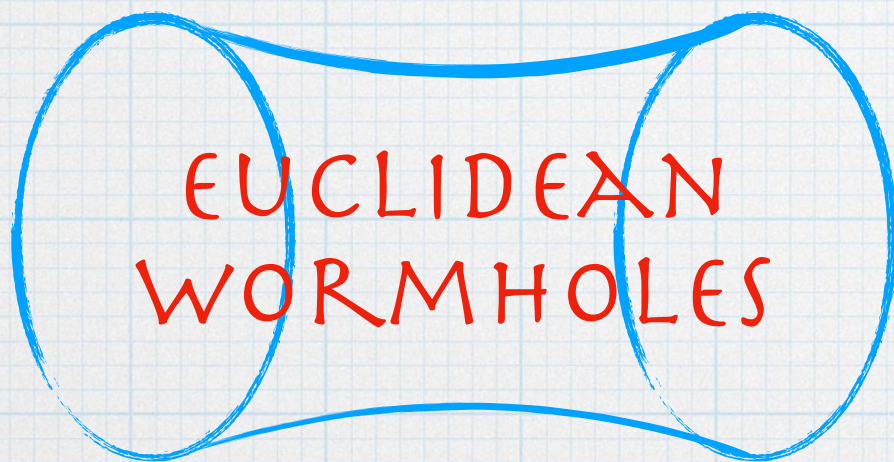
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A Partial Resolution?!

Semi-classical Gravity



Ergodic limit
of
Quantum theories



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ERGODICITY

The Connection

ETH

EUCLIDEAN
WORMHOLES

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Maldacena '02

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$$\xrightarrow{t \rightarrow \infty} \sum_i e^{-\beta E_i} |O_{ii}|^2 > 0$$

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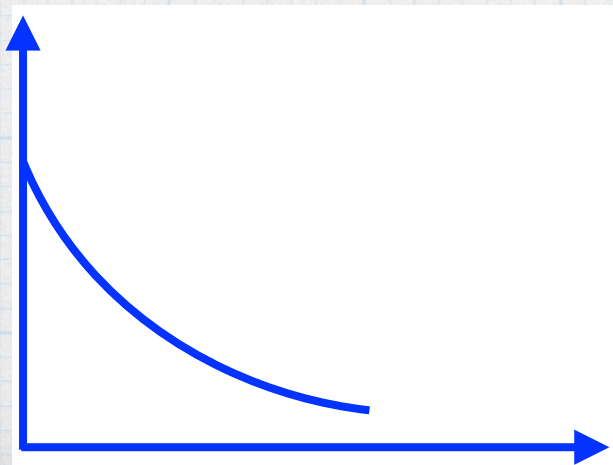
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- * The gravitational computation, however, gives a decaying answer



Gives a decaying
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Quantum Ergodicity

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Berry-Taylor '77;

Bohigas-Giannoni-Schmit '84

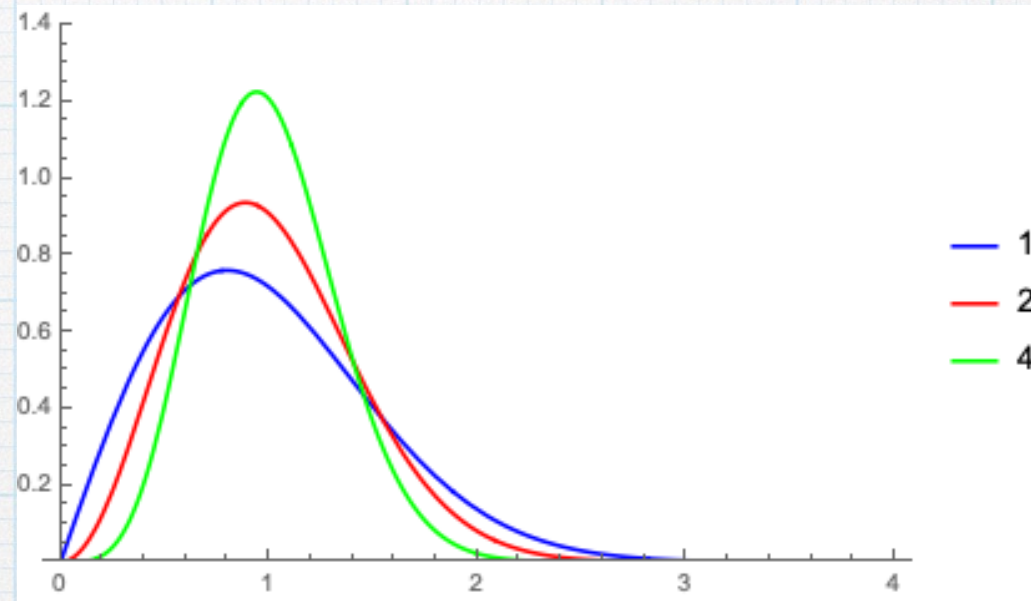
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$$P^n[\omega] = A_n \omega^n \exp[-B_n \omega^2]$$



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Quantum Ergodicity

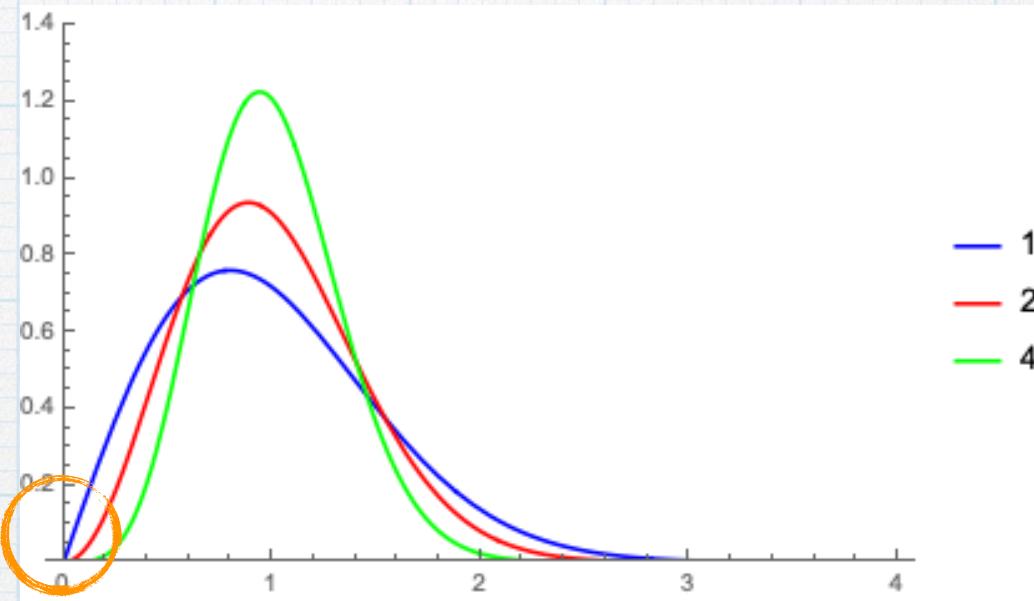
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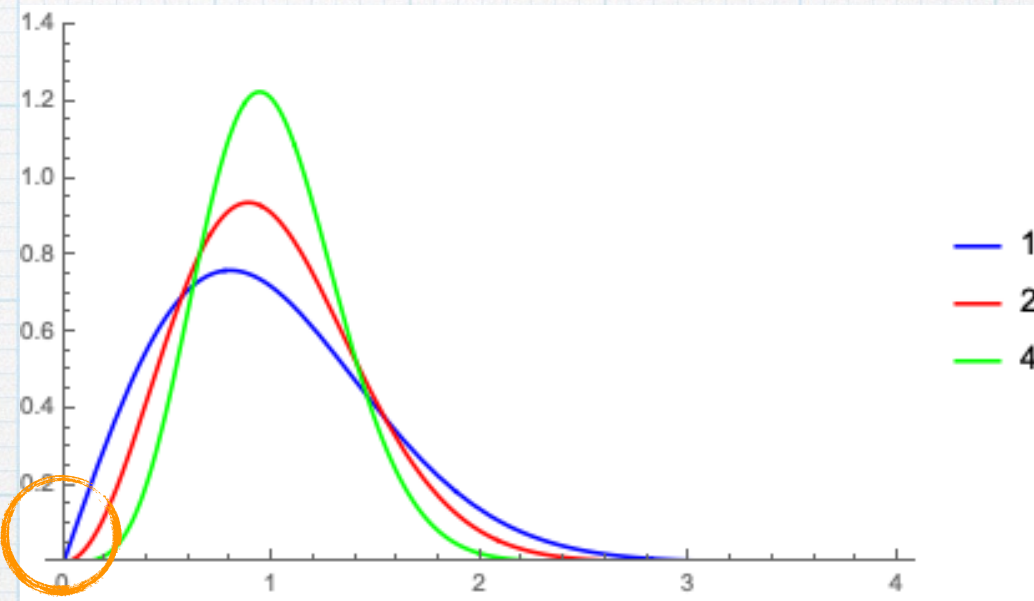
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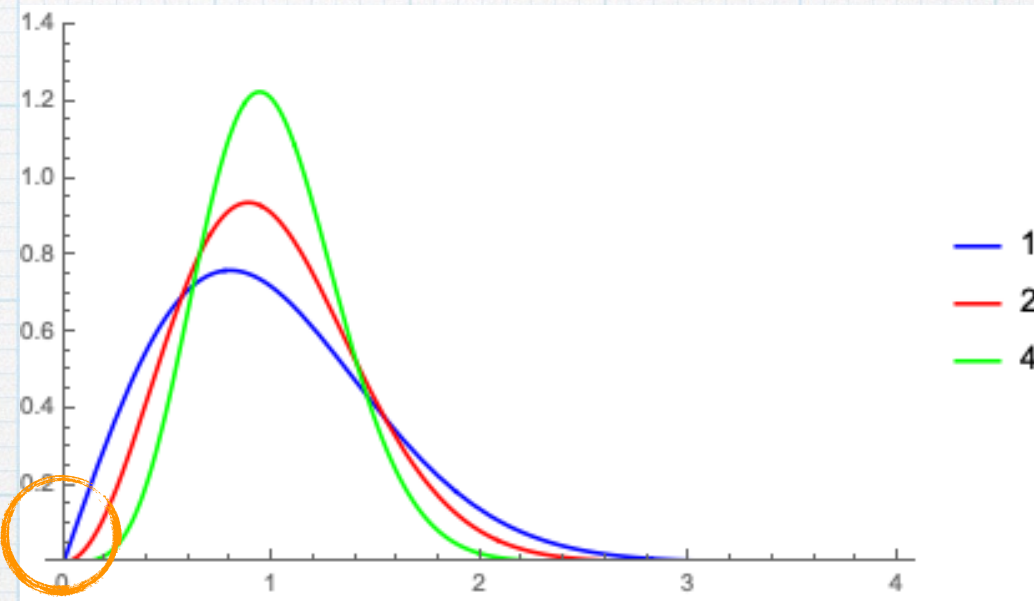
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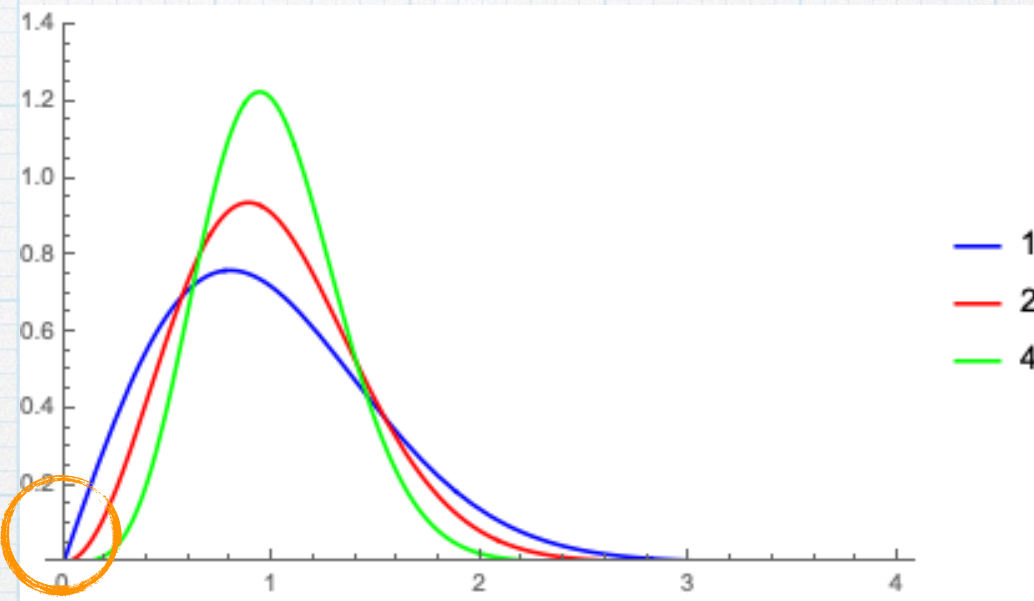
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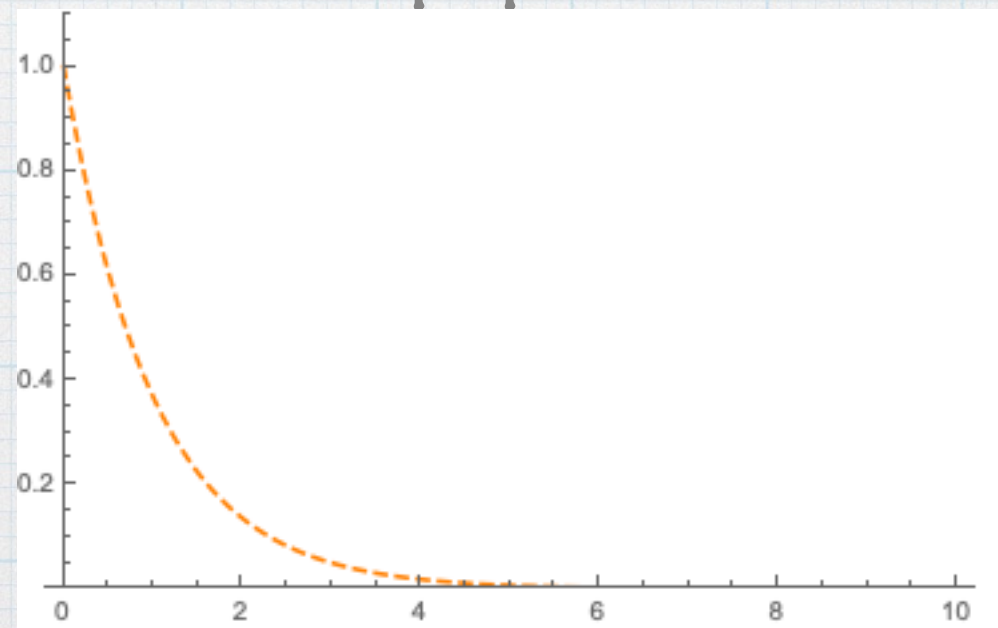
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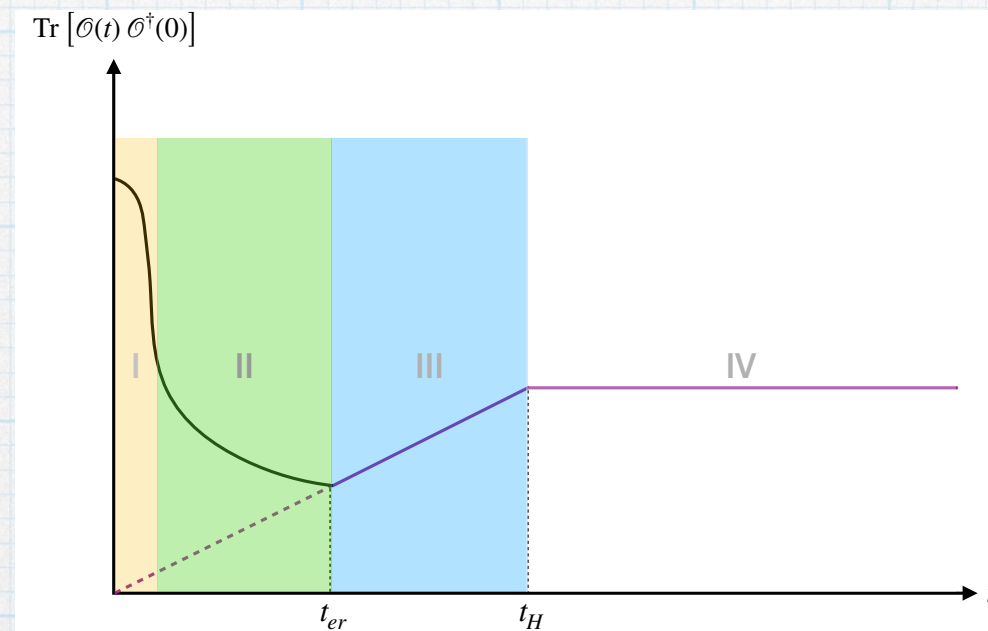
$$P^n[\omega] = \exp[-\omega]$$



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Quantum ergodicity... manifestations

- * Level repulsion and Spectral Rigidity in energy eigenstates of RMT leads to a characteristic **slope-ramp-plateau** behavior of observables like Spectral Form factor and correlation functions in RMT.



- * Ergodic limit is defined as the energy domain in which RMT statistics persists for a many-body (chaotic) system. Corresponding time is called **t_{er}**

Quantum ergodicity... a history

- * Quantum ergodicity has been studied numerically for a large class of models
- * Numerical computations of such observables in the SYK model demonstrate a similar behavior.
- * However, the **microscopic description** of this behavior in the SYK model (or any other system, for that matter) was **incomplete**.

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[Efetov '97 and references therein]
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[Cotler et. al. '16; Gharibyan, Hanada, Shenker, Tezuka '18
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Ergodic limit with σ -model

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- * Let us look at the resolvent of an operator which is defined as,

$$\begin{aligned}\tilde{R}(E, \omega) &= \sum_{\beta} |\langle \alpha | \mathcal{O} | \beta \rangle|^2 \delta(E_{\alpha} - E_{\beta} - \omega) \\ &= \int dE' \rho(E) \rho(E') |\langle E | \mathcal{O} | E' \rangle|^2 \delta(E - E' - \omega)\end{aligned}$$

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$$|\langle E | \mathcal{O} | E' \rangle|^2 = \mathcal{O}_i \mathcal{O}_j^* u_{E,i} u_{i,E'}^* u_{E',j} u_{j,E}^*$$

[Pollack, Rozali, Sully, Wakeham '20]

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The Resolvent

- * The resolvent is related to the Fourier transform of the thermal 2-point function,

$$\text{Tr}[e^{-\beta H} \mathcal{O}(t) \mathcal{O}^\dagger(0)] = \int dE d\omega e^{-\beta E + i\omega t} \tilde{R}(E, \omega)$$

ω is conjugate to the time, t

Late time $\longleftrightarrow$ small ω

- * Recall the identity,

$$\rho(E) = \pm \text{Im} \text{Tr} \left[\frac{1}{E \mp i\varepsilon - H} \right]_{\lim \varepsilon \rightarrow 0}$$

$$\rho(E_1)\rho(E_2) \sim -\text{Re} \left(\text{Tr} \left[\frac{1}{E + i\varepsilon - H} \right] \text{Tr} \left[\frac{1}{E - i\varepsilon - H} \right] \right)_{\lim \varepsilon \rightarrow 0}$$

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$$\frac{\text{Det}(z_2^+ - H) \text{Det}(z_4^- - H)}{\text{Det}(z_1^+ - H) \text{Det}(z_3^- - H)}$$

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Symmetry broken
b/w advanced &
retarded section

Also spontaneously
broken by the saddle
point in the limit
 $\dim(H) \gg 1$

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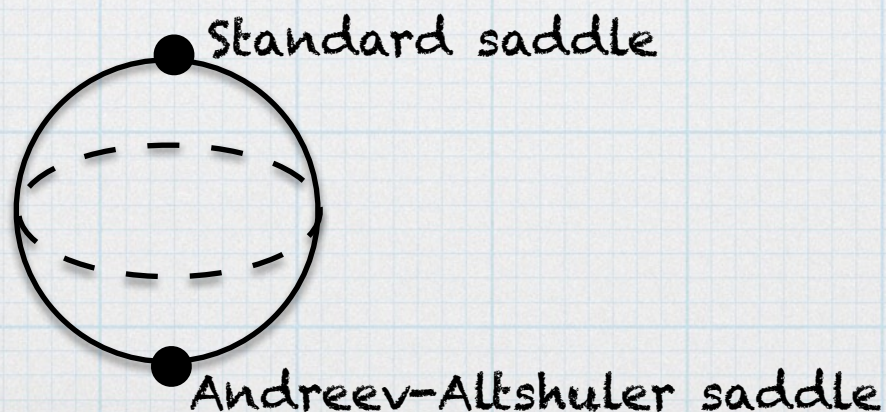
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- * There are two symmetry breaking saddle points in the limit, $\dim(H) \gg 1$:

$$\mathcal{M} = H_2 \times S^2$$



Operators in Ergodic limit

[Altland, Bagrets, *PN*, Sonner, Vielma '21]

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- * Operator correlation functions in the ergodic limit show following behaviour

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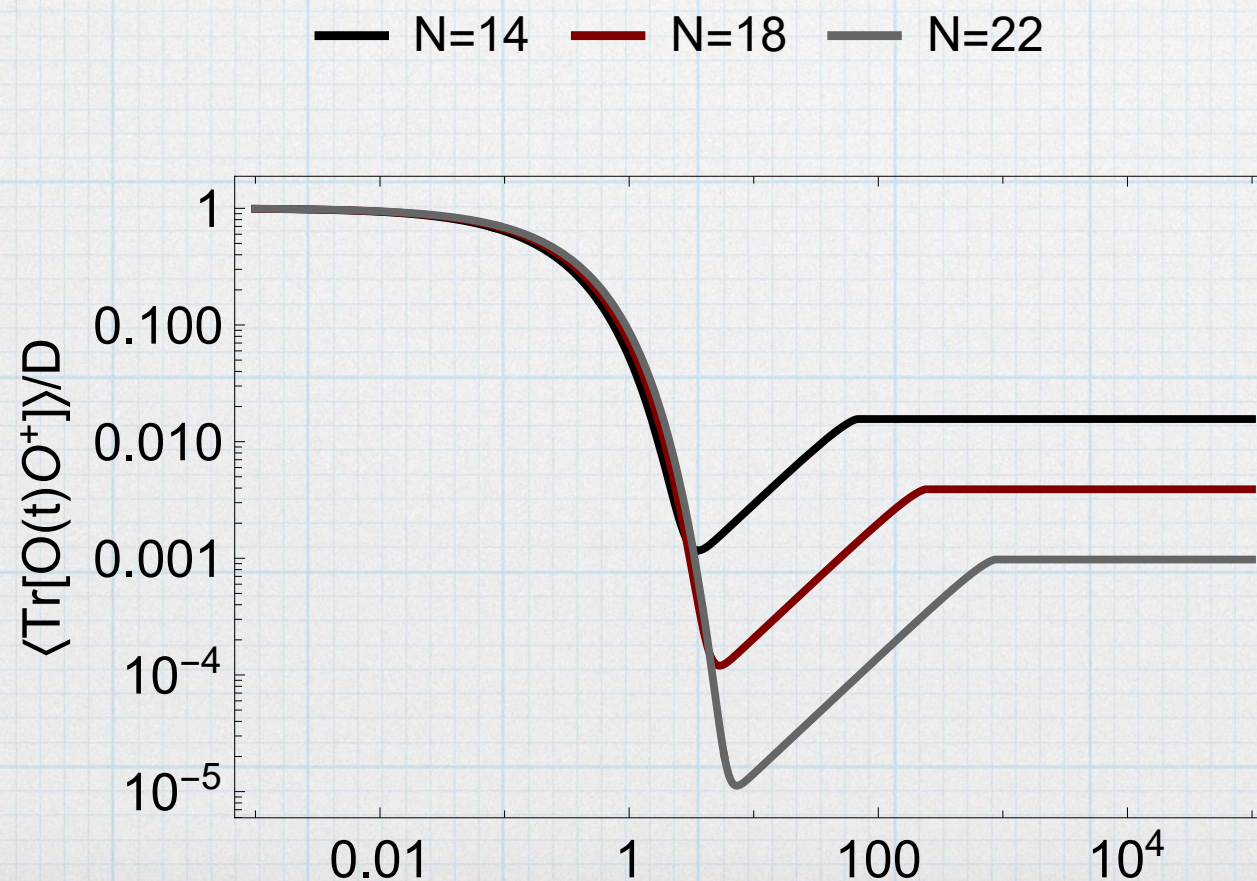
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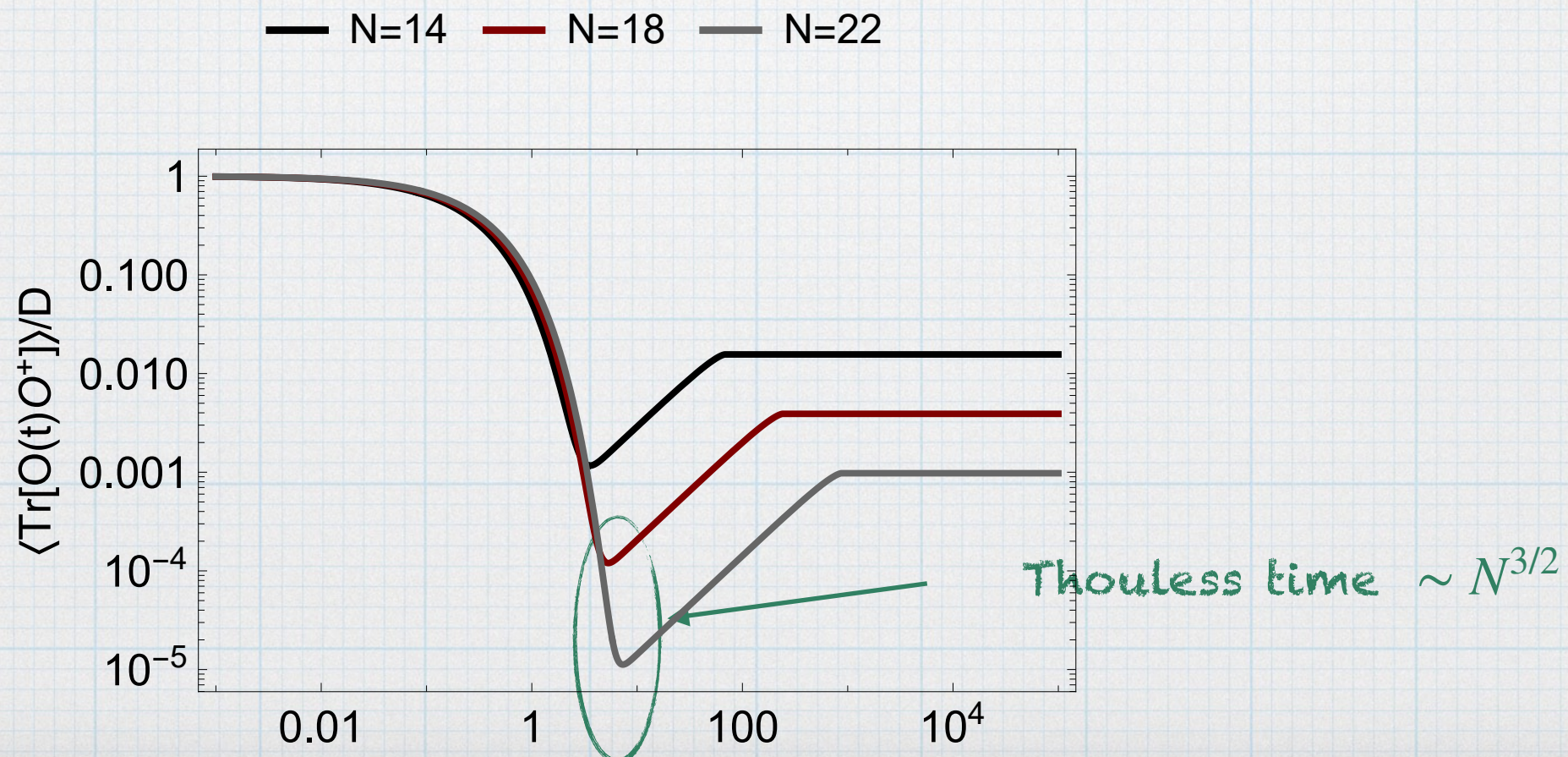
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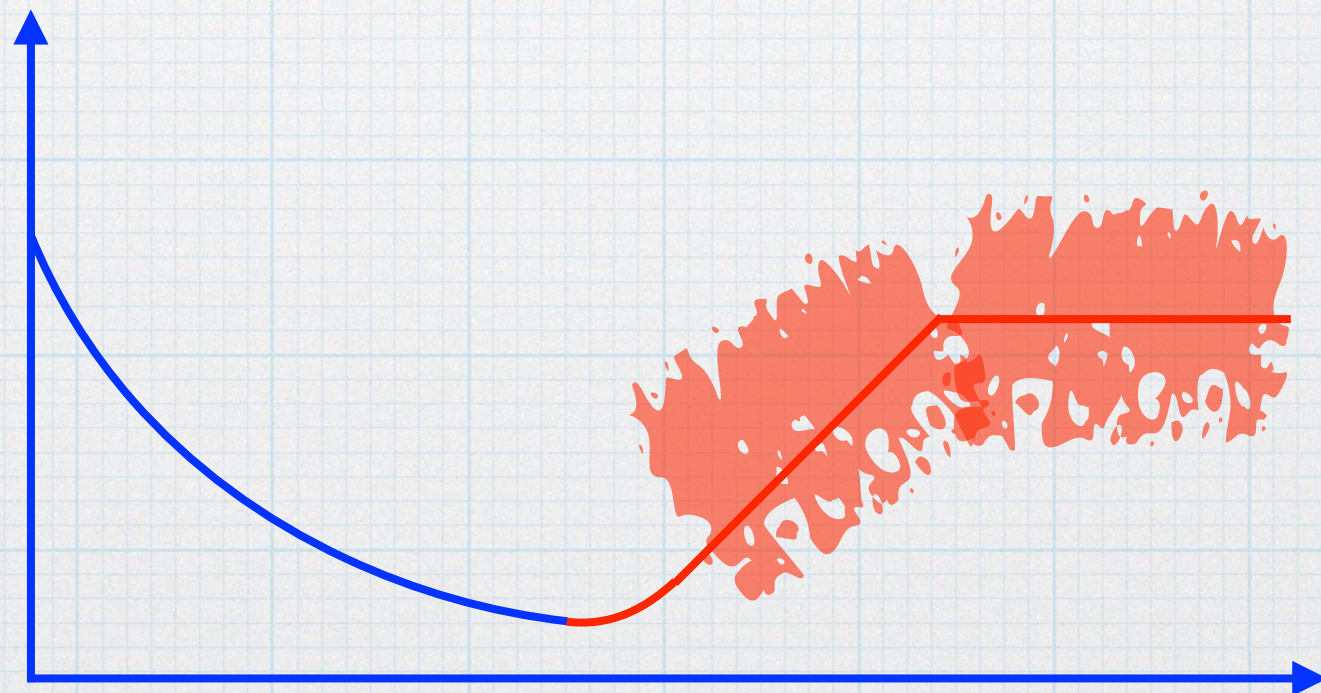
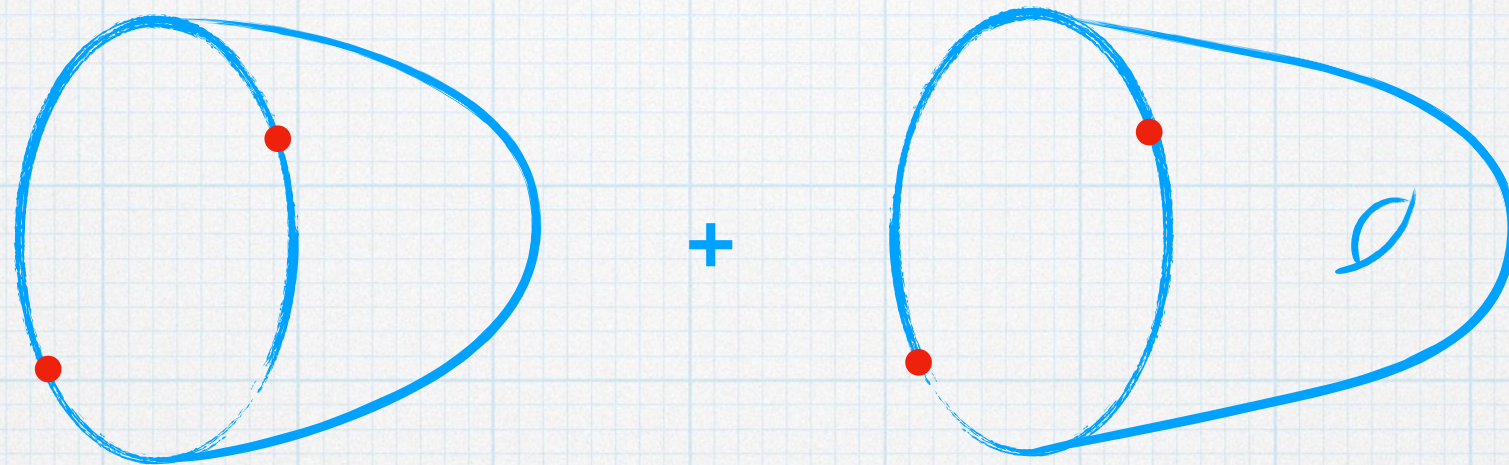
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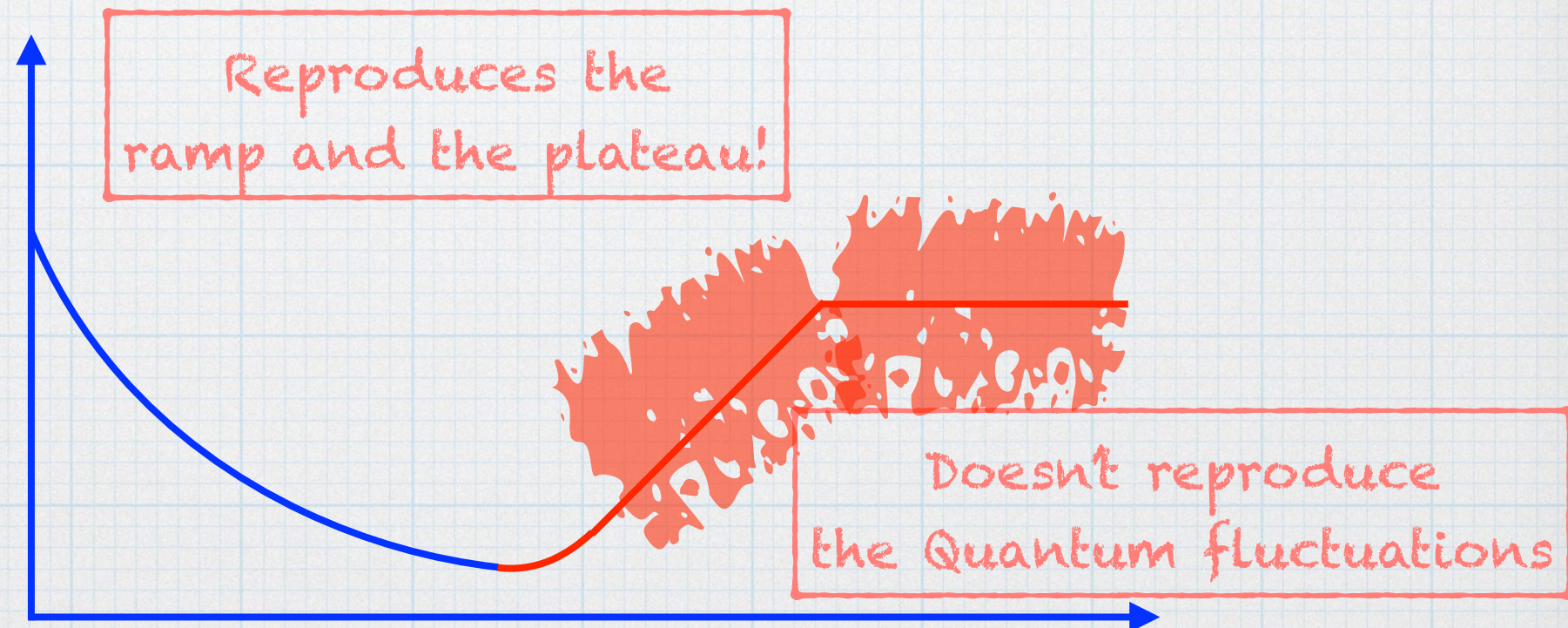
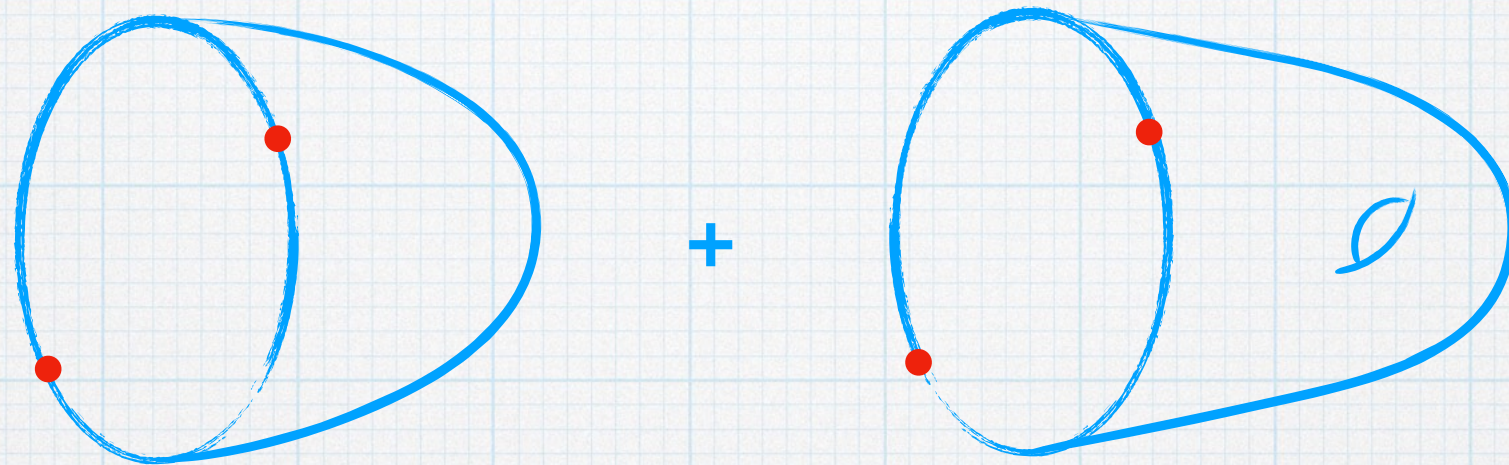
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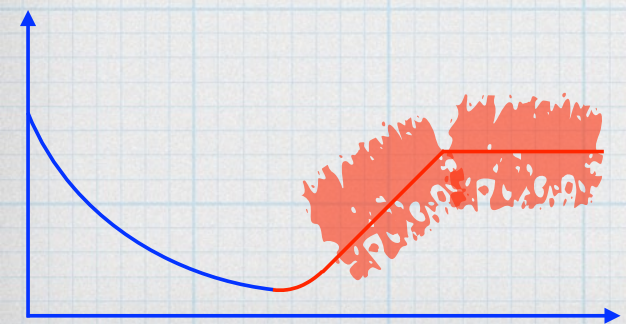
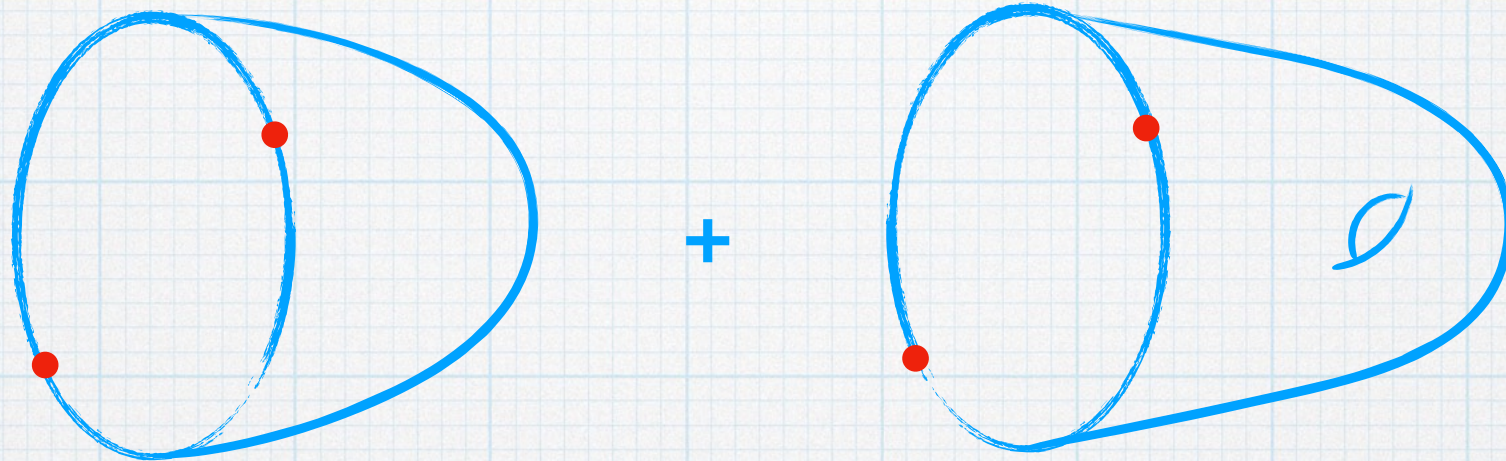
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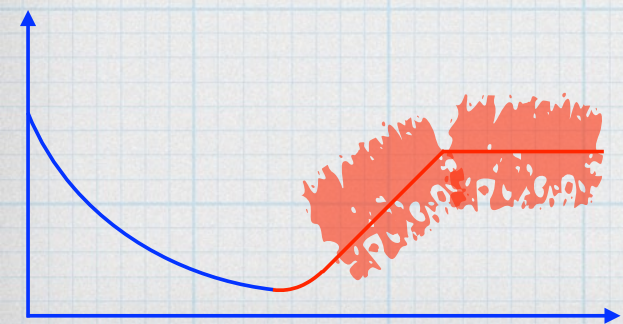
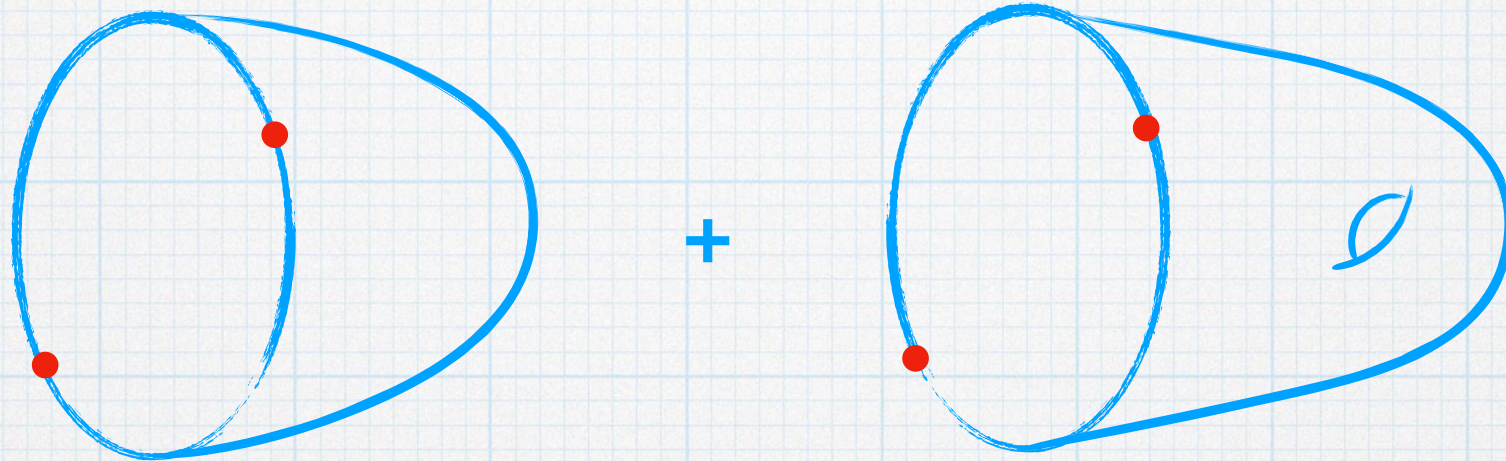
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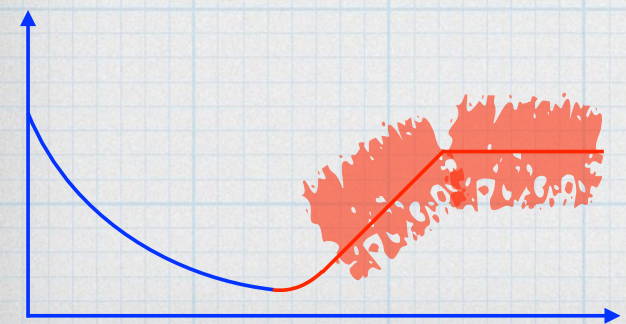
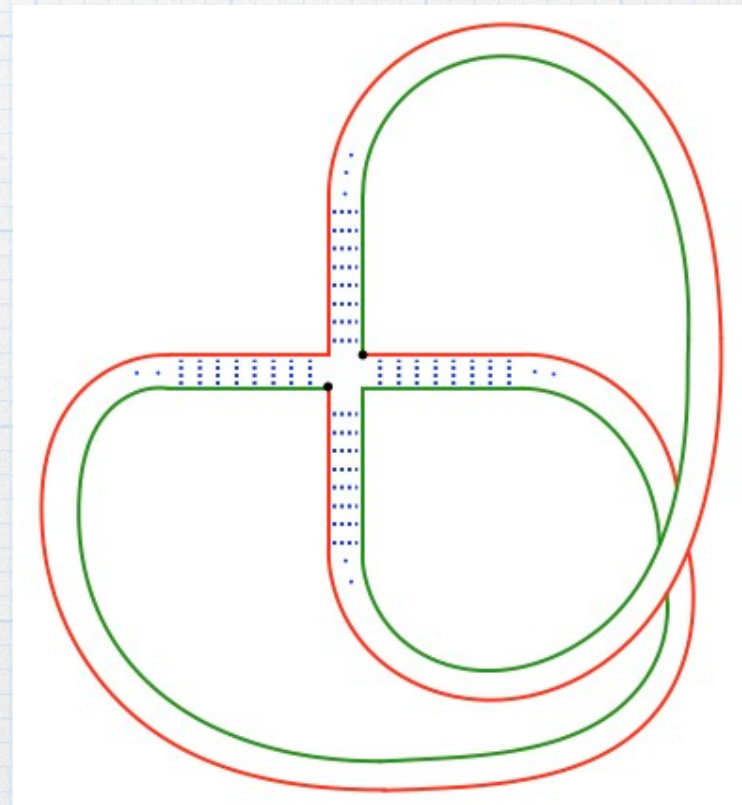
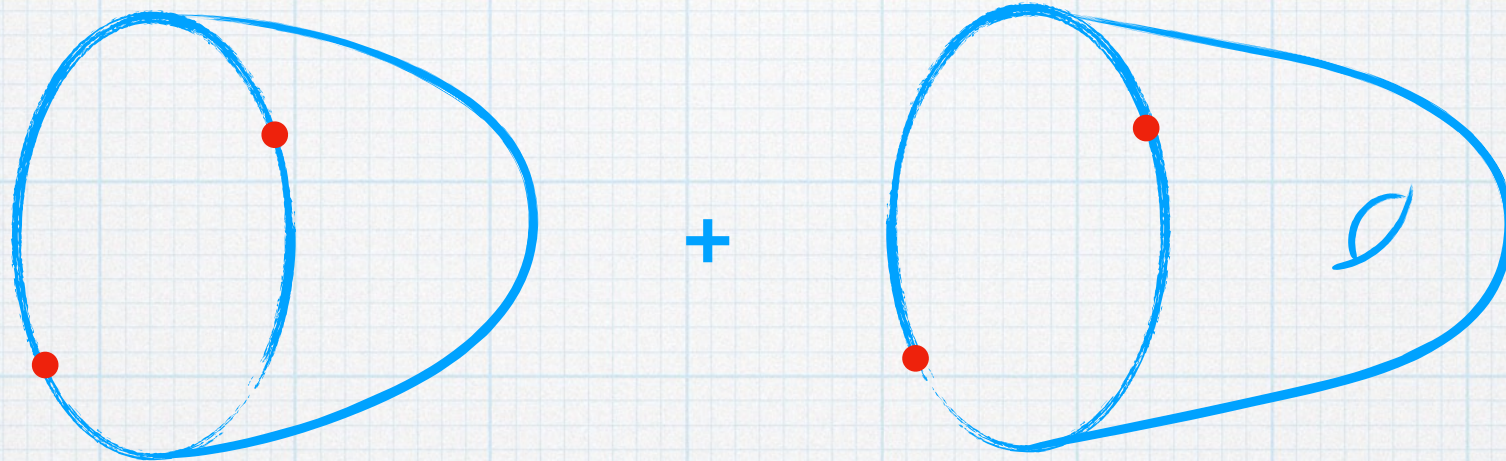
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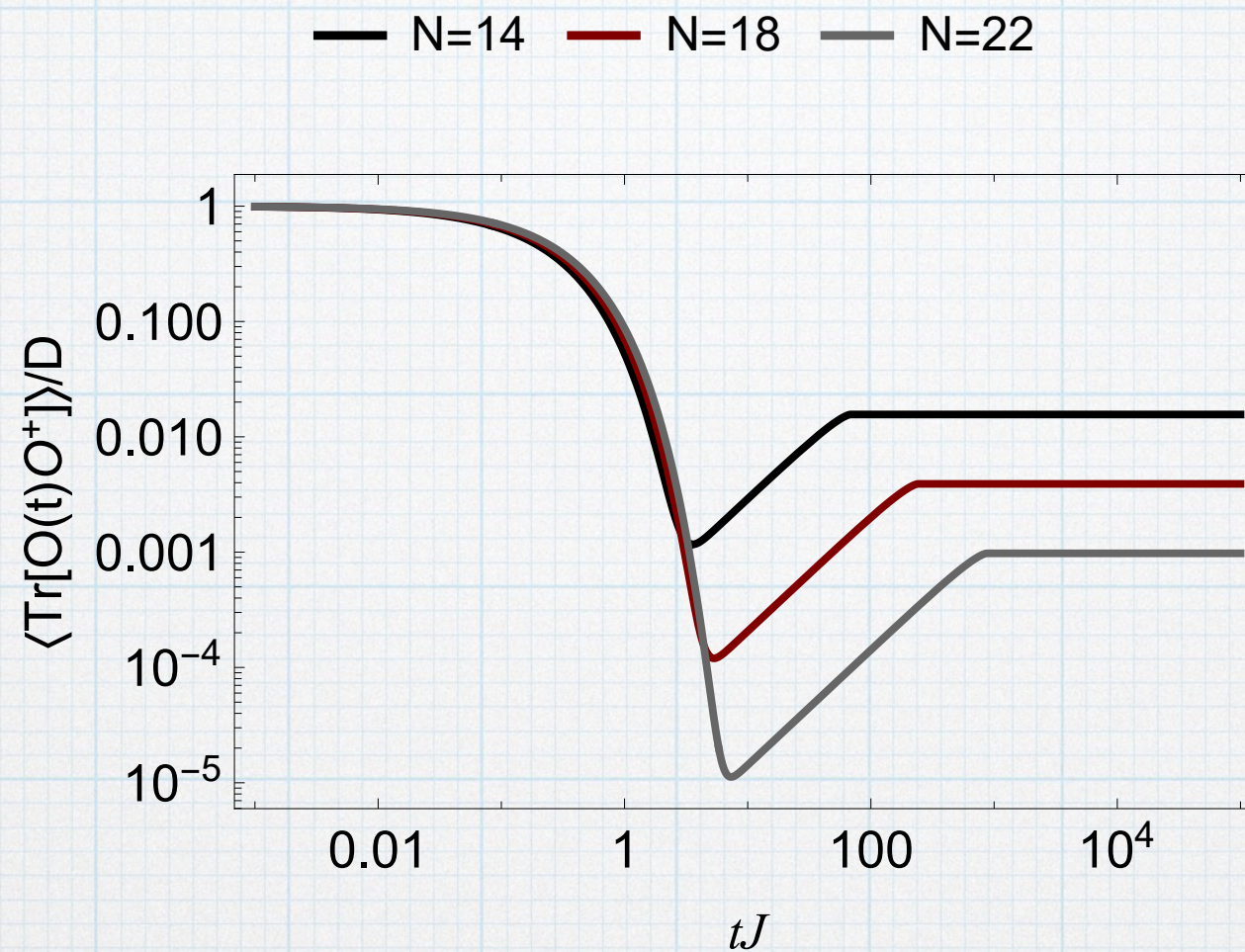
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Fewer random parameters in $H \Rightarrow$ corrections to mean field approximation,

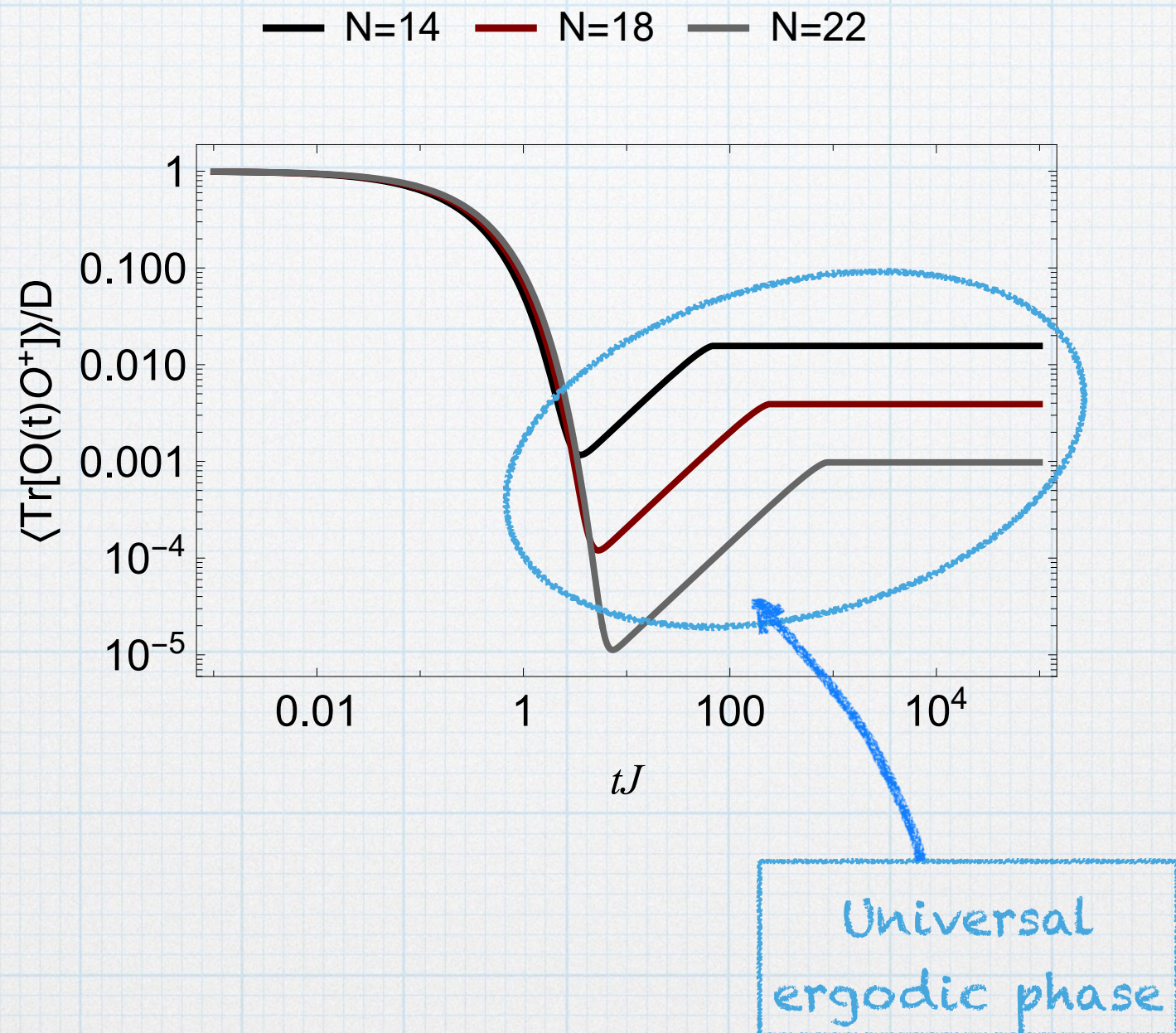
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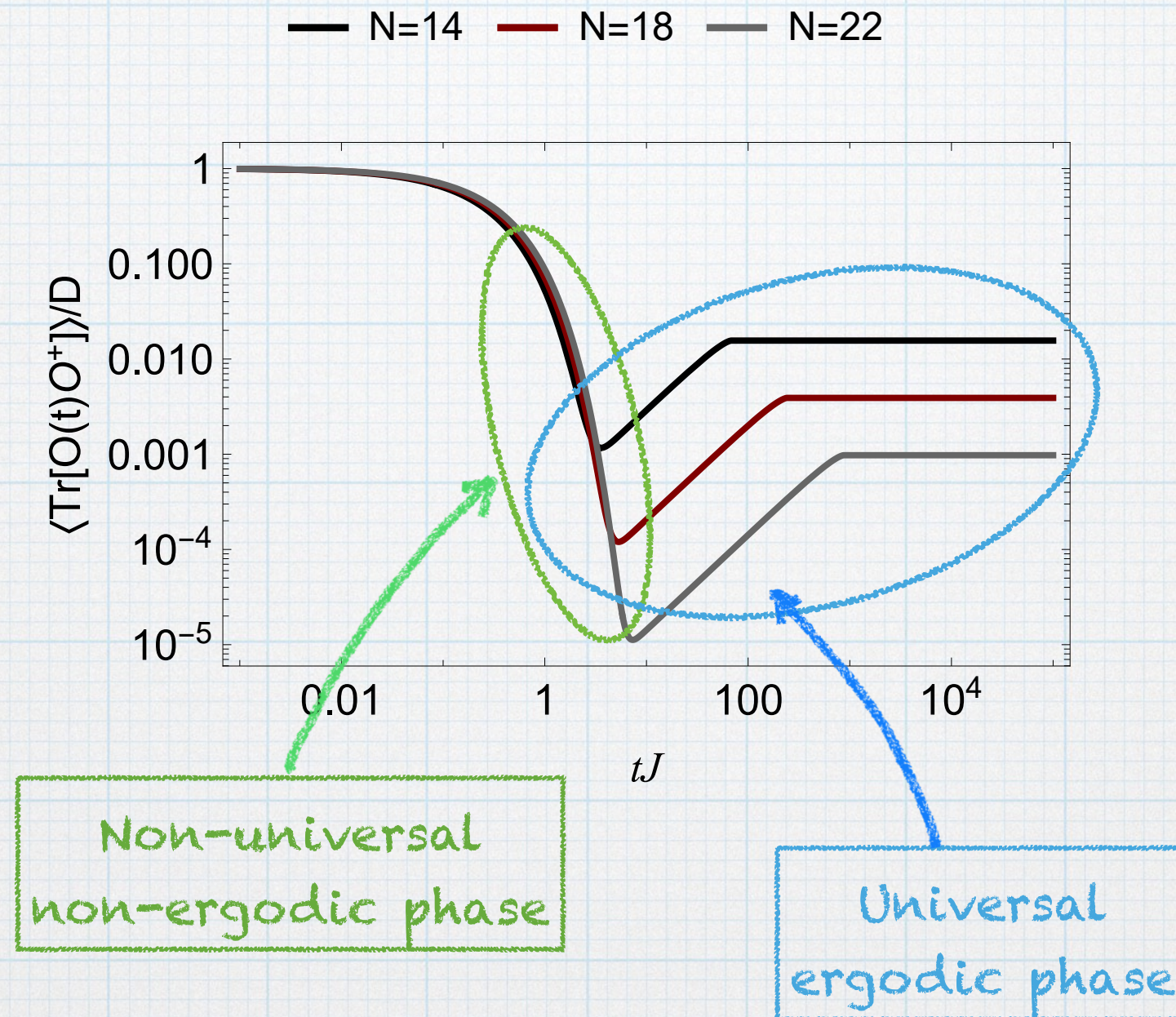
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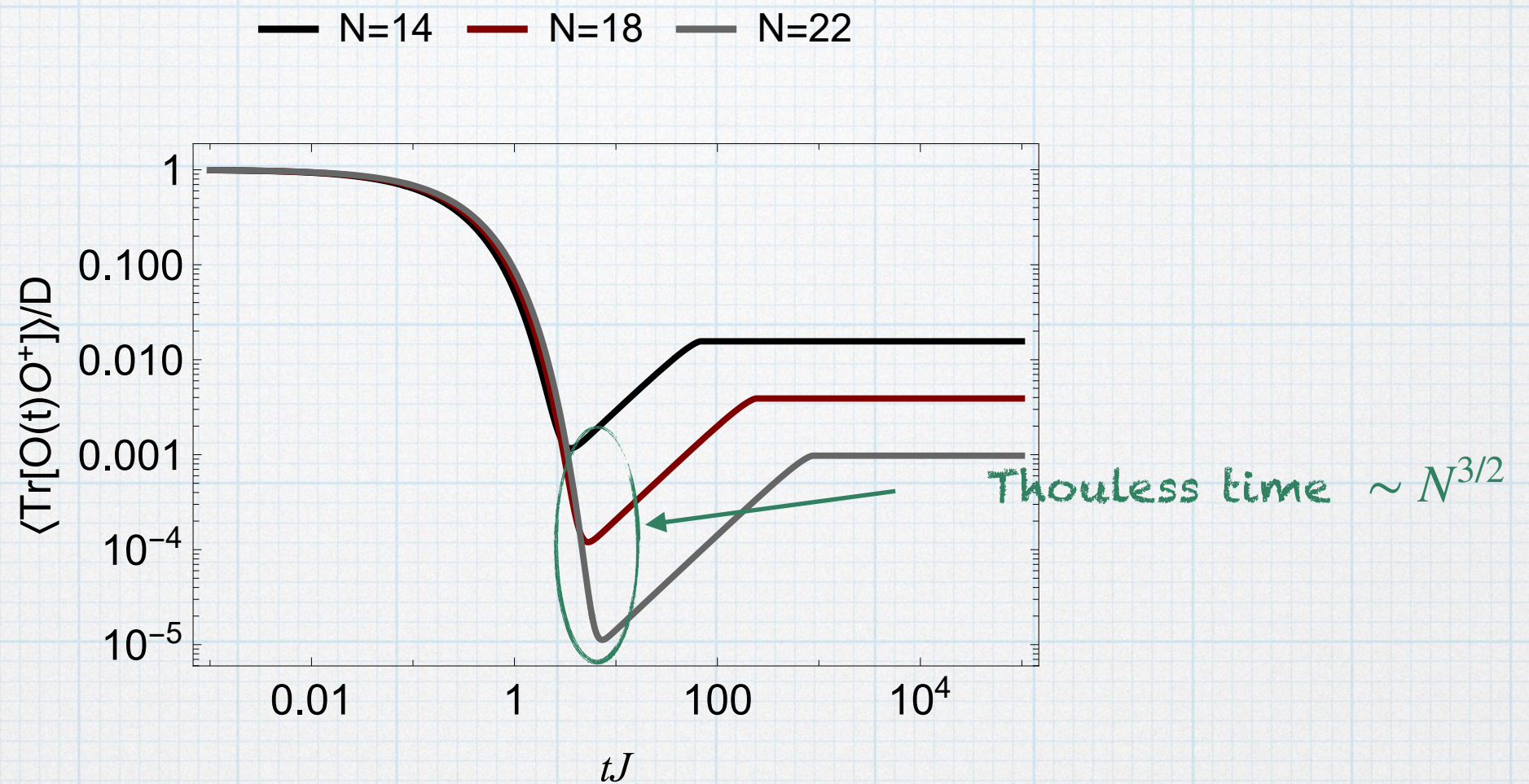
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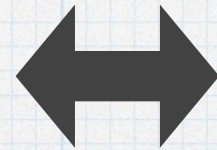


Non-ergodic Regime



New Conjecture

Semi-classical Gravity



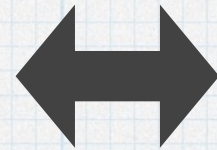
**Ergodic limit
of
Quantum theories**

[Altland, Sonner '20; Altland, Bagrets, PN, Sonner, Vielma '21]

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Semi-classical gravity has more information about the fine-grained structure of the spectrum of the Quantum theory!

Semi-classical Gravity



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Ergodic Predictions in CFT₂

Belin, de Boer '20;
Belin, de Boer, *PN*, Sonner
2111.xxxxx

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Statistics of OPE coefficients $\langle\langle c_{ijk} c_{lmn}^* \rangle\rangle$ in ergodic regime

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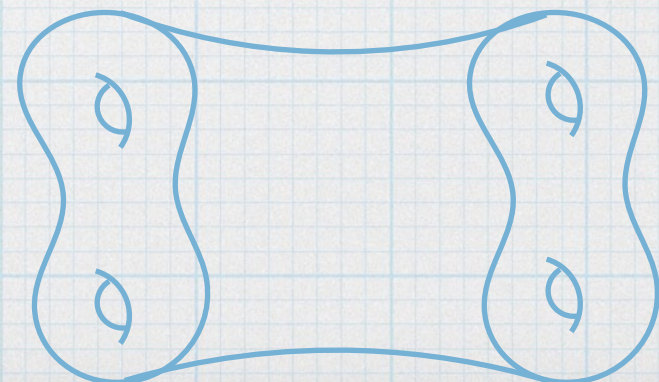
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☑ Reproduce the correct entropic suppressions as predicted

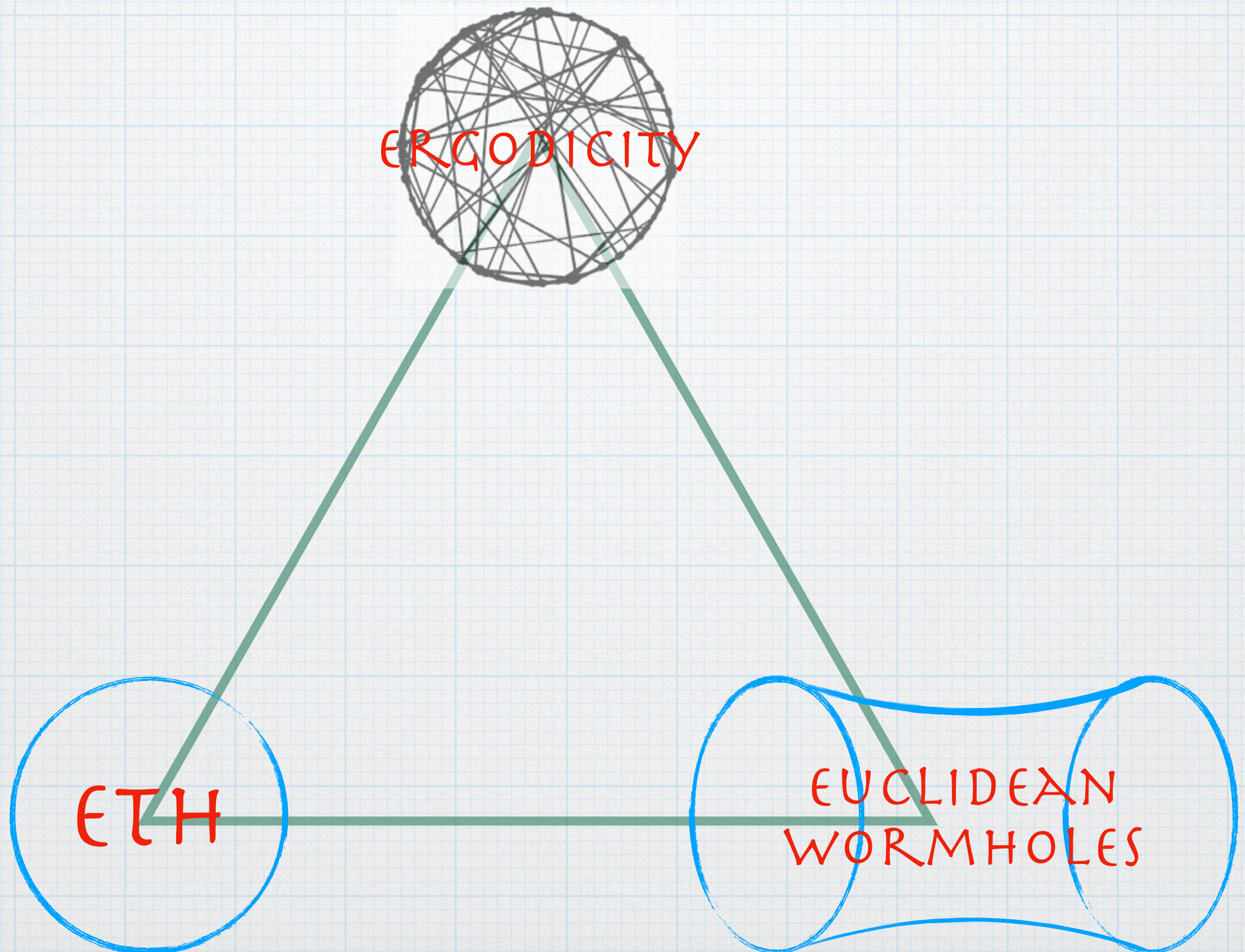
ERGODICITY

The Connection

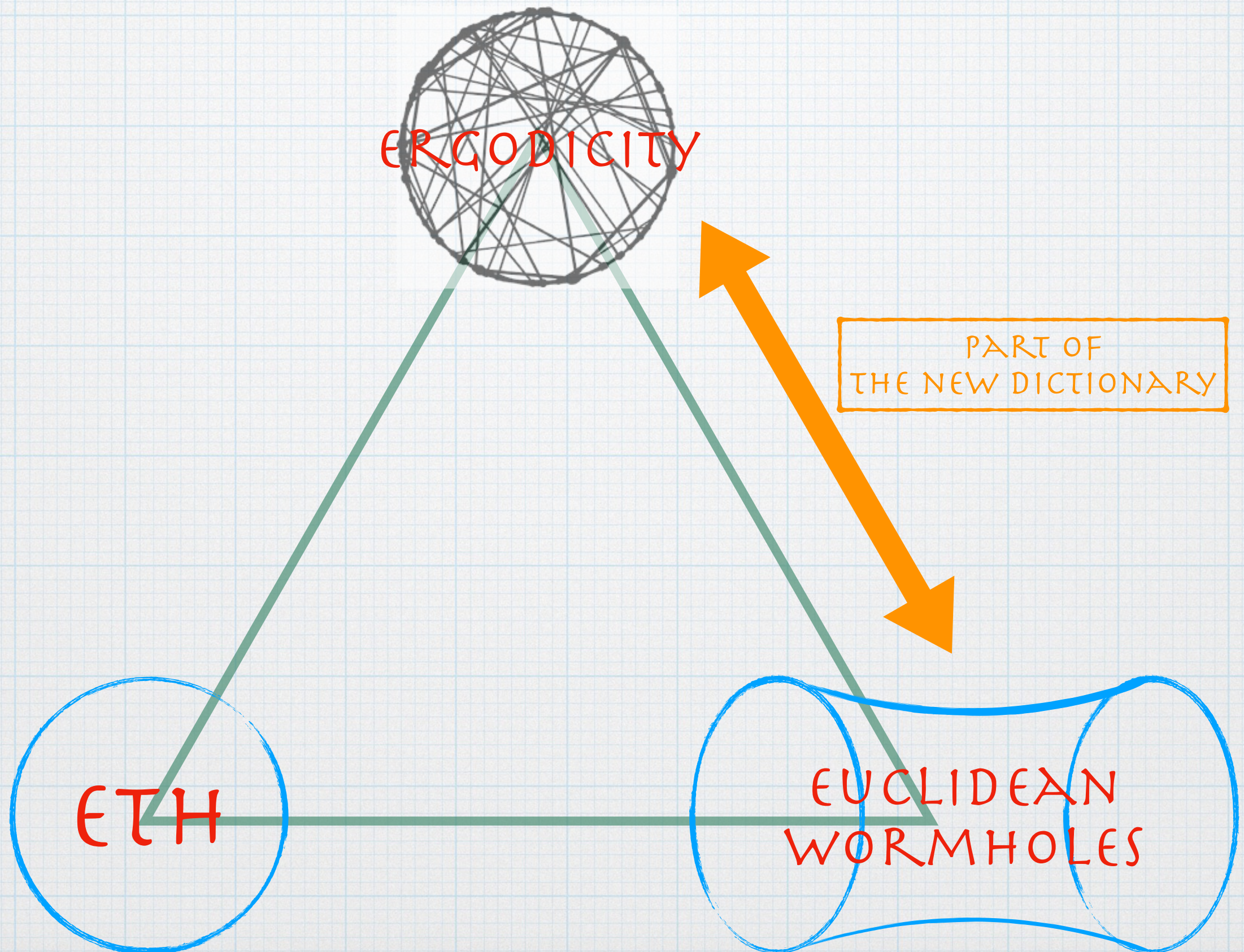
ETH

EUCLIDEAN
WORMHOLES

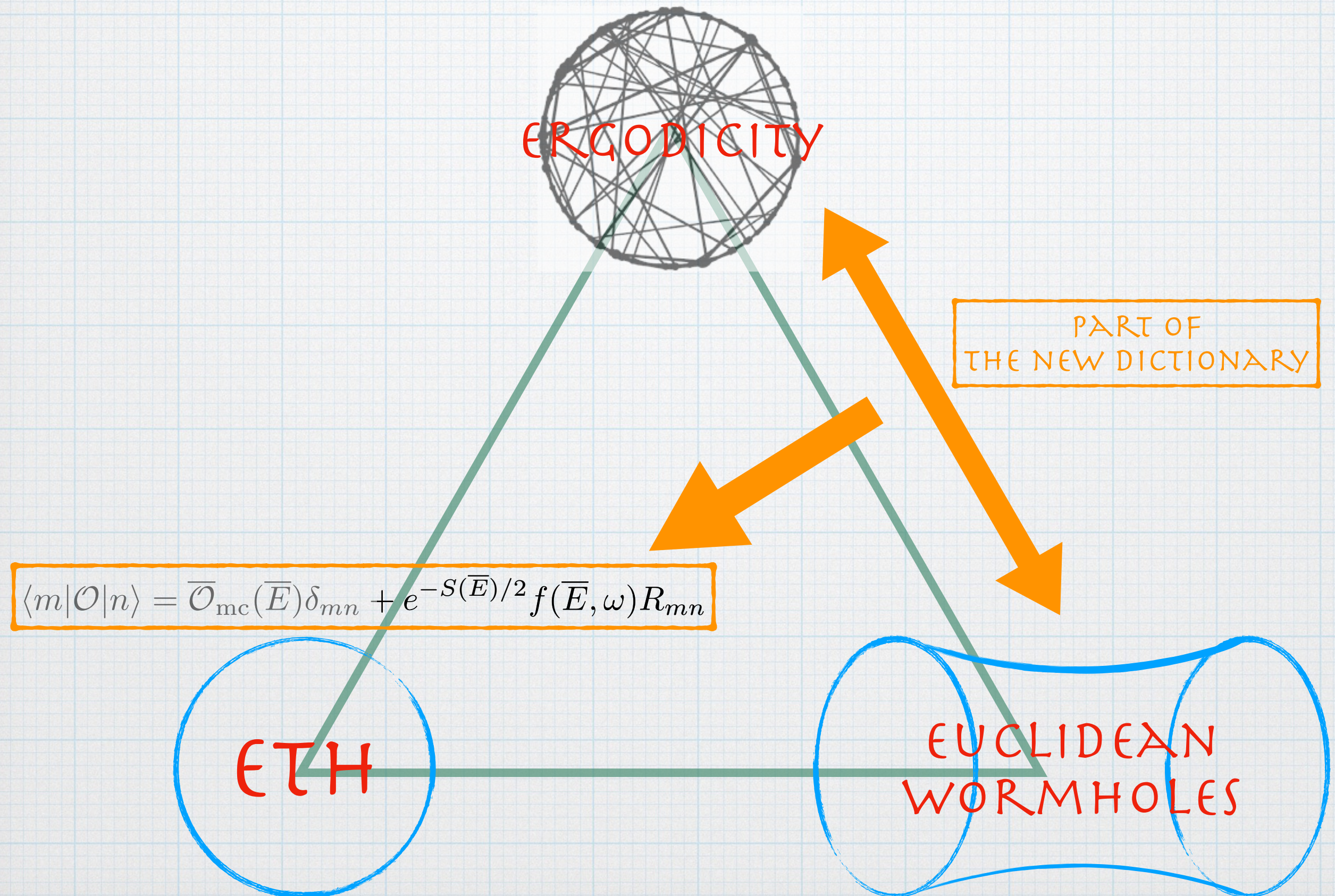
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ETH for charged operators

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Belin, de Boer, *PN*, Sonner '20

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- * A Tale of Two Conjectures:

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Belin, de Boer, PN, Sonner '20

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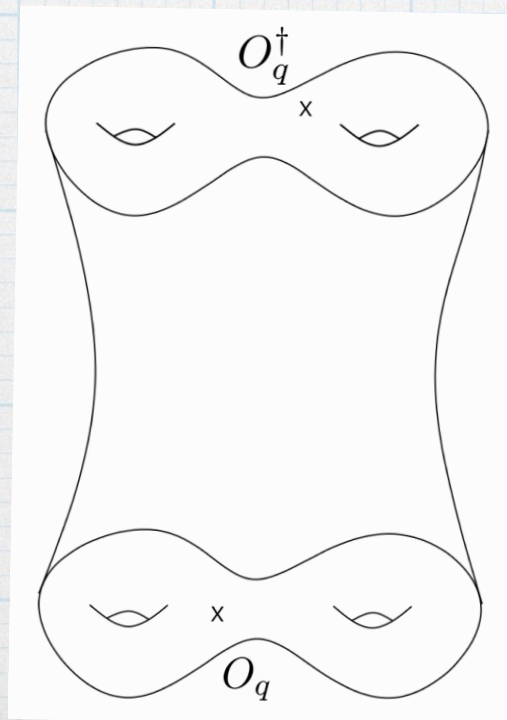
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* Wormholes or no wormholes



Summary

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TO DO

- What is the bulk interpretation of the **quantum fluctuations**?
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Arriving at sigma model

$$Z[h] = \left\langle \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \exp \left[i\bar{\Psi} \cdot (z - H + h) \cdot \Psi \right] \right\rangle_{\text{dis}}$$

$$Z[h] = \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \mathcal{D}A \exp \left[i\bar{\Psi} \cdot (z + h) \cdot \Psi - \frac{1}{2n} \sum_a \text{STr} [X^a A^a X^a A^a] + \frac{i\gamma}{n} \sum_a \text{STr} [\Psi \bar{\Psi} A^a] \right]$$

$$Z_0[h] = \int \mathcal{D}y \exp \left[-\frac{D}{2} \text{STr}(y^2) - \text{STr} \ln(K_0) \right], \quad K_0 = \left(z + h + \gamma y \otimes \mathbb{I}^{\text{Fock}} \right)$$

Sigma model

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$$\Lambda_{AA} = \sigma_3^{\text{RA}} \otimes \sigma_3^{\text{bf}}$$

$$Z[h] = \int \mathcal{D}Q \exp \left[\frac{i\pi\rho(E)}{2} \text{Str}(\omega\Lambda Q) - S_{\text{src}}(h) \right], \quad Q = T\Lambda T^{-1}$$

$$T = \exp(-W), \quad W = - \begin{pmatrix} 0 & B \\ \tilde{B} & 0 \end{pmatrix}, \quad B = \begin{pmatrix} x & \mu \\ \nu & iy \end{pmatrix}, \quad \tilde{B} = \begin{pmatrix} x^* & \bar{\nu} \\ -\bar{\mu} & iy^* \end{pmatrix}.$$

$$\partial_h^2 Z[h] \Big|_{h=0} = \int \mathcal{D}Q \exp \left[2i\pi\rho(E)\omega \text{Str} [B\tilde{B}] + \dots \right] \times \text{STr} [\tilde{B}P_b\tilde{B}BP_bB]$$

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Sources for various correlation functions

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- * Sources for the operator correlation function that we have been looking at,
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