

A KRYLOV SPACE ODYSSEY

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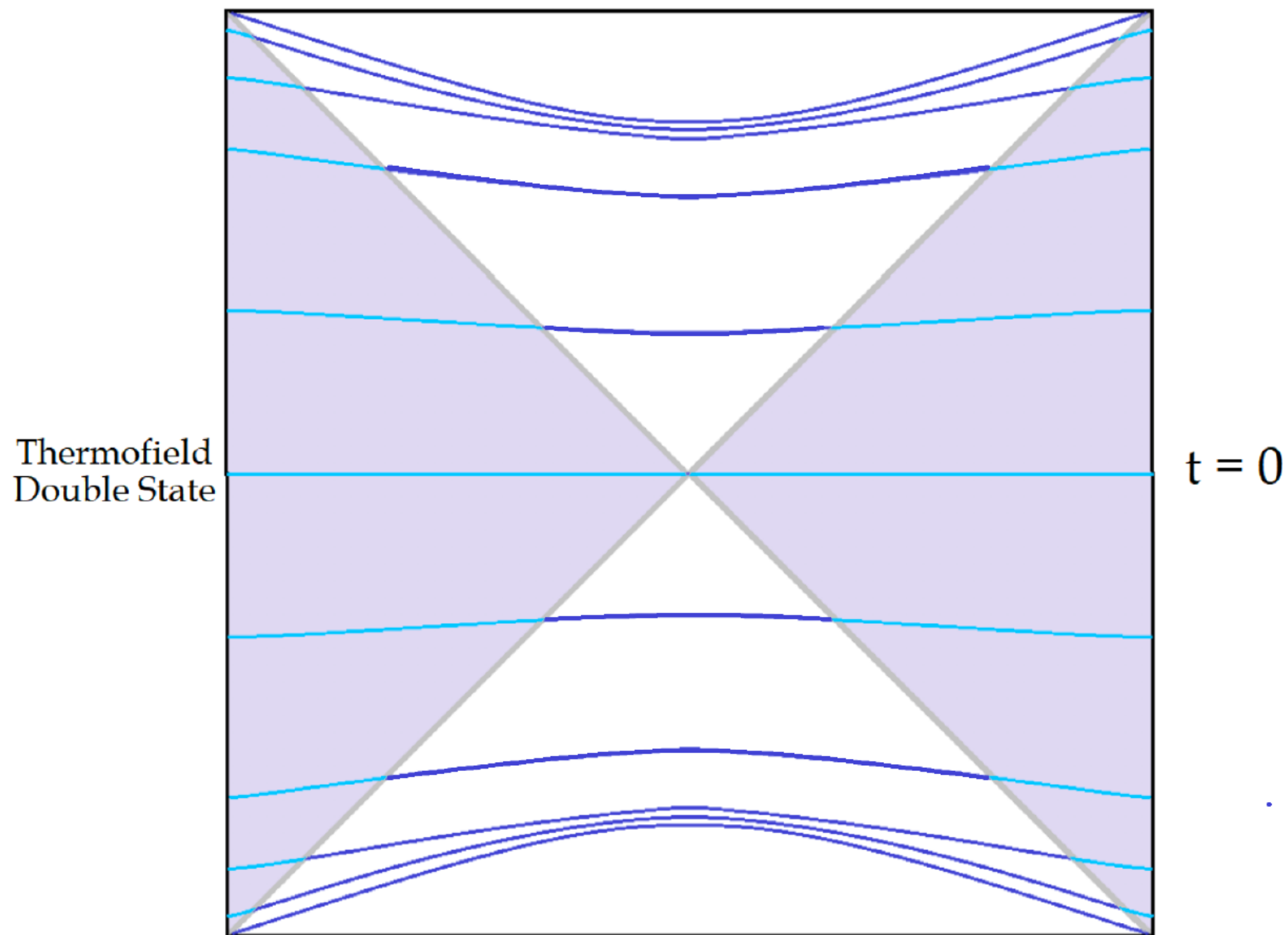
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COMPLEXITY IN HOLOGRAPHY?



(L. Susskind, 2018.)

- What boundary observable describes the growth of the Einstein-Rosen Bridge (ERB)?
- It *could* be complexity (L. Susskind, 2013).

But what complexity??

- Should saturate at Heisenberg time, $t_H \sim e^S$.
- No *external* tolerance parameter.

STATES OR OPERATORS?

- ERB of eternal AdS black hole should be captured by the evolution of $|TFD\rangle$ in the **Schrödinger picture**.
- *Butterfly effect*: backreacted geometry due to shockwave of infalling particle (Shenker & Stanford, 2013).
 - Captured by evolution of $|TFD\rangle$ perturbed by an operator insertion $\mathcal{O}(t)$.
 - Motivates the study of **operator complexity in the Heisenberg picture**.

This gives access to a wide range of time scales:

- **Scrambling time**, $t_s \sim \log S$ (*infalling particle crosses the BH horizon*).
- **Heisenberg time**, $t_H \sim e^S$ (*geometrical description fails, Hilbert space fully explored, fine-grained spectrum is probed...*)

KRYLOV SPACE

- A notion of complexity adapted to time evolution of an initial operator $\mathcal{O} \equiv \mathcal{O}(0)$.
- Take a Hilbert space of states $\mathcal{H}$ with $\dim \mathcal{H} = D$.
 $\implies$ Operator space $\widehat{\mathcal{H}}$ will have $\dim \widehat{\mathcal{H}} = D^2$.
- Time evolution generator in $\widehat{\mathcal{H}}$ is the **Liouvillian** $\mathcal{L} := [H, \cdot]$, as:

$$\mathcal{O}(t) = e^{iHt} \mathcal{O} e^{-iHt} = e^{it\mathcal{L}} \mathcal{O} = \mathcal{O} + it[H, \mathcal{O}] - \frac{t^2}{2} [H, [H, \mathcal{O}]] + \dots$$

- Define **Krylov space** as $\mathcal{H}_{\mathcal{O}} := \text{span} \left\{ \mathcal{L}^n \mathcal{O} \right\}_{n=0}^{+\infty}$

$\implies$ Always contains $\mathcal{O}(t)$. What is its dimension, K ??

KRYLOV SPACE DIMENSION AND UPPER BOUND

The spectrum of $\mathcal{L}$ is made up of all possible energy differences $E_a - E_b \equiv \omega_{ab}$:

$$|E_a\rangle\langle E_b| \equiv |\omega_{ab}\rangle \implies \mathcal{L} |\omega_{ab}\rangle = \omega_{ab} |\omega_{ab}\rangle$$

Spectral decomposition of $|\mathcal{O}\rangle = \sum_{ab} O_{ab} |\omega_{ab}\rangle \implies \mathcal{L}^n |\mathcal{O}\rangle = \sum_{ab} \omega_{ab}^n O_{ab} |\omega_{ab}\rangle$

$\implies K =$ Number of non-zero eigenspace projections of $|\mathcal{O}\rangle$!!

Immediate upper bound: $K \leq D^2 - D + 1$ $\left\{ \begin{array}{l} \text{- Saturated in chaotic systems (ETH).} \\ \text{- Degeneracies reduce } K \end{array} \right.$

LANCZOS COEFFICIENTS AND KRYLOV CHAIN DYNAMICS

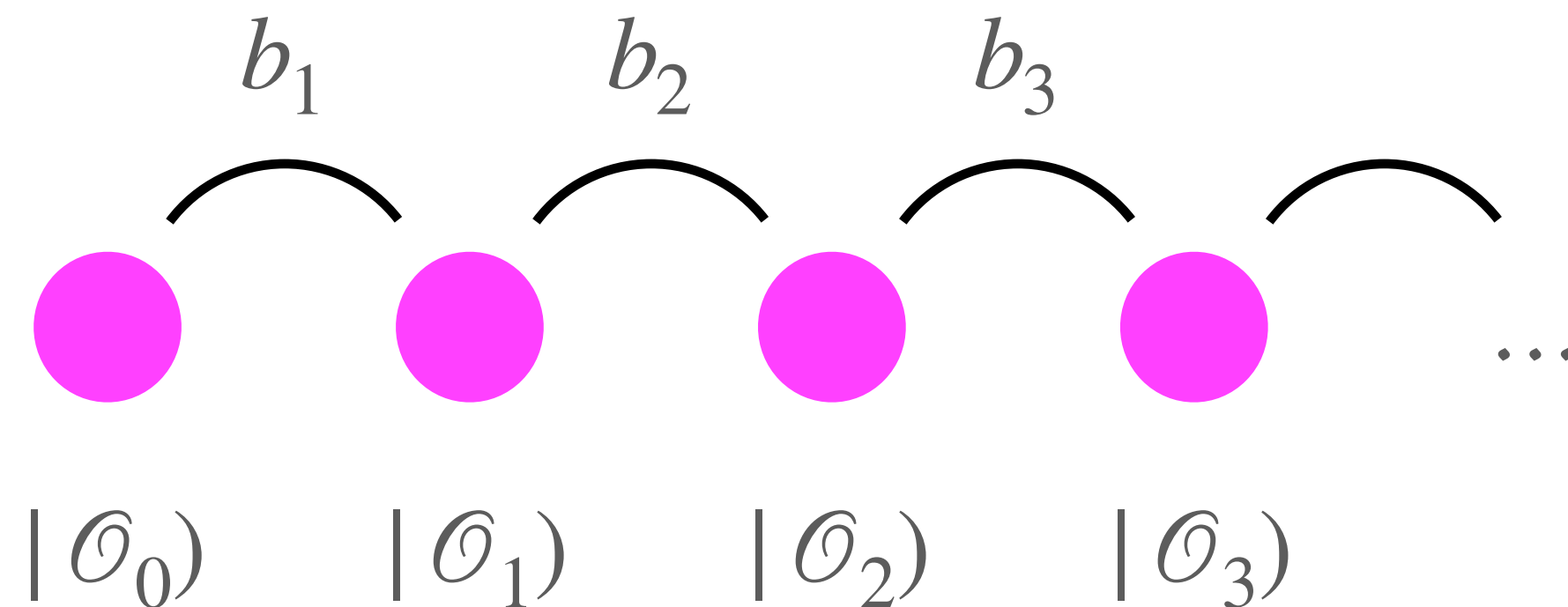
- We can build an *ordered* orthonormal basis for Krylov space (**Lanczos algorithm**).

$$\{ |\mathcal{O}\rangle, \mathcal{L}|\mathcal{O}\rangle, \mathcal{L}^2|\mathcal{O}\rangle, \dots \} \mapsto \{ |\mathcal{O}_0\rangle = |\mathcal{O}\rangle, |\mathcal{O}_1\rangle, |\mathcal{O}_2\rangle, \dots \}$$

- Re-orthogonalization coefficients: **Lanczos coefficients**, $\{b_n\}$.

Krylov elements

$\Rightarrow$ Liouvillian is tri-diagonal! $(\mathcal{L}_{mn}) = \begin{pmatrix} 0 & b_1 & 0 & 0 & \dots & 0 \\ b_1 & 0 & b_2 & 0 & \dots & 0 \\ 0 & b_2 & 0 & b_3 & \dots & 0 \\ \vdots & & \ddots & & \ddots & \vdots \end{pmatrix}$



1d hopping model!

K-COMPLEXITY AND K-ENTROPY

- Evolution of $|\mathcal{O}(t)\rangle$ encoded in wave function $\varphi(t)$:

$$|\mathcal{O}(t)\rangle = \sum_n i^n \varphi_n(t) |\mathcal{O}_n\rangle$$

K-COMPLEXITY:

$$C_K(t) = \sum_n n \left| \varphi_n(t) \right|^2$$

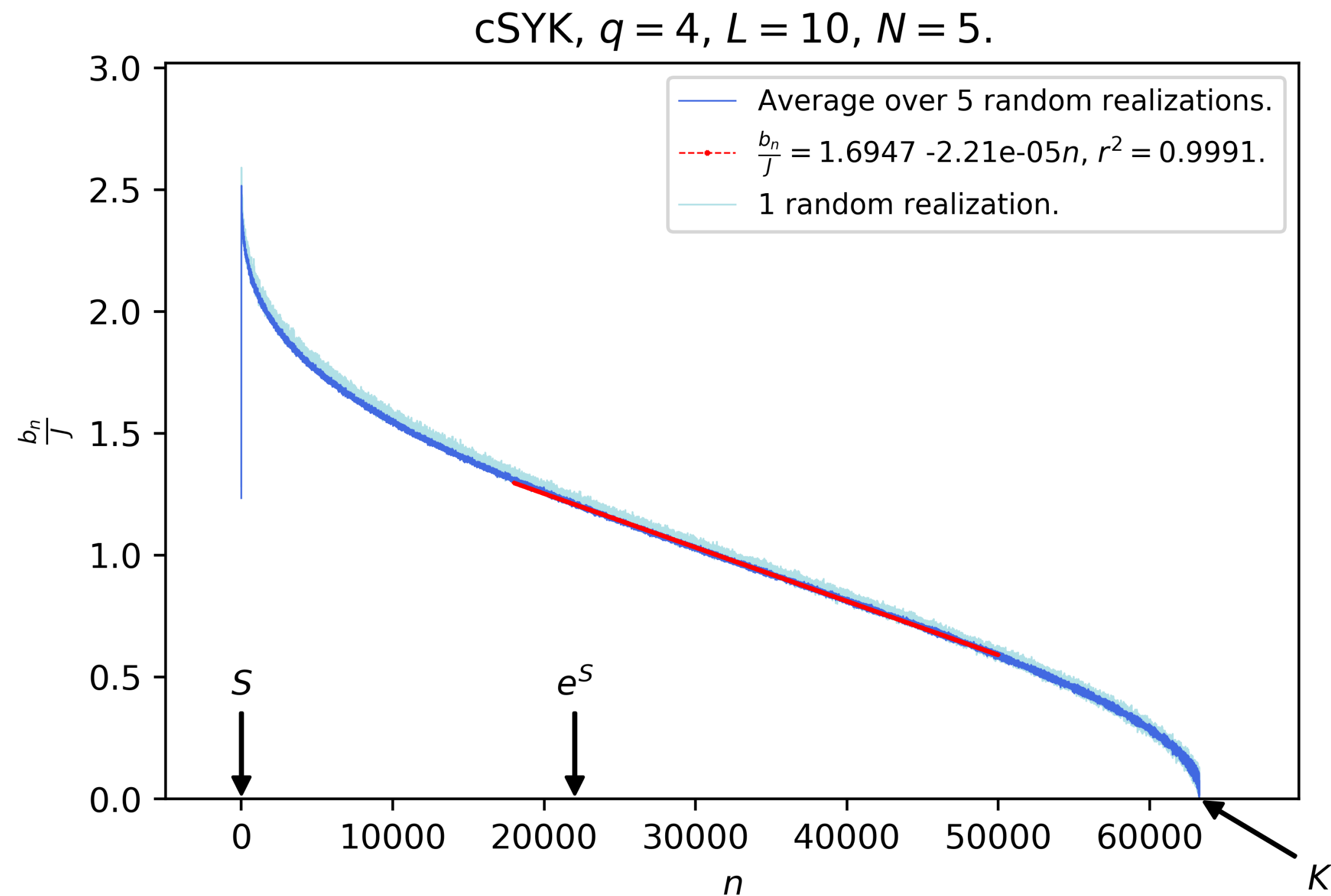
K-ENTROPY:

$$S_K(t) = - \sum_n \left| \varphi_n(t) \right|^2 \log \left| \varphi_n(t) \right|^2$$

For maximally chaotic systems:

- **Thermodynamic limit** (Altman et al., 2019): $b_n \sim \alpha n$, $C_K(t) \sim e^{2\alpha t}$, $S_K(t) \sim t$
- Finite size, **post-scrambling** (Barbón et al., 2019): $b_n \rightarrow b_\infty$, $C_K(t) \sim t$, $S_K(t) \sim \log(t)$

LANCZOS COEFFICIENTS IN SYK

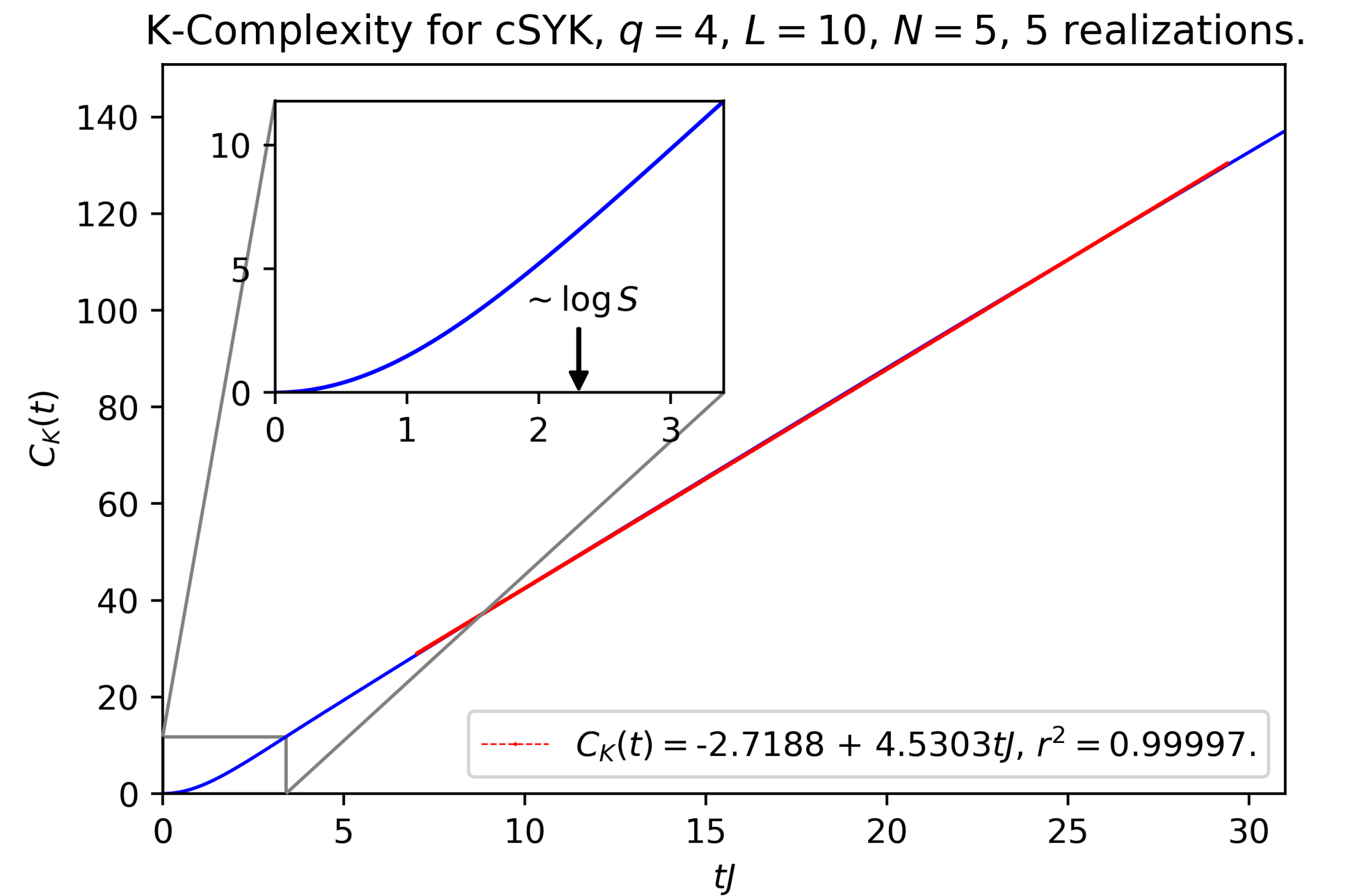
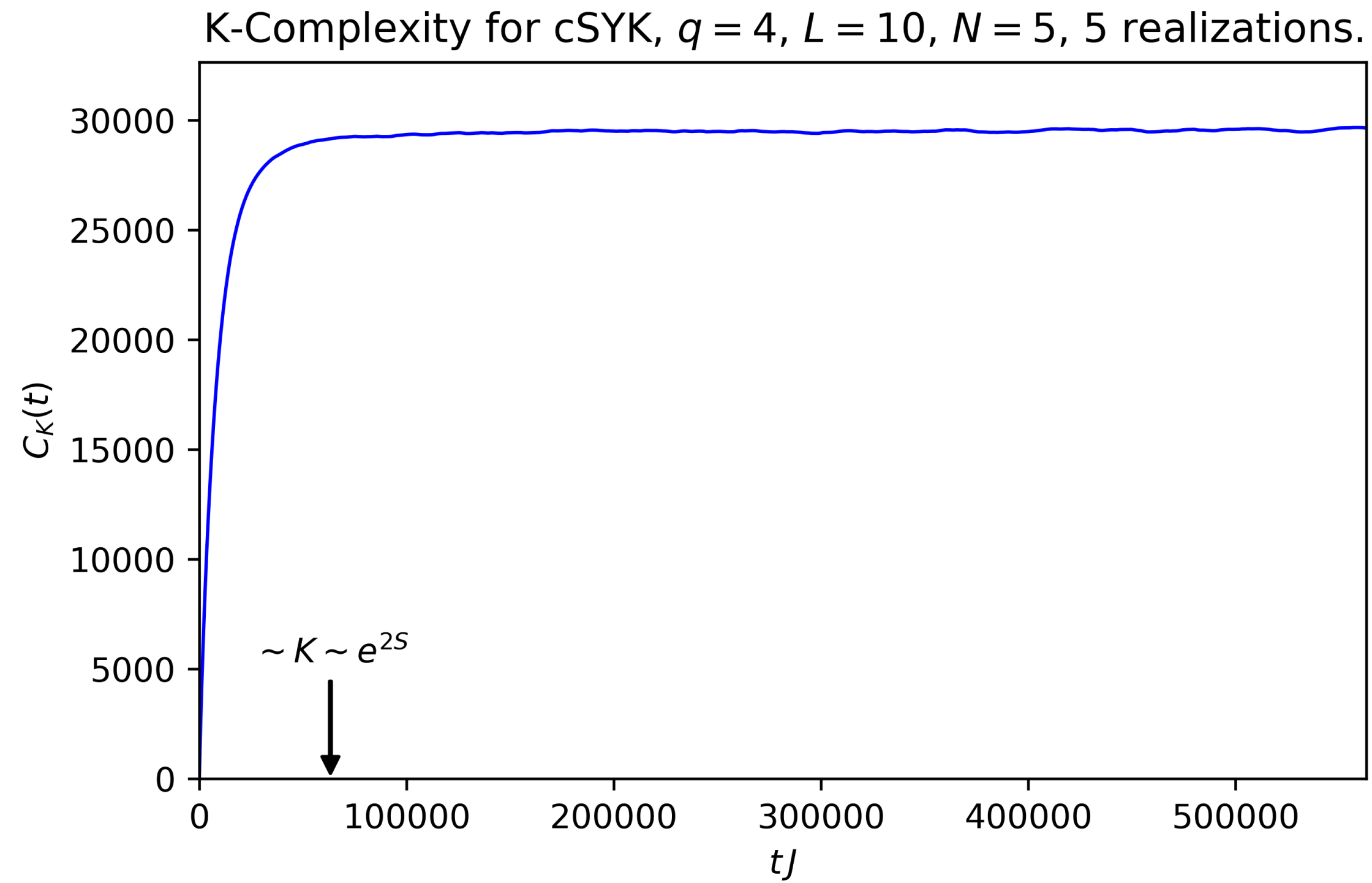


$$H = \sum_{ijkl} J_{ij;kl} c_i^\dagger c_j^\dagger c_k c_l; \quad \mathcal{O} = c_1^\dagger c_2 + \text{h.c.}$$

- Upper bound saturated, $K = D^2 - D + 1$.
- Linear growth up to $n \sim S$.
- **Non-perturbative descent** with slope $\sim e^{-2S}$.
- Reaches zero at $n = K \sim e^{2S}$.

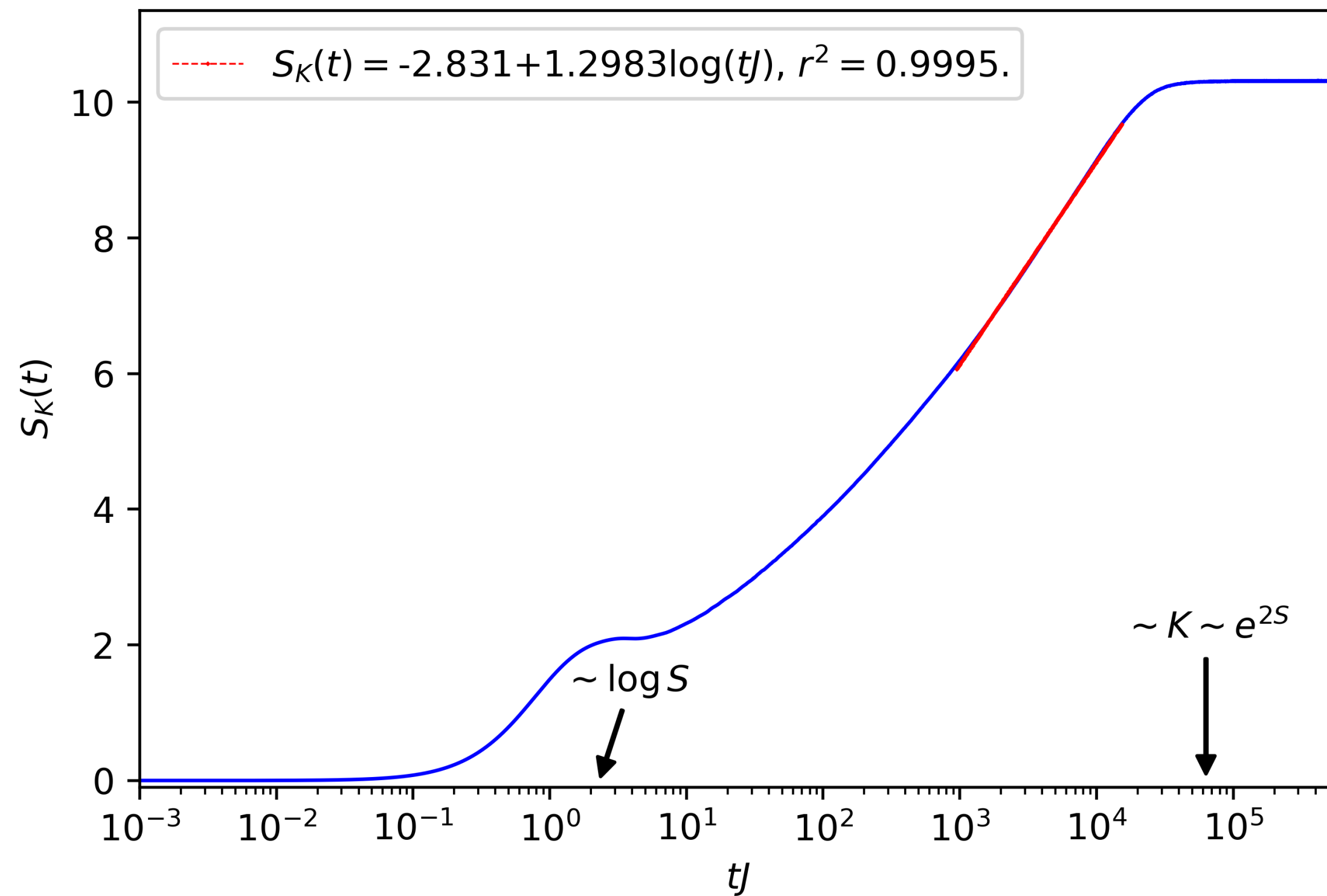
(E. Rabinovici, ASG, R. Shir, J. Sonner, 2020.)

K-COMPLEXITY IN SYK

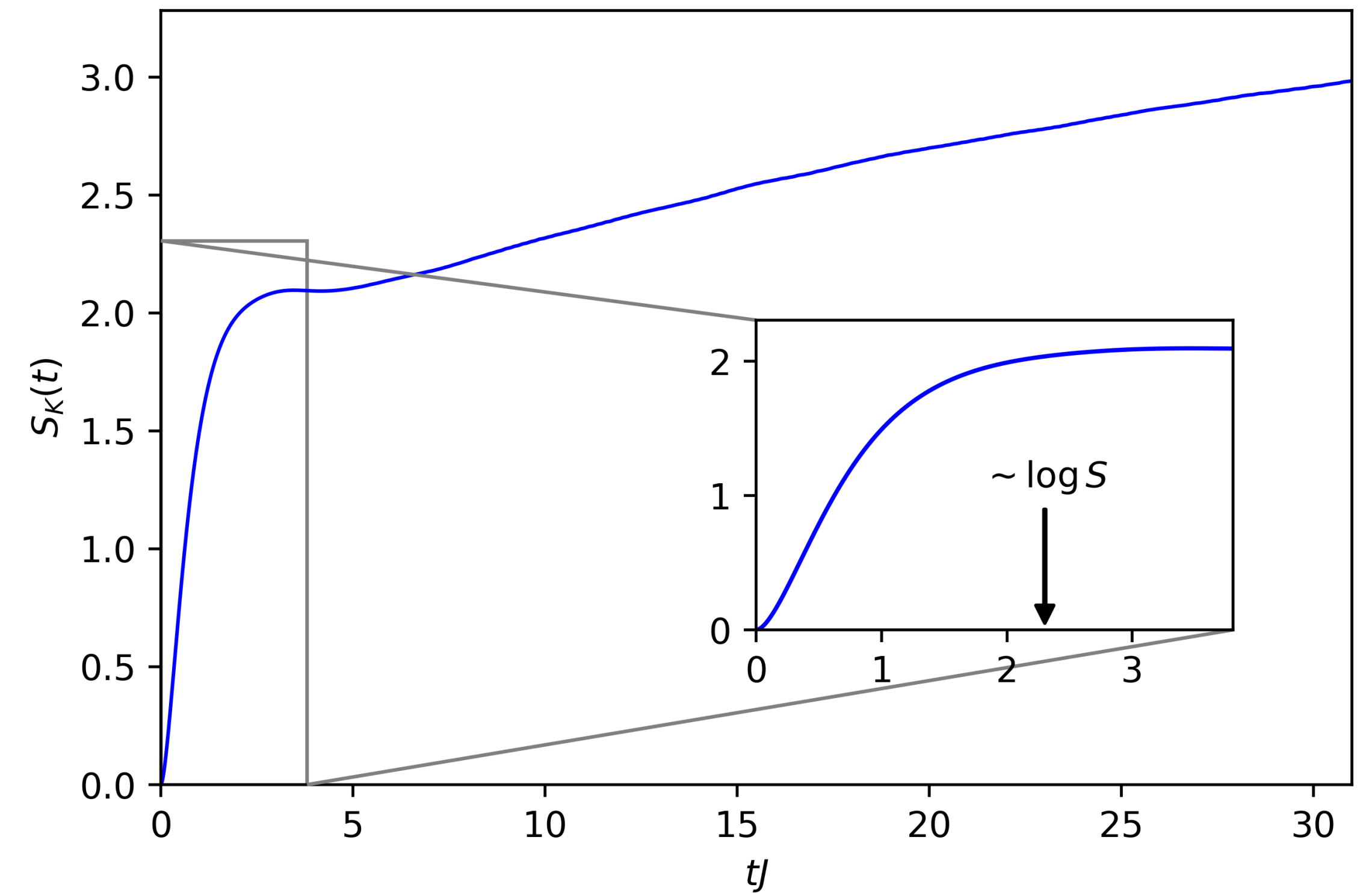


K-ENTROPY IN SYK

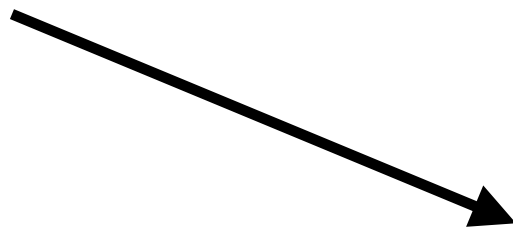
K-Entropy for cSYK, $q = 4$, $L = 10$, $N = 5$, 5 realizations.



K-Entropy for cSYK, $q = 4$, $L = 10$, $N = 5$, 5 realizations.



CONCLUSIONS

- K-complexity is a notion of complexity that doesn't depend on an external tolerance parameter.
 - It is very sensitive to the spectral structure.
 - Non-perturbative descent of $b_n \longrightarrow$ Boundedness and saturation of $C_K(t)$.
 - Some questions:
 - Bulk interpretation?
 - Problems: Initial condition dependence; time-dependent Hamiltonians...
 - Probe of quantum chaos? (*Dymarsky & Smolkin, 2021; work in progress...*)
- 
- Late time physics!

i Thanks !

KRYLOV SPACE DIMENSION – VANDERMONDE ARGUMENT

- Liouvillian eigenvalues: $\omega_{ab} = E_a - E_b \implies \omega_{aa} = 0, \forall a = 1, \dots, D.$

$$\implies \mathcal{L}^n |\mathcal{O}\rangle = \delta_{n0} \sum_{a=1}^D O_{aa} |\omega_{aa}\rangle + \sum_{a \neq b} \omega_{ab}^n O_{ab} |\omega_{ab}\rangle$$

- In coordinates:

$$\begin{pmatrix} \mathcal{O} \\ \mathcal{L}\mathcal{O} \\ \mathcal{L}^2\mathcal{O} \\ \vdots \end{pmatrix} = \begin{pmatrix} O_{11} & O_{22} & \dots & O_{DD} & O_{12} & O_{13} & \dots & O_{D-1,D} \\ 0 & 0 & \dots & 0 & O_{12}\omega_{12} & O_{13}\omega_{13} & \dots & O_{D-1,D}\omega_{D-1,D} \\ 0 & 0 & \dots & 0 & O_{12}\omega_{12}^2 & O_{13}\omega_{13}^2 & \dots & O_{D-1,D}\omega_{D-1,D}^2 \\ \vdots & \vdots & \ddots & \vdots & \vdots & \vdots & \ddots & \vdots \end{pmatrix}$$

$\implies$ **Vandermonde matrix!!**

LANCZOS ALGORITHM

Construct orthonormal basis **given some inner product**, e.g. $(A | B) = \frac{1}{D} \text{Tr} [A^\dagger B]$.

1. Set $b_0 \equiv 0$ and $|\mathcal{O}_{-1}\rangle \equiv 0$.

$$2. \quad |\mathcal{O}_0\rangle = \frac{1}{\sqrt{(\mathcal{O} | \mathcal{O})}} |\mathcal{O}\rangle$$

3. For $n \geq 1$: $|\mathcal{A}_n\rangle = \mathcal{L} |\mathcal{O}_{n-1}\rangle - b_{n-1} |\mathcal{O}_{n-2}\rangle$

$$4. \quad b_n = \sqrt{(\mathcal{A}_n | \mathcal{A}_n)}$$

5. If $b_n = 0$ stop. Otherwise $|\mathcal{O}_n\rangle = \frac{1}{b_n} |\mathcal{A}_n\rangle$

$\longrightarrow \{b_n\}_{n=1}^{K-1}$ (Lanczos coefficients)

$\longrightarrow \{|\mathcal{O}_n\rangle\}_{n=0}^{K-1}$ (Krylov basis)

-Due to linear independence:

$$|\mathcal{A}_K\rangle = 0 \implies b_K = 0,$$

And the algorithm terminates.

Problem: numerically unstable!!

RELATION BETWEEN LANCZOS COEFFICIENTS AND CORRELATION FUNCTION

Inner product $\longrightarrow$ 2-point function: $(A | B) = \frac{1}{D} \text{Tr} [A^\dagger B]$

Autocorrelation function:

$$C(t) := \langle \mathcal{O} \mathcal{O}(t) \rangle = (\mathcal{O} | \mathcal{O}(t)) = \left(\mathcal{O} \left| e^{it\mathcal{L}} \right| \mathcal{O} \right) = \left(\mathcal{O}_0 \left| e^{it\mathcal{L}} \right| \mathcal{O}_0 \right) = \sum_{n=0}^{+\infty} \frac{(it)^n}{n!} \left(\mathcal{O}_0 \left| \mathcal{L}^n \right| \mathcal{O}_0 \right)$$

Defining moments μ_{2n} such that $\langle \mathcal{O} \mathcal{O}(t) \rangle = \sum_{n=0}^{+\infty} \frac{(it)^n}{n!} \mu_{2n}$:

$$\implies \mu_{2n} = \left(\mathcal{O}_0 \left| \mathcal{L}^{2n} \right| \mathcal{O}_0 \right) = \left[\left(\begin{array}{cccccc} 0 & b_1 & 0 & 0 & \dots & 0 \\ b_1 & 0 & b_2 & 0 & \dots & 0 \\ 0 & b_2 & 0 & b_3 & \dots & 0 \\ \vdots & & \ddots & & \ddots & \vdots \end{array} \right)^{2n} \right]_{00}$$

- $\longrightarrow$ Dyck paths.
- $\longrightarrow$ Invertible $b_n \leftrightarrow \mu_{2n}$.
- $\longrightarrow$ Analytical structure of $C(t)$.