

# Operator spectrum of NRCFTs at large charge

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Based on [Orlando, VP, Reffert '20], [VP '21]

[Hellerman, Orlando, VP, Reffert, Swanson, *to appear*]

Introduction

NRCFT at large- $Q$

Conclusion

# Motivation

Broad question: can we understand theory space?

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Broad question: can we understand theory space?

- Ambitious  $\rightarrow$  restrict to constrained theories. Here, CFTs.
- Bootstrap does an amazing job at collecting CFT data.
- Otherwise: large  $R$ -charge, large spin, etc.

# Motivation

## What if we keep adding symmetries?

- Natural guess: large internal charge  $Q$  (this talk:  $U(1)$  symmetry)
- If this limit is tractable, is there a sense in which the predictions can be universal?

# Toolbox

## Effective field theory (EFT)

In the Wilsonian picture:

- Write down all possible operators.
- Wilsonian coefficients encode UV physics.
- Specify two well-separated scales  $\Lambda_{UV}$  and  $\Lambda_{IR}$ .
- Expansion parameter:  $\epsilon \equiv \frac{\Lambda_{IR}}{\Lambda_{UV}} \ll 1$ .
- Observables given by an asymptotic series in  $\epsilon$  (or  $\epsilon^2$ )

$$\langle \mathcal{O} \rangle = \# [\alpha_1 + \alpha_2 \epsilon^2 + \alpha_3 \epsilon^4 + \dots]$$

(this talk: ignore non-perturbative effects).

# Toolbox

## Relativistic state-operator correspondence

- Operator spectrum on  $\mathbb{R}^{d+1} \leftrightarrow$  Energy spectrum on  $\mathbb{R} \times S_R^d$

$$\Delta = E \cdot R.$$

- **Goal:** compute  $\Delta_Q$  of lowest op. of charge  $Q$  via the GS energy  $E_0$ .

## Context: the relativistic $O(2)$ model

In this case:

- Fixed- $Q$  sector  $\leftrightarrow$  superfluid phase for the Goldstone.
- Scales:  $\Lambda_{IR} = \frac{1}{R}$  and  $\Lambda_{UV} = \rho^{\frac{1}{d}} = \frac{Q^{\frac{1}{d}}}{R}$ .
- Expansion parameter:

$$\epsilon = \frac{\Lambda_{IR}}{\Lambda_{UV}} = Q^{-\frac{1}{d}}.$$

- With  $Q \gg 1$ , we have  $\epsilon \ll 1$  and the EFT regime is well-defined.

Context: the relativistic  $O(2)$  model

Moreover,

$$\begin{aligned}\text{charge density} &\sim \frac{Q}{\text{Vol}} \sim \frac{Q}{R^d} \\ \text{energy density} &\sim \frac{E_0}{\text{Vol}} \sim \frac{\Delta_Q}{R^{d+1}}\end{aligned}$$

are finite, even if  $R \rightarrow \infty$ , hence

$$\Delta_Q \sim Q^{\frac{d+1}{d}}$$

to leading order.

## Context: the relativistic $O(2)$ model

We conclude that [Hellerman, Orlando, Reffert, Watanabe '15] [Cuomo '20]

$$\Delta_Q = Q^{\frac{d+1}{d}} \left[ \alpha_1 + \frac{\alpha_2}{Q^{\frac{2}{d}}} + \frac{\alpha_3}{Q^{\frac{4}{d}}} + \dots \right] \\ + Q^0 \left[ \beta_0 + \frac{\beta_1}{Q^{\frac{2}{d}}} + \frac{\beta_2}{Q^{\frac{4}{d}}} + \dots \right] + \dots,$$

in  $(d + 1)$ -dimensions.

Second line given by the Casimir energy, based on the spectrum

$$\omega_l = \sqrt{\frac{l(l+d-1)}{d}} + \mathcal{O}(Q^{-\frac{2}{d}}),$$

with multiplicity  $\frac{(2l+d-1)\Gamma(l+d-1)}{\Gamma(l+1)\Gamma(d)}$  on the  $d$ -sphere.

## Context: the relativistic $O(2)$ model

Typically,

$$\Delta_Q^{(d=2)} = \alpha_1 Q^{\frac{3}{2}} + \alpha_2 \sqrt{Q} - 0.0937 + \dots,$$

and

$$\Delta_Q^{(d=3)} = \alpha_1 Q^{\frac{4}{3}} + \alpha_2 Q^{\frac{2}{3}} - \frac{1}{48\sqrt{3}} \log Q + \alpha_3 + \dots$$

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# Nonrelativistic conformal invariance

Consider the so-called *Schrödinger* group:

- Galilean algebra
- Central extension: particle number symmetry ( $U(1)$ )
- Scale transformation:  $(t, \vec{x}) \rightarrow (e^{2\tau} t, e^\tau \vec{x})$
- SCT:  $(t, \vec{x}) \rightarrow \left( \frac{t}{1+\lambda t}, \frac{\vec{x}}{1+\lambda t} \right)$

## Nonrelativistic state-operator correspondence

Schrödinger algebra also has an automorphism such that

- Operator sp. on  $\mathbb{R}^{d+1} \leftrightarrow$  Energy sp. in  $A_0(\vec{x}) = \frac{m\omega^2}{2\hbar} |\vec{x}|^2$

[Werner, Castin '05] [Nishida, Son '07] [Goldberger, Khandker, Prabhu '14]

$$\Delta = \frac{E}{\hbar\omega}$$

(from now on:  $\hbar = m = 1$ ).

- Defines a "turning point" region  $\rightarrow$  spherical cloud/droplet.
- Focus on  $\Delta_Q$  of lowest op. of charge  $Q$  via GS energy  $E_0$ .

## Effective field theory

Let  $\chi$  be the Goldstone associated with  $Q$ .

- Fixed- $Q$  sector  $\leftrightarrow$  superfluid phase for  $\chi$  in the trap.
- IR scale: radius of the cloud  $R_{cl} \sim \frac{Q^{\frac{1}{2d}}}{\sqrt{\omega}}$ .
- UV scale: charge density (in the center)  $\rho \sim \omega^{\frac{d}{2}} \sqrt{Q}$ .
- Expansion parameter:

$$\epsilon = \frac{R_{cl}^{-1}}{\rho^{\frac{1}{d}}} \sim Q^{-\frac{1}{d}}.$$

Well-defined low-energy regime for  $Q \gg 1$ !

## Effective field theory

Since  $\text{Vol} \sim \frac{\sqrt{Q}}{\omega^{\frac{d}{2}}}$  and  $E_0 = \omega \Delta_Q$ ,

$$\text{charge density} \sim \frac{Q}{\text{Vol}} \sim \omega^{\frac{d}{2}} \sqrt{Q}$$

$$\text{energy density} \sim \frac{E_0}{\text{Vol}} \sim \frac{\omega^{\frac{d}{2}+1}}{\sqrt{Q}} \Delta_Q.$$

Then, the limit  $\omega \rightarrow 0$  tells us that

$$\Delta_Q \sim Q^{\frac{d+1}{d}}$$

to leading order.

## Intermediate conclusion

So far, it seems that

$$\Delta_Q = Q^{\frac{d+1}{d}} \left[ a_1 + \frac{a_2}{Q^{\frac{2}{d}}} + \frac{a_3}{Q^{\frac{4}{d}}} + \dots \right]$$

plus quantum corrections (Casimir energy).

- Surprising? Perhaps.
- Disappointing? Somewhat...

## Intermediate conclusion

### But what about the boundary?

- Associated with the vanishing of the charge density ("Dirichlet").
- EFT not well-defined in its proximity.

## Intermediate conclusion

In fact,

- Operators can be placed on the boundary.
- The corresponding expansion parameter is

$$\tilde{\epsilon} \sim Q^{-\frac{2}{3d}}.$$

- Contributions to  $\Delta_Q$  start at  $Q^{\frac{2d-1}{3d}}$ .

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## Results

A careful computation leads to

[Son, Wingate '05] [Kravac, Pal '18] [Orlando, VP, Reffert, '20] [VP '21]

[Hellerman, Orlando, VP, Reffert, Swanson, *to appear*]

$$\begin{aligned} \Delta_Q &= Q^{\frac{d+1}{d}} \left[ a_1 + \frac{a_2}{Q^{\frac{2}{d}}} + \frac{a_3}{Q^{\frac{4}{d}}} + \dots \right] \\ &+ Q^{\frac{2d-1}{3d}} \left[ b_1 + \frac{b_2}{Q^{\frac{2}{3d}}} + \frac{b_3}{Q^{\frac{4}{3d}}} + \dots \right] \\ &+ Q^{\frac{d-3}{3d}} \left[ c_1 + \frac{c_2}{Q^{\frac{2}{3d}}} + \frac{c_3}{Q^{\frac{4}{3d}}} + \dots \right] + \dots \end{aligned}$$

In particular

$$\Delta_Q^{(d=2)} = d_1 Q^{\frac{3}{2}} + d_2 \sqrt{Q} \log Q + d_3 \sqrt{Q} + d_4 Q^{\frac{1}{6}} - 0.2942 + \dots,$$

and

$$\Delta_Q^{(d=3)} = d_1 Q^{\frac{4}{3}} + d_2 Q^{\frac{2}{3}} + d_3 Q^{\frac{5}{9}} + d_4 Q^{\frac{1}{3}} + d_5 Q^{\frac{1}{9}} + \frac{1}{3\sqrt{3}} \log Q + d_6 + \dots$$

# Results

What I haven't told you:

- Some of the  $b$ 's contain  $\log Q$ -terms.
- Computation of the Casimir energy.
- Miraculous connection with experiments!

# Outlook

- Include spin [Kravec, Pal '19]
- Gravity dual [Son '08] [Balasubramanian, McGreevy '08]
- BCS-BEC crossover
- Non-Abelian  $Sp(N)$  at large- $N$   
[Veillette, Sheehy, Radzihovsky '06] [Sachdev, Nikolic '06]